

Perovskite semiconductor-engineered cascaded molecular energy levels in naturally-sensitized photoanodes

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ABSTRACT

Naturally-sensitized photoanodes in dye-sensitized solar cells (DSSCs) are promising alternatives to enhance photoabsorption, electron excitation/injection, but voltage loss remains a challenge. Here, we focus on understanding the cascading of energy levels in perovskite semiconductor cosensitized naturally-sensitized photoanodes to leverage forward charge transport addressing the voltage loss arising from ITO/TiO₂ heterojunction's built-in potential. The β -carotene-sensitized TiO₂ photoanode modified with methylammonium lead iodide (MAPbI₃) co-sensitizer causes an upward shifting in TiO₂ Fermi level (E_F). This phenomenon is predominantly attributed to increased initially injected electrons due to low MAPbI₃ bandgap and high visible-light absorption. Enhanced charge separation and injection mechanisms at the TiO₂/MAPbI₃ interface increase the effective density-of-states ($DOS > 2.46 \times 10^{21} \text{ cm}^{-3}$) in the TiO₂ conduction band (CB) and hence decrease its work function to 4.82 eV. The decrease in TiO₂ work function suppressed CB bending at ITO/TiO₂ heterojunction, which minimized the photoinduced electrostatic potential barrier up to 13.1%. The reduced Schottky barrier ($\phi_{SBH} < 0.52 \text{ eV}$) only allows electrons tunneling, while inhibited back-electron transport reduced both current leakage and voltage loss yielding in high open-circuit voltage (V_{oc} increase by 120%) and power conversion efficiency (PCE increase by 240%). The MAPbI₃ incorporation also broadened photoanode absorbance by 2-fold, paving the way towards perovskite semiconductor cosensitization to avoid voltage loss from bio-integrated photoanodes for photovoltaic and other optoelectronic and photonic applications. Future works will focus on studying the series resistance, cathode electrode coating uniformity, recombination kinetics, solid electrolytes, as well as other aspects typically related to increasing the photocurrent levels of this β -carotene-sensitized solar cell.

1. Introduction

Photosensitization is the basis for designing efficient dye-sensitized solar cells (DSSCs) capable of initiating electron injection and charge transfer from the LUMO energy level of dye molecules to the conduction band (CB) of the wide band gap (WBG) semiconductor [1–3]. Different kinds of natural and environmental-friendly photosensitizers [4] extracted from light-harvesting (LH) complexes of anthocyanin, carotenoid, and chlorophyll biomolecules [5] have emerged as an attempt to substitute the expensive and toxic [6] ruthenium polypyridyl dyes [7]. However, the power conversion efficiency (PCE) of bio-sensitized DSSCs is low (<0.05–1.7%) [8–10] and requires a thorough understanding of the role of pigment's molecular structure, interfacial interactions,

anchoring groups, and conjugated double bonds (n) in dye attachment, improved carriers transport, and decreased recombination [11–13]. Surface hydroxylation [14–16], perovskite cosensitization [17–20], and/or absorbance expansion were previously utilized to improve electron injection and photoanode sensitization, respectively. Enhanced photoanode sensitization and broadened photoanode absorption spectrum [21] is attainable by stacking semiconducting materials [22–24] with different bandgaps, or more viably by utilizing low-bandgap energy cosensitizers [17–20]. Methylammonium lead iodide (MAPbI₃) perovskite semiconductors [25,26] are appropriate for cosensitization functionalities owing to their high visible-light absorption ($\alpha = 0.7 \times 10^7 \text{ cm}^{-1}$) [27], composition-engineered bandgap tunability ($E_g = 1.5\sim 2.43 \text{ eV}$) [28], high charge carriers mobility ($>5 \text{ cm}^2 \text{ V}^{-1} \text{ S}^{-1}$) [29], long diffusion lengths ($>100 \text{ nm}$) [30], and easy processability

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Nomenclature**Abbreviation/unit Jargon**

DSSCs	Dye-sensitized solar cells
LH	Light-harvesting
PCE (η)/%	Power conversion efficiency
MAPbI ₃	Methylammonium lead iodide
D- π -A	Donor- π -acceptor
CB	Conduction band
V_{oc} /V	Open circuit voltage
ϕ_{SBH} /eV	Schottky barrier height
WBG	Wide bandgap
MAI	Methylammonium iodide
CEs	Counter electrodes
ETL	Electron transport layer (TiO ₂)
P	Perovskites (MAPbI ₃)
HTL	Hole-transport layer
TE	Thermionic emission
J-V	Current density-voltage
n	Diode ideality factor
RFs	Rectification factors

J_0 /nA	Dark saturation current (leakage)
FF	Fill factor
R_s /k Ω	Series resistance
R_{sh} /k Ω	Shunt resistance
J_{sc} /μA cm ⁻²	Short circuit current density
J_1 /μA cm ⁻²	Apparent current density in both TiO ₂ and MAPbI ₃
J_2 /μA cm ⁻²	Total apparent current density in the electrolyte solution [I^-/I_3^-]
α /cm ⁻¹	Absorption coefficient
$h\nu$ /eV	Photon energy
E_g /eV	Bandgap
t /mm	Thickness of the quartz cuvette
OD /cm ⁻¹	Optical density
λ_{max} /cm ⁻¹	Maximum absorption wavelength
V_1 /V	Voltage loss
E_F /eV	Fermi level
$[I^-/I_3^-]$	Liquid redox electrolyte
ϕ_{TiO_2} /eV	TiO ₂ work function
THF	Tetrahydrofuran

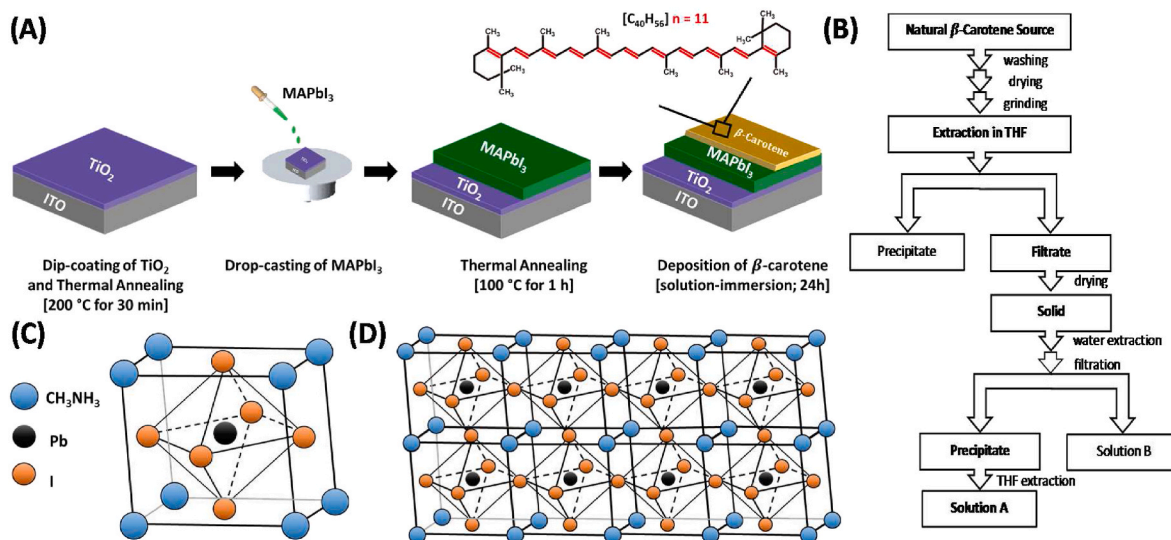


Fig. 1. (A) Experimental steps for the fabrication of [ITO/TiO₂/MAPbI₃/β-carotene] photoanode heterostructure; (B) Flowchart showing the THF-extraction procedure and step-wise purification for β-carotene; (C) Organic-inorganic halide perovskite (MAPbI₃) unit cell; (D) Ideal cubic structure for MAPbI₃ when (tolerance factor = 1) indicating a perfect cation fit.

[31,32]. Combining CH₃NH₃PbX₃ (X = Br or I) in a TiO₂-based photoanode has shown its potential to achieve high efficiency with Ru-dyes in a liquid-state (3.81%) [17], quantum-dot-sensitized (6.54%), and a solid-state mesoscopic DSSC (15%) [33]. However, such cosensitization has not been yet studied with natural biomolecular dyes, according to the authors' knowledge, for potential improvements in bio-sensitized DSSCs.

In this work, we focus on understanding the role of perovskite cosensitization and its interfacial interaction with TiO₂ and β-carotene in leveraging forward charge transport. Such cosensitization can address the voltage loss arising from ITO/TiO₂ heterojunction's built-in potential in bio-sensitized DSSCs. We elucidated the photoabsorption mechanisms [34] pertaining cosensitization of β-carotene pigments. The integration of organic-inorganic MAPbI₃ halide perovskites with β-carotene dyes have been studied using a novel naturally-sensitized

photoanode in [ETL/Perovskite/HTL] configuration in liquid-state DSSCs. Perovskite cosensitizers improved PCE via enhancing visible-light absorption, excitons generation, electron injection, and reducing voltage loss at ITO/TiO₂ junction. Broadening of the absorbance spectrum is expected for the cosensitized photoactive layers increasing the probability of absorbing more photons within (400–700 nm) [35] from the available 42% visible-light energy in solar radiation [36]. Experimental procedures for the design of [ITO/TiO₂/MAPbI₃/β-carotene] photoanode are shown in Fig. 1-A. The THF-solvent dye extraction and step-wise purification of non-treated lab-grade β-carotene powders (Sigma Aldrich, USA) are illustrated in Fig. 1-B, to obtain "Solution A" which was used for β-carotene deposition atop MAPbI₃. The MAPbI₃ unit cell and crystal structure consist of a cubic array with A-site corners sharing BX₆ octahedra are shown in Fig. 1-C and Fig. 1-D.

2. Carotenoids and perovskites cosensitization

Poly-isoprenoid carotenoids (e.g. xanthophylls, lycopene, and β -carotene) hold great potential as natural photosensitizers in DSSCs [11,21]. Carotenoids contain double π -conjugated bonds ($n > 11$) responsible for visible-light photoexcitation (400–550 nm and 620–680 nm) [37] of free electrons to higher energy levels (π to π^*) [38]. However, donor– π –acceptor (D– π –A) conjugated carotenoid dye molecules [9,39] have poor surface interactions and/or chemisorption onto TiO_2 . The poor dye attachment hinders effective electron injection into the TiO_2 CB resulting in low PCE for carotenoids $< 0.58\%$ [8] ($< 0.074\%$ for β -carotene [9]). Onodera et al. showed that the open-circuit voltage (V_{oc}) decreased with increasing Schottky barrier height (ϕ_{SBH}) at the interface of TiO_2 and transparent conducting oxide (TCO) [40–42]. In DSSCs, TCO is the base substrate on which the cell is built and it allows the light to enter into the solar cell. TCOs used in DSSCs are indium tin oxide (ITO), fluorine-doped tin oxide (FTO), aluminum-doped zinc oxide (AZO), and gallium-doped zinc oxide (GZO). In this study, when TiO_2 comes in contact with ITO, band bending in TiO_2 CB occurs influencing the current-loss mechanism from inhibited electron collection. Perovskite cosensitization can promote the β -carotene sensitized-photoanode light-harvesting ability addressing both TiO_2 CB bending and voltage loss problems. This is because perovskite cosensitization mitigates interfacial recombination and CB bending in TiO_2 (reduced ϕ_{SBH}) due to the $\text{CH}_3\text{NH}_3\text{PbI}_3/\text{TiO}_2$ interactions. However, weak dye attachment lowers the performance of bio-DSSCs from electron recombination at the dye/semiconductor interface [8,21,43]. One may be able to study changes in the MAPbI_3 film and/or β -carotene anchoring either by hydrogen bonding or chemisorption from the FT-IR analysis of the available bonds. Interfacial recombination with the oxidized dye or oxidized redox species in the electrolyte [44] facilitates current leakage that is the main limitation for improving PCE of naturally-sensitized DSSCs [11]. Missing and/or few dye anchoring groups ($-\text{OH}$ or $-\text{COOH}$) in the acceptor segments [12] lead to high series resistance and reduced current density ($J_{sc} \sim \mu\text{A}/\text{cm}^2$) [45]. These observations contribute in increasing the voltage loss from ϕ_{SBH} at the ITO/ TiO_2 interface (n-type semiconductor/metal) resulting in internal resistance and recombination current with the electrolyte. Thus, strong van der Waals interactions possibly introduced by MAPbI_3 ensure fast electron injection completed within 100 fs (with possible improved dye attachment on MAPbI_3), where $\text{CH}_3\text{NH}_3\text{PbI}_3/\text{TiO}_2$ covalent bonding should facilitate photoexcitation of charge-separated states leading to ultrafast interfacial charge separation prior to electron injection [46].

Recent studies (2018–2020) showed the potential of using carotenoids including β -carotene as biomolecular sensitizers in photoanode structures. For instance, Zanjanchi and Beheshtian (2019) showed in their recently published theoretical work, based on density functional theory (DFT), that carotenoids are considered as an optimal natural pigment class to achieve high open-circuit photovoltage [47]. Further, experiments have backed-up this theory showing the possibility to have a PCE as high as 2% from using a cocktail of naturally-sensitized dyes (e.g. chlorophyll, betanins, carotenoids, anthocyanins, and tannins), which enhanced photoanode visible-light absorption abilities [48]. Moreover, the use of cobalt solid-state redox with novel metal-free dihexyloxy-substituted triphenylamine (DHO-TPA) dye has shown a significant improvement in dye attachment and cell stability allowing achievement of a cell performance of up to 10.3% [49]. However, when focusing on studying β -carotene alone as a natural β -carotene photosensitizer in DSSCs, recent works from the last three years showed a maximum PCE of 0.074% in $[\text{FTO}/\text{TiO}_2/\beta\text{-carotene}/[\text{I}^-/\text{I}_3^-]/\text{Pt}/\text{FTO}]$ with a current density of $270 \mu\text{A cm}^{-2}$ and $V_{oc} = 400 \text{ mV}$ [9]. This relatively high PCE from β -carotene only (i.e. considered high when compared to DSSCs made of natural pigments) was achieved from using EL-HSE (triiodide/iodide imidazole compound in 3-methoxypropionitrile for high stability redox, Dyesol) interfaced with β -carotene dye molecules with a

bandgap around 2.47 eV, where the dye was extracted from Syngonium leaf [9]. Another study (2018) showed that β -carotene extracted from yellow sweet potato was capable of absorbing visible-light energy in the range 400–550 nm within a designed $[\text{FTO}/\text{TiO}_2/\beta\text{-carotene}/[\text{I}^-/\text{I}_3^-]/\text{Pt}/\text{FTO}]$ DSSC heterostructure achieving a maximum PCE of 0.057% [6]. A much recent study (2019) used readily available lab-grade β -carotene powders in a regular $[\text{ITO}/\text{TiO}_2/\beta\text{-carotene}/[\text{I}^-/\text{I}_3^-]/\text{C}/\text{ITO}]$ liquid-redox-based DSSCs which showed PCE of $1 \times 10^{-4}\%$ [50], meanwhile, Supriyanto et al. (2018) found that extracted β -carotene dyes from orange fruit and tomatoes applied in $[\text{FTO}/\text{TiO}_2/\beta\text{-carotene}/[\text{I}^-/\text{I}_3^-]/\text{Pt}/\text{FTO}]$ would show PCEs of 0.02% and 0.03%, respectively, with maximum absorption wavelength around 441 and 466.5 nm [21]. Furthermore, the blending of chlorophyll with carotenoids has shown a much improvement in carotenoids sensitizing function in DSSCs since chlorophyll provides intensified visible-light absorption capabilities with enhanced layer protection [51,52] and formed radical cations [53]. Protein complexes (LH2, BR, RC) and chlorophyll-a combined with carotenoids showed a high DSSCs PCE of 0.16–0.57% and 4%, respectively.

3. Materials and methods

3.1. TiO_2 deposition (dip-coated or spin-coated) on ITO photoanodes

ITO substrates (1" transparent and conductive sheets with sheet resistance of $\sim 23 \Omega/\square$) were sonicated in IPA for 15 min and followed by deionized (DI) water washing. Scotch tapes were applied on the conductive sides to have a $2 \text{ cm} \times 1.5 \text{ cm}$ exposed surface area ready for TiO_2 dip-coating (or spin-coating) with no prior treatments for TiO_2 nanoparticles paste (Sigma Aldrich, USA). With taped ITO backsides, ITO sheets were dip-coated (or spin-coated @ 3000 rpm) successively, removed immediately, and dried in ambient air for 5 min before annealing the ITO/ TiO_2 structure at 400°C for 30 min [54]. TiO_2 crystals forming above $>270^\circ\text{C}$ may induce a significant increment of R_s of ITO layer yielding in extremely low photocurrents. Crystallized TiO_2 nanoparticles atop ITO films were removed from the hot plate, rinsed with DI water, and stored in a dry petri dish.

3.2. Synthesis of MAPbI_3 and deposition

MAPbI_3 was synthesized by mixing 2 g of PbI_2 powders with 0.5 g of MAI powders in a glass beaker with continuous addition of 5 mL of DMF under standard conditions [31,32]. The mixed solution was then heated on a hot plate for 1 h at 60°C under vigorous stirring. The obtained precursor was drop-casted (2–3 drops) onto sintered ITO/ TiO_2 electrodes and crystallized at 100°C for 1 h to obtain the ITO/ TiO_2 / MAPbI_3 heterostructure.

3.3. THF-extraction of β -carotene sensitizers and deposition

Since the solubility of β -carotene carotenoids is greatest in tetrahydrofuran (THF), 10,000 mg/L [55], among other organic solvents, THF solvent was utilized for dye extraction and step-wise purification of non-treated lab-grade β -carotene powders (Sigma Aldrich, USA) [56]. β -carotene powders (1.5 mg) were added to THF organic solvent (50 mL) under continuous stirring for 1 h in ambient conditions. Dye extraction was confirmed from the observed color change in the THF (yellowish/reddish). THF solvent was then vacuum filtered, dried, and washed with DI water following the step-purification using the procedure shown in Fig. 1(B) to obtain THF– H_2O –extracted carotenoid precipitates. Purified carotenoids were then washed again with THF (50 mL) to obtain "Solution A" used for dye deposition. Solution A was concentrated at 60°C (THF volume reduced to 25 mL). All sintered photoanodes were soaked in THF– β -carotene concentrated solution for $>24 \text{ h}$ to ensure dye deposition and molecules anchoring onto TiO_2 via covalent bonding

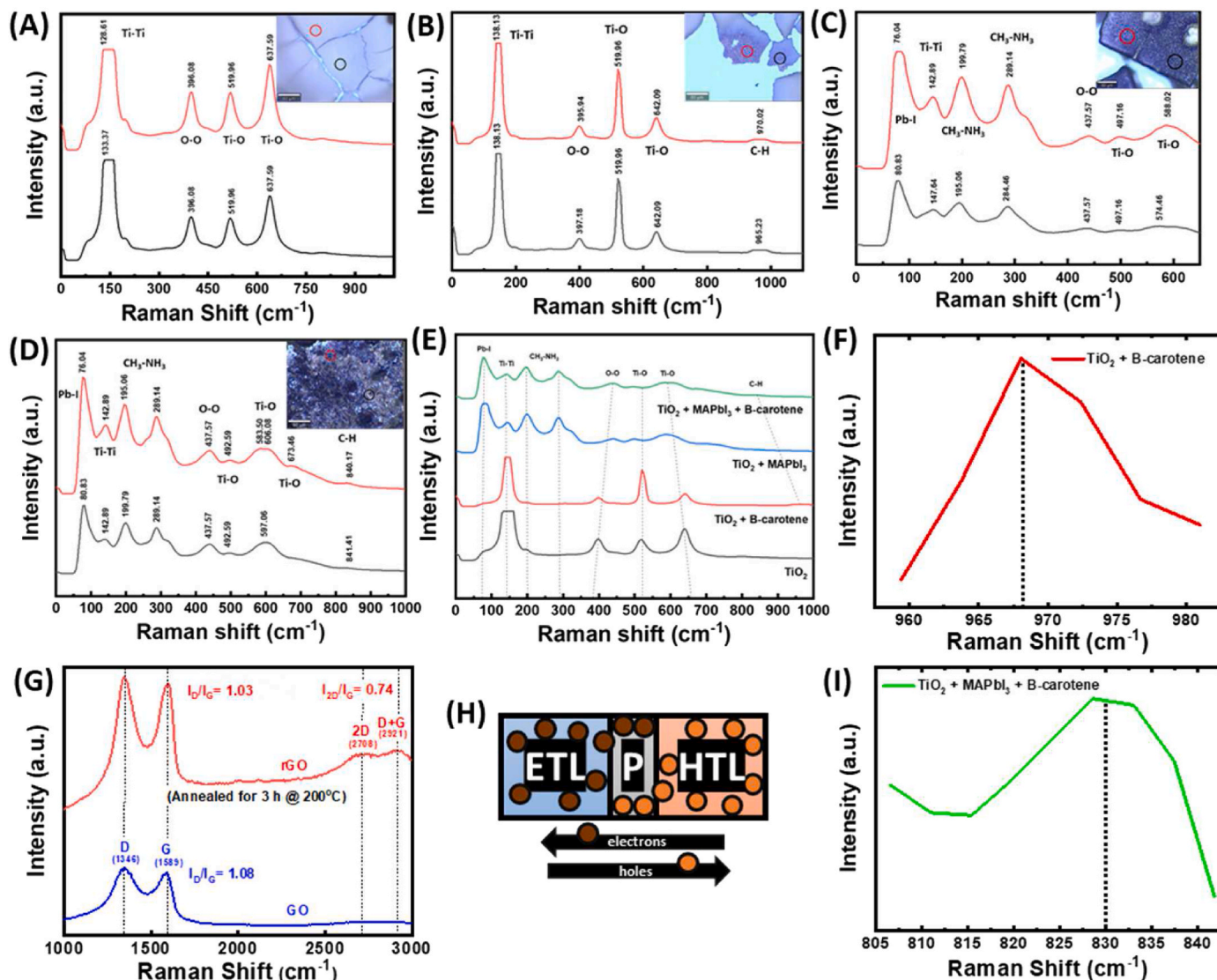


Fig. 2. Single-spectrum Raman spectroscopy characterization of different nanostructured naturally-sensitized photoanodes: (A) ITO/TiO₂; (B) ITO/TiO₂/β-carotene; (C) ITO/TiO₂/MAPbI₃; and (D) ITO/TiO₂/MAPbI₃/β-carotene, insets in A-D: optical images and Raman scan locations; (E) Comparison of Raman bands observed in the four photoanode surfaces showing energy-shifts (red- or blue-shifted) in TiO₂ and β-carotene when combined with MAPbI₃ co-sensitizer; (F) and (I) Zoom-in on the (ν₃) vibrational bands of β-carotene from C-H rocking before and after MAPbI₃ interfacing; (G) Single-spectrum Raman spectroscopy characterization for the synthesized coal-derived GO and the annealed ITO/GO yielding in rGO-cathode electrodes; (H) Schematic showing the possible excitons generation and electron/hole transport layers involved in perovskite-co-sensitized DSSCs; ETL: electron transport layer (TiO₂); P: Perovskites (MAPbI₃); HTL: hole-transport layer (β-carotene/[I⁻/I₃⁻]).

and/or possible chemisorption.

3.4. Synthesis of coal-derived rGO for cathode

NaNO₂ was used as a weak oxidant to avoid structural damages while breaking the van der Waals interlayer forces that exist between the coal-based graphitic layers. 2.5 g of coal sample was oxidized with 2 g of NaNO₂ in 100 mL of H₂SO₄ (98%). The suspension was agitated for 24 h at room temperature conditions; followed by heating up to 80 °C and stirring for 6 h. HNO₃ (1 M) was added to the suspension and then diluted with DI water (volume ratio of 1:2) to promote chemical exfoliation. The resultant graphene oxide (GO) suspension was centrifuged (1000 rpm) for 15 min to discard the undesired ions. The supernatant was dispersed in DI water and centrifuged (12,000 rpm) to collect precipitates (3-times) and to remove the remaining impurities. The counter electrode (cathode) was created by drop-casting (2–3 drops) of the coal-derived GO suspension on the conductive side of ITO. Drop-casted GO

(post-deposited on ITO) was then chemically reduced in a vacuum oven by annealing at 200 °C for 3 h to remove the undesired oxygen (=O), hydroxyl (–OH), and carboxyl (–COOH) functional groups yielding in reduced graphene oxide (rGO) blackish films.

3.5. Cell fabrication and photovoltaic measurement

Four DSSC cell designs are considered for the performance analysis: (a) ITO/TiO₂/[I⁻/I₃⁻]/rGO/ITO; (b) ITO/TiO₂/β-carotene/[I⁻/I₃⁻]/rGO/ITO; (c) ITO/TiO₂/MAPbI₃/[I⁻/I₃⁻]/rGO/ITO; and (d) ITO/TiO₂/MAPbI₃/β-carotene/[I⁻/I₃⁻]/rGO/ITO; as illustrated in Figure S4. A drop of [I⁻/I₃⁻] liquid redox electrolyte was drop-casted on the TiO₂-photoanode. Both photoanode and cathode were brought together while keeping the conductive ITO sides facing each other with a careful line-up of TiO₂-anode on rGO-cathode. Once the line-up was satisfied, binder clips were applied on the cell sides to hold the DSSCs in place. An offset was kept between the two electrodes to connect Cu-tapes electrical

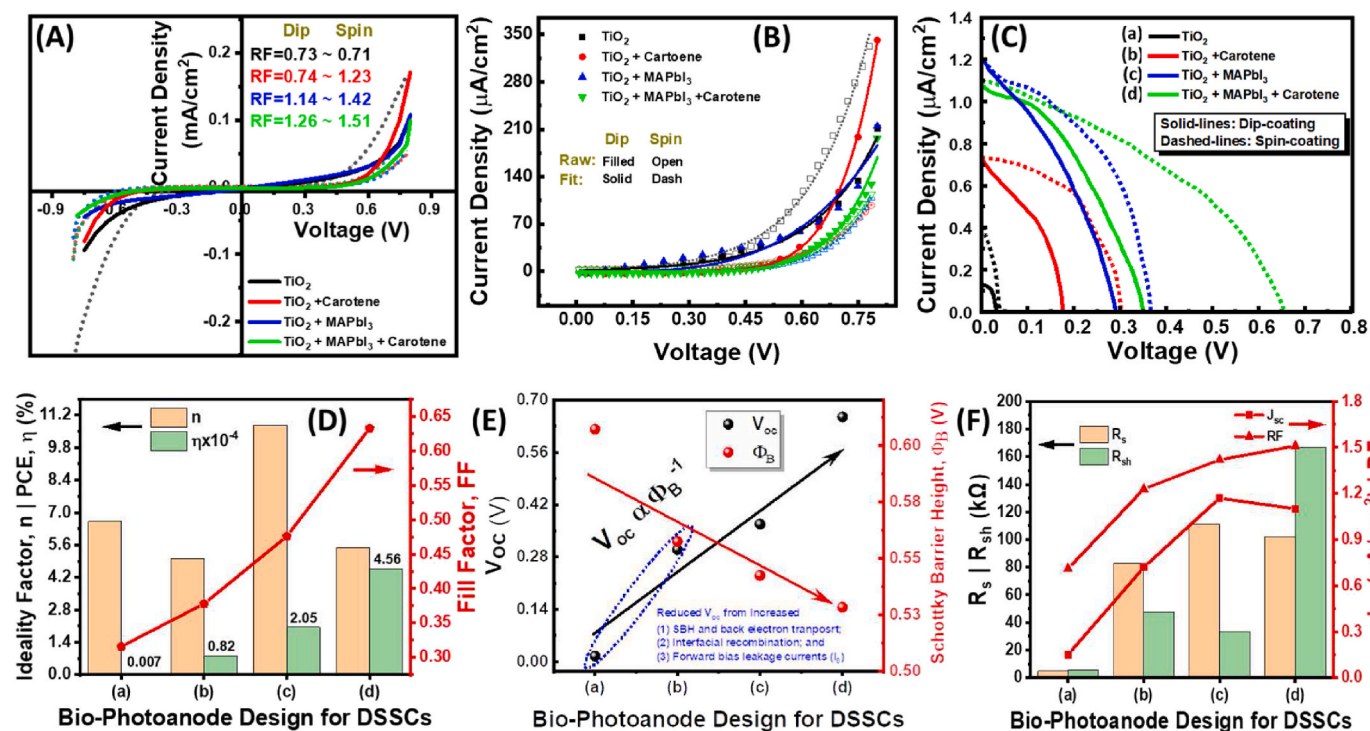


Fig. 3. Different nanostructured photoanodes for β -carotene-sensitized and/or MAPbI₃ cosensitized bio-DSSCs: (A) Diode behavior and rectification factors (RFs) with highest RF \approx 1.5 reserved for spin-coated photoanode with sensitized/cosensitized DSSC reducing dark saturation currents (J_0); (B) Exponential fitting of forward-current using the TE diode model for analyzing diode ideality (n) factors; (C) J-V curves and photovoltaic effect with highest V_{oc} \approx 0.66 V reserved for DSSC with spin-coated sensitized/cosensitized photoanode; (D) Calculated ideality factor (n), power conversion efficiency (PCE), and fill factor (FF) for DSSCs with different spin-coated photoanodes (a)–(d); (E) Relationship between V_{oc} and ITO/TiO₂ ϕ_{SBH} in spin-coated samples showing reduced ϕ_{SBH} from addition of MAPbI₃ co-sensitizers resulting in high built-in potential and upward shifting of TiO₂ Fermi level; (F) R_s , R_{sh} , J_{sc} , and RF observed in DSSCs with different spin-coated photoanodes (a)–(d).

contacts at the positive and negative terminals creating a complete electric circuit ready for current-voltage (J-V) analysis [54]. Diode and photovoltaic measurements were carried out in the room temperature conditions with a source meter (Keithley 2612); current density versus voltage (J-V) measurements of designed solar cells were performed under 100 mW/cm² illuminations (AM1.5G) using a solar simulator.

3.6. Raman spectroscopy

Raman spectroscopy characterization of photoactive electrode materials was performed using a confocal Raman microscope (Raman-AFM, WITec alpha 300 RA, laser wavelength of 532 nm). Single-spectrum analyses were carried out using a laser spot-size of about 700 nm and 20X objective lens.

3.7. Absorption spectrum of MAPbI₃, TiO₂, and β -carotene

Absorption spectrums of photoactive materials were measured using a UV-Vis spectrophotometer (SSEYL UV-5100, \pm 2 nm). DI water was used as a reference solution for the baseline of the absorption spectrum of MAPbI₃ and TiO₂, whereas THF-solvent was used as a reference solution for β -carotene absorption measurements. First, DI-water (10 mL) reference solution spectrum was measured for baseline correction. Then, a drop of MAPbI₃ or TiO₂ was added to each of the two separate cuvettes filled with DI (10 mL), located in front of the UV-Vis light to measure the optical absorption spectrum. Similarly, a drop of β -carotene was added to THF-solvent (10 mL) reference solution for the β -carotene absorption spectrum analysis. Three runs of each solution were measured for accuracy. Lastly, calculations of the cumulative fit peak for all the above three solutions' absorption spectrum were performed to estimate the photoanode optimal wavelength absorption abilities.

4. Results and discussion

Raman spectroscopy characterization (532 nm laser excitation) of bulk TiO₂ in the anatase phase reveals evident vibrational modes at 638, 520, 396, and a very distinct peak at 129 cm⁻¹ (Fig. 2) confirming the deposition of TiO₂ layer on ITO substrate. The space group D_{4h}¹⁹ is attributed to anatase phase with assumed site symmetries for the Ti and O atoms within the unit cell (D_{2d} for Ti; C_{2v} for O). The group-theory analysis shows six Raman-active "lattice vibrations" assigned as follows: A_{1g} (517 cm⁻¹), B_{1g} (640 cm⁻¹), B_{1g} (397 cm⁻¹), E_g (640 cm⁻¹), E_g (147 cm⁻¹), and E_g (197 cm⁻¹) [57]. The distinct peaks at 128~134 cm⁻¹ indicate the vibrations of Ti-Ti bonding in the octahedral chains. The band at 396 cm⁻¹ represents the O-O covalent interactions in TiO₂, where covalence/length/frequency relation calculations confirm the Ti-O band occurrence at 638 and 520 cm⁻¹ [58]. Broader and shifted peaks were observed in Fig. 2-E for Ti-O (638 cm⁻¹) and O-O (396 cm⁻¹) vibrations from the addition of MAPbI₃ and/or β -carotene photosensitizers. The O-O band is blue-shifted while the Ti-O band is red-shifted indicating changes in the energy requirements for Raman-active polarizable vibrations. In Fig. 2-(C-E), we observed the Raman spectra of MAPbI₃ with two prominent peaks: (i) the vibrational mode at 76~80 cm⁻¹ (theoretically ~110 cm⁻¹) is attributed to the Pb-I stretching and bending of the Pb-I-Pb bond with either A_g or B_{2g} symmetry, and (ii) the bands at 200 and 300 cm⁻¹ (theoretically ~250 cm⁻¹) indicate the liberations of the MA cations (MA torsional modes) [59,60]. No red-shifts were observed in perovskite peaks when MAPbI₃ was interfaced with other photoanode materials confirming stress-free/undeformed perovskite crystals and preserved MAPbI₃ photosensitivity. Thampany et al. [61] proved that the use of [MAPbI₃ + TiO₂] ~ MAPbI₃ provides the highest interfacial stability, indicated by decomposition temperature, as compared to [MAPbI₃ + SnO₂] >

Table 1

Determined diode parameters and φ_{SBH} at the ITO/TiO₂ interface in the different designed DSSCs using various naturally-sensitized and/or cosensitized photoanode architectures.^a

(SN) Photoanode Architecture	Coating	J ₀ (μA)	n	RF	φ_{SBH} (V)	J ₁ (pA)	J ₂ (μA)
(a) TiO ₂	Dip	2.280	5.97	0.73	0.607	0.274	0.150
	Spin	1.870	6.62	0.71	0.607	0.143	0.150
(b) TiO ₂ + β -carotene	Dip	0.364	5.40	0.74	0.559	N/A	N/A
	Spin	0.002	5.00	1.23	0.557	N/A	N/A
(c) TiO ₂ + MAPbI ₃	Dip	0.006	4.03	1.14	0.542	423.2	1.330
	Spin	0.041	10.76	1.42	0.542	3.988	1.170
(d) TiO ₂ + MAPbI ₃ + β -carotene	Dip	0.173	4.67	1.26	0.544	N/A	N/A
	Spin	0.081	5.47	1.51	0.528	N/A	N/A

The bolded numbers refer to the optimal found design parameters for achieving high solar cell efficiency based on the selected photoanode and coating method.

^a $R^2 > 0.99$ in the different cells using TE diode equation against curves exponential fitting [$y = a \exp(bx)$]; φ_{SBH} at ITO/TiO₂ from Eq. (S7) in Supplementary (see section 4) from the measured J_{sc} and V_{oc} = V₁ reported in Table 2; J₁ is the apparent current density in both TiO₂ and MAPbI₃, J₂ is the total apparent current density in the electrolyte solution [I^-/I_3^-] from electron diffusion, calculated @ corresponding reported diffusion coefficients (D_e) and charge/electron densities or number of electrons (n_e) of materials from literature (Supplementary, see section 4).

[MAPbI₃ + NiO] contacts. Methylammonium iodide (MAI) DFT calculations and experiments showed the strongly bound MAPbI₃ molecules on the TiO₂ surface preventing further reactions [61]. This comes in agreement with the observed stress-free/undeformed MAPbI₃ crystals deposited on TiO₂ suggesting equivalent MAPbI₃ bulk decomposition rates for [MAPbI₃ + TiO₂] ~ MAPbI₃ for reduced interfacial recombination and improved device stability. It turned out that the interactions between perovskite and TiO₂ at 0 K are purely van der Waals, whereas I-Ti covalent bonds can form transiently at the interface at ambient temperatures according to time-domain DFT calculations. The strong interactions between the donor and acceptor species at CH₃NH₃PbI₃/TiO₂ interface ensure fast electron injection completed within 100 fs [46]. The covalent bonding contributes to the high efficiencies of perovskite-sensitized TiO₂ solar cells due to facilitated photoexcitation of charge-separated states and adiabatic electron transfer leading to ultrafast interfacial charge separation [46].

The Raman spectra of β -carotene show a weak band at 970 cm⁻¹ (v₃) due to the in-plane rocking modes of the CH₃ groups attached to the polyene chain and C-H bending, Fig. 2-(B-E) [62,63]. Interestingly, when β -carotene was combined with either TiO₂ and MAPbI₃, we observed that there is a strong red-shift of about 130 cm⁻¹ in the weak band of β -carotene as shown in Fig. 2-(F,I). This red-shifting might indicate stresses in the β -carotene structure induced by hydrogen-bonding and/or dipole-dipole van der Waals interactions of C-H bond with MAPbI₃ crystal lattice.

To build a cost-effective bio-DSSC, we have further employed coal as a natural carbon source to create rGO-based counter electrodes (CEs). Coal-derived GO and rGO CEs were characterized from the D-peak at 1346 cm⁻¹, which corresponds to the A_{1g} breathing vibration mode, activated by the structural defects and disordered sp² domains in graphene lattice, the G-peak at 1589 cm⁻¹, which corresponds to the in-plane E_{2g} vibration mode of neighboring carbon atoms of sp²-hybridized C=C bonds, and the 2D peak at 2708 cm⁻¹, which corresponds to the overtone of the D-band resulting from the double-resonance two-phonon in-elastic scattering process [64,65]. The intensity ratio of 2D

and G-bands (I_{2D}/I_G ≈ 0.74) confirming the existence of a multi-layered graphene structure (Fig. 2-G). The charge transport layers involved in perovskite-cosensitized bio-DSSCs, modified with [ETL/Perovskite/HTL] configuration, are shown in Fig. 2-H.

To quantify the diode and photovoltaic parameters, thermionic emission (TE) [66,67] and Cheung's methods [68] are typically used. The cell diode parameters were determined from the exponential fitting of dark current density-voltage (J-V) data (dark current) following the

TE model: $J_D = J_0 \exp\left(\frac{qV}{nk_B T}\right)$, (Fig. 3-B), where symbols definitions can

be found in Supplementary (see sections 1 and 2). The diode ideality factor (n) approached unity by MAPbI₃ and/or β -carotene photoanode sensitization reaching optimal n = 4–5, due to reduced current leakage (J₀). Dark saturation currents (J₀) are found to approach the nano-amperes (nA) range when adding MAPbI₃ and/or β -carotene yielding less leakage current as shown in Table 1 or Fig. 3-A and Fig. 3-D for spin-coated photoanode cells. A lower n value indicates better electron transport from reduced recombination at the interfacial contacts. Reproducibility of the diode curves was obtained for the same photoanode (see Supplementary, Figure S1). Sensitized/cosensitized DSSCs diodes showed the highest RF ≈ 1.2–1.5 indicating a two-fold increase in [forward/reverse] photocurrents compared to only-TiO₂ photoanodes (Fig. 3-A). The φ_{SBH} of ITO/TiO₂ junction was also diminished matching the expected high V_{oc} (Fig. 3-E). However, the very low photocurrents ~ μA/cm² in fabricated devices might be correlated to the perovskite crystals decomposition (suggesting higher MAPbI₃ bandgap). Such decomposition of deposited layers is possible due to the introduced interfaces in the photoanode structure: (i) TiO₂/MAPbI₃, (ii) TiO₂/MAPbI₃/dye, (iii) TiO₂/MAPbI₃/electrolyte, and (iv) TiO₂/MAPbI₃/dye/electrolyte. The complicated structure would alter the existing material's structural bonds yielding possible films degradations and/or corroded MAPbI₃ from being in contact with the electrolyte. In an optimal photoanode heterostructure, the MAPbI₃ should be completely sandwiched by the deposited TiO₂ nanoparticles and β -carotene to avoid undesired interactions between the MAPbI₃ layer

Table 2

DSSCs photovoltaic parameters in DSSCs using different photoanode heterostructures.

(SN) Photoanode Architecture	Coating	R _s (kΩ)	R _{sh} (kΩ)	V _{oc} (V)	J _{sc} (μA/cm ²)	FF	η (%)
(a) TiO ₂	Dip	<33.3	55.56	0.015	0.15	0.24	5.47 × 10 ⁻⁷
	Spin	<4.76	5.56	0.015	0.15	0.32	7.11 × 10 ⁻⁷
(b) TiO ₂ + β -carotene	Dip	83–129	37.04	0.171	0.80	0.30	4.16 × 10 ⁻⁵
	Spin	83–129	47.62	0.300	0.72	0.38	8.22 × 10 ⁻⁵
(c) TiO ₂ + MAPbI ₃	Dip	55–84	33.33	0.290	1.33	0.33	1.29 × 10 ⁻⁴
	Spin	111–168	33.33	0.368	1.17	0.48	2.05 × 10 ⁻⁴
(d) TiO ₂ + MAPbI ₃ + β -carotene	Dip	51–83	41.67	0.341	1.14	0.36	1.40 × 10⁻⁴
	Spin	102–167	166.67	0.655	1.10	0.63	4.56 × 10⁻⁴

The bolded numbers refer to the optimal found design parameters for achieving high solar cell efficiency based on the selected photoanode and coating method.

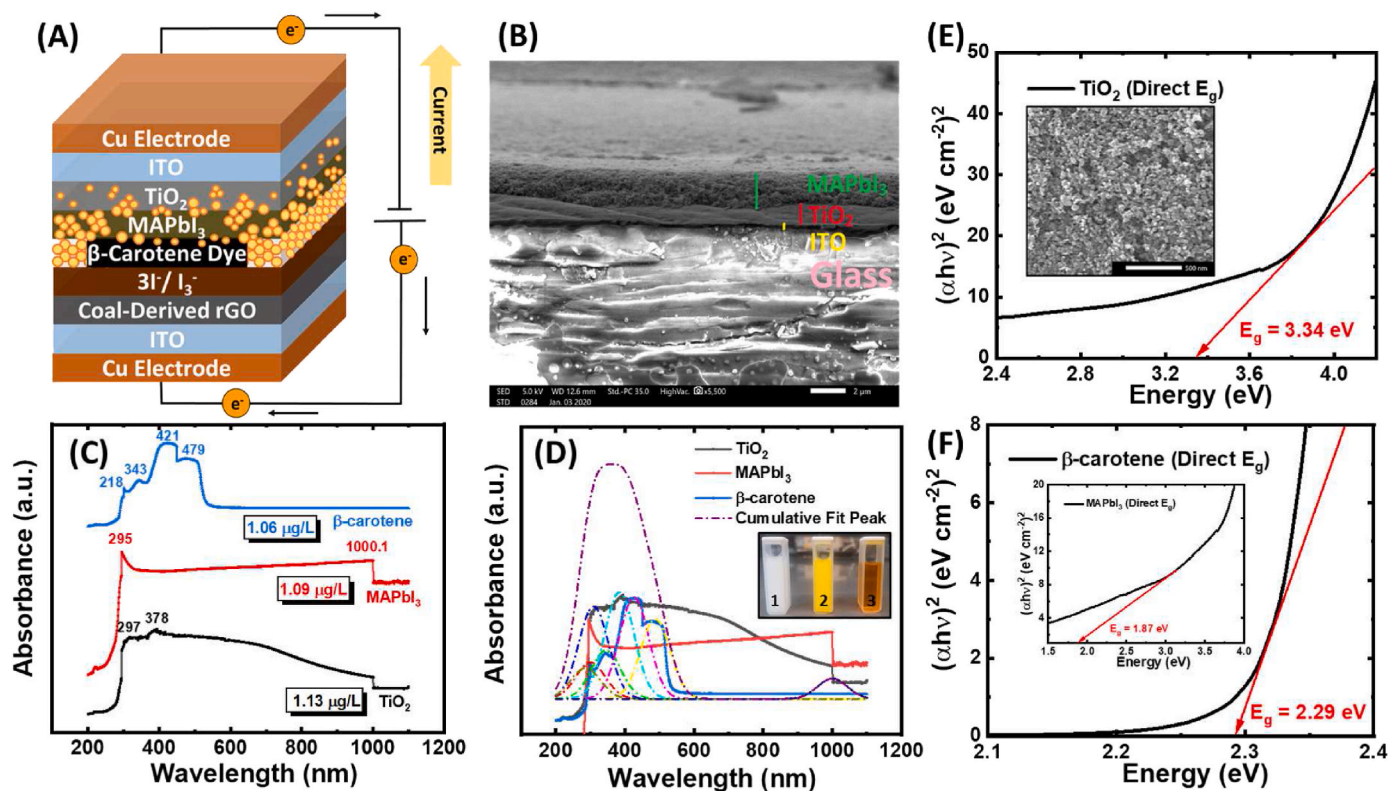


Fig. 4. β -carotene-sensitized and MAPbI₃-cosensitized bio-DSSCs: (A) Cartoon design of the DSSC with both β -carotene natural sensitizer and perovskite cosensitizer; (B) FESEM image showing sensitized/cosensitized spin-coated photoanode cross section and deposited layers with their approximated thicknesses; (C) Absorbance curves for TiO₂ and sensitizers in H₂O solutions showing absorbance maximum of each material; (D) Cumulative peaks Gaussian fitting of the sub-fitted peaks of TiO₂ and sensitizers showing 2-fold increase in the photoanode light absorption within 200–600 nm UV–Vis spectrum, **inset of (D)**: Solutions of (1) TiO₂, (2) MAPbI₃, and (3) β -carotene; (E) Tauc plot for TiO₂ semiconductor with direct bandgap = 3.34 eV, **inset of (E)**: surface morphology of ITO/TiO₂ showing TiO₂ nanoparticles with average-diameter of 30–40 nm; (F) Tauc plot and direct bandgap for β -carotene (2.29 eV), **inset of (F)**: Tauc plot and direct bandgap for MAPbI₃ (1.87 eV) sensitizers.

and the liquid electrolyte, which can easily damage the perovskite structure from the introduced wetting of the surface via iodide/triiodide electrolyte. From a comparison of our work with previous studies (2018–2020) [6,9,21,47,49,50] which utilized β -carotene in TiO₂-based DSSCs, none of the previous studies have been able to achieve $V_{oc} > 0.5$ V, but with higher photocurrent densities than those observed here in the designed perovskite-cosensitized bio-DSSCs. This can be attributed to the role of MAPbI₃ in reducing the voltage loss from the presumed upward shifting in TiO₂ Fermi level (E_F) due to the increased number of injected electrons via MAPbI₃ cosensitization, resulting from the decreased ϕ_{SBH} at ITO/TiO₂ junction blocking back-electron transport. However, the observed low current densities may be a result of poor electron collection abilities of the GO/rGO-based CE to capture the incoming forward electrons from the photoanode side. Thus, utilizing Ag might allow to increase the photocurrent density and achieve high PCEs.

The J-V curves of the architecturally designed DSSCs with different photoanodes are plotted in Fig. 3-C, with the obtained photovoltaic parameters and cell performance reported in Table 2. The involvement of both MAPbI₃ + β -carotene as a co-sensitizer and a sensitizer, respectively, resulted in increasing both J_{sc} and V_{oc} leading to a higher fill factor (FF) and PCE (η). The reason behind the enhanced voltage is attributed to the presumed upward shifting in TiO₂ Fermi level (E_F) due to the increased number of injected electrons via MAPbI₃ cosensitization (lower TiO₂ work function decreases ϕ_{SBH} at ITO/TiO₂ junction blocking back-electron transport for minimized voltage loss). Spin-coated photoanodes showed higher V_{oc} due to the reduced R_s and enhanced tunneling decreasing voltage loss based on the proposed transport mechanisms. However, the high currents in perovskites-sensitized cells are attributed to (i) low J_0 or less leakage current, (ii) enhanced exciton

generation (e-h pairs) by the dye/perovskite molecules, and (iii) facilitated electron injection or charge separation at [MAPbI₃+dye]/TiO₂. Cosensitization with MAPbI₃ resulted in higher J_{sc} owing to the low R_s allowing easy e^- injection from enhanced exciton generations and photon absorption. Reproducibility of the J-V curves has been also considered (see Supplementary, Figure S3). The FF values of non-sensitized photoanode-based DSSCs are found low because FF is known to be associated with both V_{oc} and I_{sc} as well as the squareness of the J-V curve. Such low produced voltages have resulted in low FF. Also, cosensitized bio-DSSCs showed relatively low FF~0.3–0.6 due to the high current leakage and poor current collection from the undesired interfacial recombination introduced by added interfaces. FF values have increased from 0.24 to 0.36 and from 0.32 to 0.63 when TiO₂ nanoparticles were interfaced with [MAPbI₃ + β -carotene] using dip-coated and spin-coated (Fig. 3-D) photoanodes, respectively, in the different designed cell architectures (Figure S4). This results into a maximum PCE of $(1.40\text{--}4.56) \times 10^{-4}\%$ (three-order-of-magnitudes higher than that of non-sensitized DSSCs, and 4-fold higher than similar β -carotene powder-based DSSCs [15]; see Table 3).

Series resistance (R_s) and shunt resistance (R_{sh}) were estimated from the linear relations obtained from the dark diode curves using Cheung's methods [68,69] equation: $\frac{dV}{d\ln(I)} = IR_s + \left(\frac{nk_B T}{q}\right)$, and the J-V curves as $R_s = -[\Delta V_{oc}/\Delta I_{oc}]$ and $R_{sh} = -[\Delta V_{sc}/\Delta I_{sc}]$, respectively, where $I = J$ here [70]. From the J-V curves, maximum power points were determined and taken as a reference to obtain the two linear lines at V_{oc} and J_{sc} locations. For efficient cells and high currents, R_s should be low for better electron injection and R_{sh} must be high enough to prevent current leakage. Both [MAPbI₃+ β -carotene] enhanced R_{sh} (but with increased

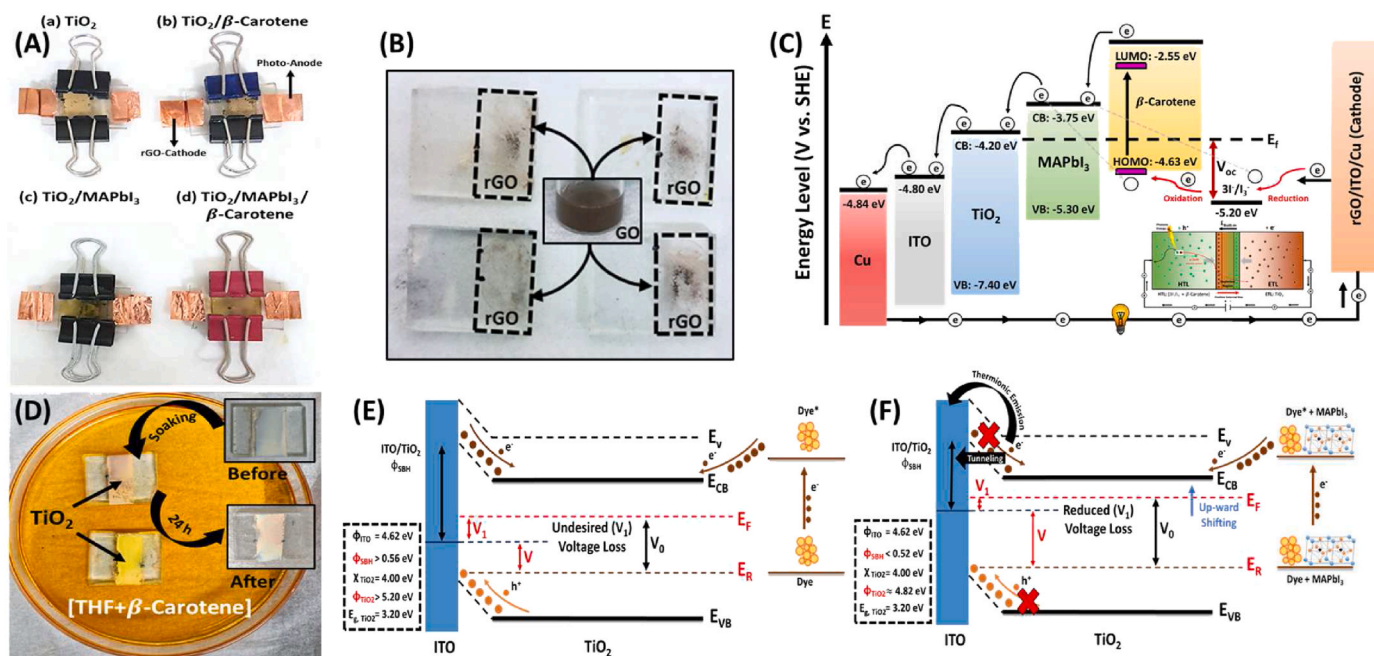


Fig. 5. (A) Fabricated bio-DSSCs using different photoanode designs with natural-sensitization and/or cosensitization; (B) Post-annealed ITO/rGO counter electrodes for lower R_s and efficient current (electron) collection; (C) Energy levels and band diagram alignment for the sensitized/cosensitized photoanode with perfect energy cascading for enhanced carriers mobility and lower interfacial recombination owing to the visible light absorption and low bandgaps of β -carotene and MAPbI₃, **inset of (C):** hypothesized HTL/ETL (or HTL/Perovskite/ETL) theory for exciton generation and charge transport across the built-in potential; (D) Crystallized TiO₂ dip-coated thin-film (2 cm \times 1.5 cm) on ITO substrate before and after β -carotene overnight deposition/soaking; (E) TiO₂ Fermi level and ITO/TiO₂ SBH before MAPbI₃ cosensitization; (F) Upward shifting of TiO₂ Fermi level after MAPbI₃ cosensitization minimizing ITO/TiO₂ SBH for reduced voltage loss (V_1) and prevented back-electron transport to TiO₂.

R_s) yielding in relatively high RF and J_{sc} (Fig. 3-F). V_{oc} has been improved from the noticeable decrease in Φ_{SBH} at ITO/TiO₂ interface (from 0.61 to 0.53 eV) from addition of both [MAPbI₃ + β -carotene] sensitizers which prevented back-electron transport, and increased excitons generation and electron injection rates (Table 1). The sensitized/cosensitized photoanode active layers and cell architecture including β -carotene are illustrated in Fig. 4-A (Supplementary, see Figure S4 for the other designs).

The optimally designed cell structures with perovskite cosensitization are illustrated in Fig. 4-A. FESEM and SEM images (Fig. 4-B, Fig. 4-E inset, Figures S5-S8, EDS Figures S9, S10) show the uniformly deposited/annealed TiO₂ morphology on ITO and photoanode cross-section layers characterization for quantifying each layer thickness: ITO \sim 260 nm, TiO₂ \sim 1.2 μ m, and MAPbI₃ \sim 2 μ m. Since the annealing temperature of TiO₂ films was at 400 $^{\circ}$ C, anatase crystals are expected to form (300–700 $^{\circ}$ C) [71] (observed as small and scattered crystals with typical \sim 40 nm grains), rather than rutile and brookite. Tauc plots for the WBG semiconductor and the sensitizers (Fig. 4-E and Fig. 4-F), obtained from UV–Vis absorbance curves (Fig. 4-C) and the linearized Tauc Equation $ah\nu = A(h\nu - E_g)^{1/2}$ [72], are used to fit the absorption spectrum for the direct allowed transition bandgaps of TiO₂, MAPbI₃, and β -carotene. Despite that MAPbI₃ is considered a direct bandgap semiconductor due to the high VB density of state, an indirect bandgap exists from spin-orbit coupling owing to CB Rashba-splitting [73]; where α is the absorption coefficient, $h\nu$ is the photon energy, A is a constant, and E_g is the bandgap; α is calculated from: $\alpha = (2.303t)^{-1} \log(1/OD)$ where t is the thickness of the quartz cuvette (light path), 2.303 is the conversion constant between $\log(x)$ and $\ln(x)$ and OD is the optical density (UV/Vis absorbance) [73].

First-principle DFT calculations showed that anatase behaves as an indirect transition bandgap semiconductor with high photocatalytic activity. Indirect transition ensures a longer photogenerated exciton lifetime. Anatase phase of TiO₂ is characterized with the lightest average

effective mass phase yielding faster electron migration from TiO₂ to ITO with minimum recombination rates within TiO₂ [74]. WBG TiO₂ was observed with a direct bandgap of 3.34 eV (Fig. 4-E) that is equivalent to the reported 3.3 eV for TiO₂ nanoparticles utilized by Kumar et al. [75] in naturally-sensitized DSSCs [Jambolana + Eosin dye] with 0.13% PCE, where absorbance lies in the UV region. The direct bandgap of MAPbI₃ was found to be 1.87 eV (inset in Fig. 4-F), which agrees with the Guo et al. [76] reported literature range (1.53–2.1 eV). An estimated direct bandgap of 2.29 eV (Fig. 4-F) was obtained for β -carotene; this is smaller than the maximum reported direct bandgap of 3.74 eV (1.59 eV for indirect transition) for β -carotene in the fabricated x (SnO₂)-y (Ag)- β -carotene/ITO photoanode [5,77]. However, other works estimated the HOMO/LUMO energy difference for various classes of carotene to be 2.2–6.7 eV [78]. Tunability of β -carotene's energy bandgap is possible via incorporation of Ag/SnO₂ nanoparticles yielding β -carotene nanocomposite with lower E_g for its application in DSSCs, giving better molecular-dye electron excitation [77]. Regarding tunability of bandgap energy in TiO₂, increasing the semiconductor annealing temperature can blue-shift energies (decreasing bandgap); while, a decrease in semiconductor grain size decreases E_g due to electron quantum confinement [79].

Absorbance curves from Fig. 4-C showed maximum absorption (λ_{max}) for TiO₂ at 297, 378 cm^{-1} , MAPbI₃ at 295, 1000.1 cm^{-1} , and β -carotene at 218, 343, 421, 479 cm^{-1} confirming that TiO₂ absorbs light mostly from the UV-region whereas sensitization fosters visible-light absorption. Cumulative peaks fitting (Gaussian fitting in ORIGIN, Fig. 4-D) from the sub-fitted peaks of TiO₂ and both β -carotene and MAPbI₃ sensitizers confirmed a minimum of 2-fold increase in the photoanode light absorption capabilities in the 200–600 nm UV–Vis spectrum.

Maximum theoretical PCE of <0.1% (spin-coated) and <0.7% (dip-coated) are estimated from the maximum theoretical $V_{oc} = 1$ V plugged into $V_{oc} = \frac{nk_B T}{q} \ln \left(\frac{J_{ph}}{J_0} \right)$ and $[\eta^* = \frac{P_m}{P_{in}}]$ according to DSSCs circuit and TE models (Supplementary, see section 1) with an assumption that there

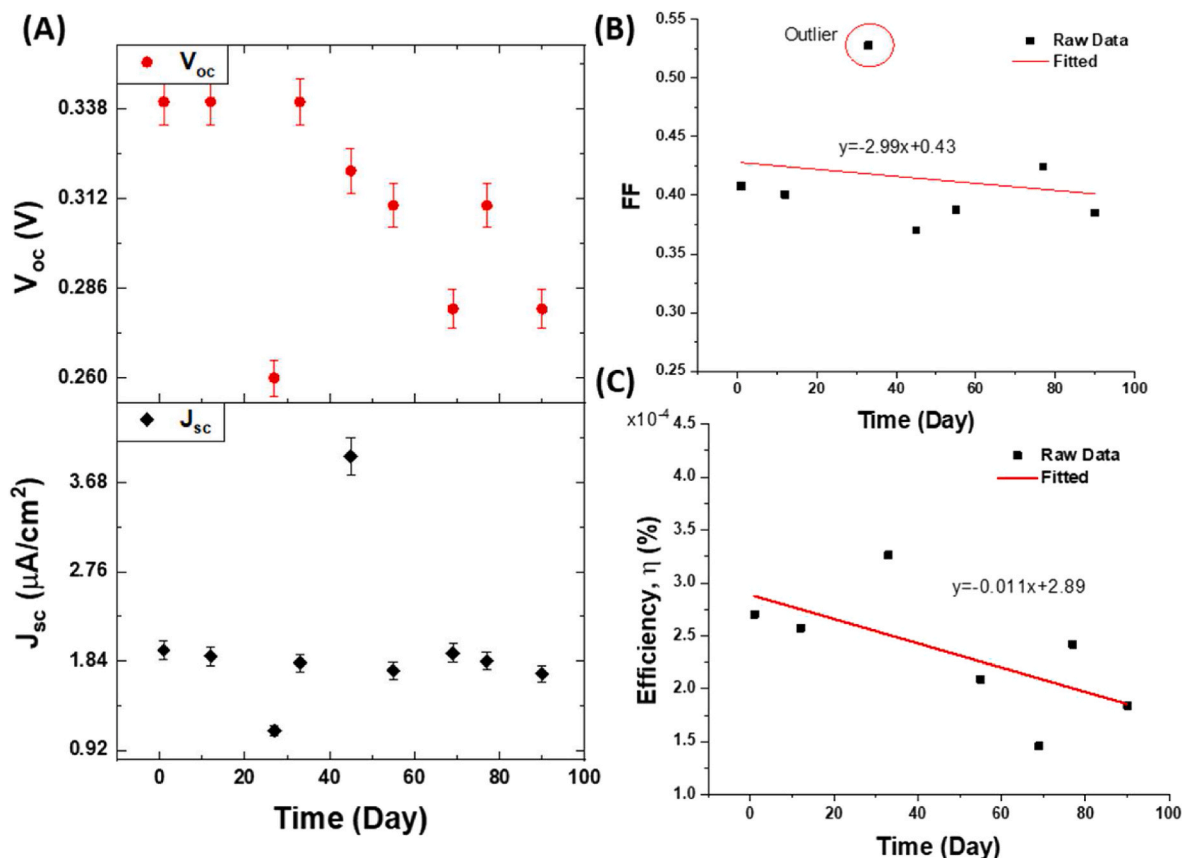


Fig. 6. Cell stability over time for naturally-sensitized bio-DSSCs using β -carotene dip-coated photoanodes: (A) Changes in the observed V_{oc} (2% error bars) and J_{sc} (5% error bars); (B) Slight reductions in FF with the estimated linear fitted relationship considering an outlier value; (C) Slight reductions in efficiency with the estimated linear fitted relationship after removing any outlier values.

will be similar observed current densities; where symbols definitions can be found in the [Supplementary, see section 1](#). If current densities rise $\times 100$ from using a silver (Ag) counter electrode, the maximum PCE can then theoretically go up to $\approx 10\%$.

The β -carotene dye deposited on TiO_2 -photoanodes and their use in the different designed/fabricated DSSCs are shown in [Fig. 5-D](#) and [Fig. 5-A](#), respectively. Post-annealed and prepared ITO/rGO CE for lower R_s are shown in [Fig. 5-B](#). The uniformity of the rGO film determines the capability of the CE for current collection; however, the random deposition and non-uniform coated rGO negatively influenced the device performance by sharply decreasing J_{sc} to $\sim \mu A/cm^2$. Such uniform coating of graphite-based materials is mandatory for easy electron collection and enhanced current density. Energy band diagram alignment in [Fig. 5-C](#) illustrates perfect energy cascading for enhanced photoanode electron mobility, reduced interfacial recombination, and continuous forward electron pathways. The low bandgaps of both β -carotene (2.29 eV) and $MAPbI_3$ (1.87 eV) assure visible-light absorption, with their LUMO (-2.55 eV) and CB (-3.75 eV) relatively in higher states than TiO_2 CB (-4.2 eV) for more exciton generation and easy electron injection with a maximum ideal $V_{oc} = 1$ V. The shift in the bandgap of $MAPbI_3$ crystals from known bandgap around 1.6 eV–1.87 eV might be attributed to the decomposition of perovskites when interfaced with the liquid redox electrolyte $[I^-/I_3^-]$. Low electron excitation energy requirements in the sensitizers ensure easy electron excitation and injection above the Fermi level of TiO_2 . The large difference between HOMO/redox potentials drives electron-diffusion which effectively controls the charge injection and dye reduction rates [80].

In DSSCs, a Schottky diode exists at the metal-semiconductor interface in the photoanode structure as in ITO/ TiO_2 or $[Cu/ITO]/TiO_2$ junctions. Onodera et al. showed that V_{oc} decreased with increasing ϕ_{SBH}

[40–42]. When TiO_2 comes in contact with ITO, band bending in TiO_2 CB occurs influencing the current-loss mechanism from inhibited electron collection or increased recombination in agreement with the thermionic-emission theory [42]. We suspect that there are three mechanisms responsible for the improved PCE from $MAPbI_3$ cosensitization: (i) absorbance broadening and increase in the probability of absorbing more visible-photons (400–700 nm) [35], (ii) upward shifting of TiO_2 Fermi level due to increased injected electrons and CB effective density of states $> 2.46 \times 10^{21} cm^{-3}$ reducing the TiO_2 work function to 4.82 eV (Supplementary, see section 9), consequently preventing back-electron transport allowing only quantum tunneling of photoelectrons at ITO/ TiO_2 , rather than transport over a Schottky barrier [81, 82] ([Fig. 5-E,F](#)), and (iii) hindered ϕ_{SBH} at the ITO/ TiO_2 interface with a 13.1% reduction in photoinduced electrostatic potential barrier due to the minimized band bending of TiO_2 CB and reduced TiO_2 work function resulting in low built-in potential and lower voltage loss (V_1). All of the described mechanisms are attributed to possible ITO/ TiO_2 electron tunneling from the increased number of initially injected electrons from perovskite and dye molecules into TiO_2 CB [42,83]. Briefly, this upward shifting in E_F level for TiO_2 promotes both electron tunneling and V_{oc} while decreasing the minimum TiO_2 work function (ϕ_{TiO_2}) required for electron excitation, injected electrons tunnel from TiO_2 to ITO overcoming the suppressed ϕ_{SBH} at the presumed junction with hindered TiO_2 CB band bending resulting in less recombination and high V_{oc} [84–86] from the $MAPbI_3$ cosensitization.

The stability of the designed devices and photoanodes have been studied over time (3-month period) as shown in [Fig. 6](#). This would allow us to study the durability of utilizing perovskite semiconductors to cosensitize naturally-sensitized bio-DSSCs. The long-term stability of the devices is a critical indicator to assess the practicability and operability

Table 3

Performance comparison and photovoltaic parameters of various literature works on β -carotene-based DSSCs in contrast with the proposed bio-DSSCs system under AM1.5G radiation.*

Pigment: Source/Absorbance [Architecture]	J_{sc} (μA cm^{-2})	V_{oc} (mV)	FF	η^* (%)	Ref.
Daucus Carota/415–508 nm [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/C/FTO]	33	239	–	1.2×10^{-3}	[39]
Carrot/400–510 nm; peaks: 433, 472.5 nm [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/Pt/FTO]	360	40	0.64	9×10^{-3}	[21]
Orange fruit; 400–510 nm; peaks: 441, 466.5 nm [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/Pt/FTO]	370	60	0.58	2×10^{-2}	[21]
Tomatoes/400–510 nm; max: 441 nm and 466.5 nm [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/Pt/FTO]	510	140	0.37	3×10^{-2}	[21]
Syngonium sp leaf/440–480 nm; peaks: 450, 480 nm; bandgap = 2.47 eV; EL-HSE (High stability redox, Dyesol) [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/Pt/FTO]	270	400	0.68	7.4×10^{-2}	[9]
Yellow sweet potato/400–500 nm; peaks: 401 nm [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/Pt/FTO]	0.27	515	–	5.7×10^{-2}	[6]
Purple sweet potato/400–500 nm; peaks: 532 nm, [FTO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/Pt/FTO]	0.17	378	–	3.3×10^{-2}	[6]
β-carotene powder; 380–500 nm [ITO/TiO ₂ /β-carotene/[I ⁻ /I ₃ ⁻]/C/ITO]	162	123	0.039	1×10^{-4}	[50]
β-carotene powder; 218–479 nm [ITO/TiO ₂ /MAPbI ₃ /β-carotene/[I ⁻ /I ₃ ⁻]/rGO/ITO]	1.10	655	0.63	4.6×10^{-4}	This work

* J_{sc} = Short-circuit current density; V_{oc} = Open-circuit voltage; FF = Fill factor; η = Incident photon-to-current efficiency (IPCE) = Quantum efficiency (QE). Note that we have used the β -carotene powder source similar to the last reported study [50], Sigma US, and recorded better η .

of such naturally-sensitized devices simply assembled under ambient conditions. Both FF and J_{sc} values found to be almost the same over 3 months, as illustrated in Fig. 6-A and Fig. 6-B. This indicates the capability of the deposited light-absorbing layers in maintaining forward electron transport and charge separation. It is worth mentioning that each of the nine measurements was carried under a light incident of 100 mW/cm² with three values averaged over a 2 h radiation period with surrounding humidity \sim 35%, and room temperature of \sim 23 °C. The stable current generation (even at low values) suggests a stable photoanode layer and stable electron collection at the cathode side. However, the overall PCE decrease rate was found to be 10%/month, which is due to the poor stability of MAPbI₃ for maintaining high V_{oc} based on the suggested transport mechanisms and the hindered loss of voltage from SBH, as obtained in Fig. 6-A and Fig. 6-C. In other words, the observed low voltage is linked to possible degradation of the MAPbI₃ layer from excess oxygen and light stresses. Also, MAPbI₃ crystals are more susceptible to degrade and lose their octahedral stable structures when coming in contact with a liquid-state electrolyte, resulting in decreasing both V_{oc} and device efficiency over time.

Apparently, employing environmental-friendly sensitizers including biomolecular dyes (carotenoids) with perovskite semiconductors for hybrid photoanode architectures will facilitate exciton generation and electron forward transport owing to the highly conjugated dye double bonds structure (π -electrons in their molecular structure) [12,13,87] and broadened UV–Vis absorption. The designed solar cells are of a

conventional type where electrons move to n-type (ETL) and hole to p-type (HTL) layers ensuring fast charge separation. Both MAPbI₃ and β -carotene absorb incident photons which initiate electron excitation before the charge separation step. Further photoelectrochemical characterizations such as EIS, IPCE, IMVS, and IMPS are important to essentially understand the proposed mechanisms and the reasons behind the low current and/or reduced cell efficiency. Future works and further steps may include the following: (i) utilization of acidic dye solutions and/or surface hydroxylation of TiO₂ semiconductor (to introduce more –OH anchoring groups) for strong dye molecules acceptor-moieties attachment yielding in lower series resistance, (ii) designing of 3D graphene-based counter electrodes with interconnected networks to provide fast electron transport from the cathode, diminished contact resistance at the graphene/electrolyte interface [88], and high electrocatalytic activity towards reduction of I₃⁻, and (iii) incorporation of low-dimensional WBG 1D-ZnO nanowires ensuring large surface area for better dye coverage, high V_{oc} , and enhanced photoanode stability against excessive UV radiation. Table 3 shows a performance comparison of β -carotene-based DSSCs with previous recent works (2018–2020) [6,9,21,47,49,50], as have been earlier discussed in the second section of this work.

5. Conclusion

We demonstrated natural photosensitization using bio-sensitized TiO₂ photoanodes which are further cosensitized by MAPbI₃ perovskite semiconductors for enhanced photon absorption, exciton generation, charge separation, and electron injection resulting in reduced voltage loss. The hybrid biomolecular dye-perovskite photoanode heterostructure [TiO₂/MAPbI₃/β-carotene] can broaden the optical absorption spectrum and reduce interfacial recombination, and most importantly facilitate interfacial charge separation, electron injection, and electron tunneling which reduce voltage loss at the presumed ITO/TiO₂ heterojunction for enhanced V_{oc} . The observed results were attributed to the perovskite engineered cascaded molecular energy levels in the hybrid biophotonode, which enabled upward shifting in Fermi level (E_F) for TiO₂, reduced work function in TiO₂, that would hinder TiO₂ CB bending. Perovskite semiconductor-integrated biophotoanode structures found to enhance the diode ideality factor, reduce the leakage current, improve the cell rectification factors, increase the FF, and noticeably show high V_{oc} (>0.65 V), from shifting of E_F level and lowering of ϕ_{SBH} .

The use of β -carotene as a light-absorbing material instead of the Ru-based complexes is already known to result in DSSCs with low PCEs, which can be further improved via perovskite semiconductor cosensitization. The thickness of the β -carotene and the number of absorbed photons in the dye layer needs to be further investigated for understanding the charge transport mechanisms and quantifying carrier lifetime. The very low device efficiency (in the order of 10^{-2} – $10^{-4}\%$) might be attributed to the poor devices stability and degradation of MAPbI₃ layers when interfaced with the liquid electrolyte [I⁻/I₃⁻]. The designed solar cells are of a conventional type where electrons transport to n-type (ETL) and holes to p-type (HTL) layers ensuring fast charge separation. Both MAPbI₃ and β -carotene absorb incident photons which initiate electron excitation before the charge separation step. Further photoelectrochemical characterizations such as EIS, IPCE, IMVS, and IMPS are important to essentially understand the proposed mechanisms and the reasons behind the low photocurrent and/or cell efficiency. The low photocurrent densities could be due to the unfavorable high series resistances that need further investigation on the engineered interfacial contacts with the possibility of using Ag and/or noble-Pt-like CE for better electron collection. Nevertheless, cosensitization prevents back-electron transport which also facilitates exciton generation, electron injection, and the possibility of a 2-fold increase in photoanode UV–Vis absorption capabilities. This work paves the way towards designing efficient bio/co-sensitized photoanodes, which can be further improved

from the incorporation of WBG low-dimensional materials, solid-state HTL redox electrolytes, and Ag/Au plasmonic nanoparticles leveraging photon absorption, electron injection and hopping, dye regeneration, and interfacial charge transport for improved overall device performance.

Credit author statement

Hisham A. Maddah: Conceptualization; Data curation; Formal analysis; Investigation; Methodology; Resources; Visualization; Roles/Writing - original draft; Writing - review & editing. **Lila Aryadwita:** Formal analysis; Investigation; Writing - review & editing. **Vikas Berry:** Resources; Supervision; Writing - review & editing. **Sanjay K. Behura:** Conceptualization; Investigation; Data curation; Formal analysis; Methodology; Project administration; Resources; Supervision; Visualization; Roles/Writing - original draft; Writing - review & editing

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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Appendix A. Supplementary data

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