


Lecture 4 · Existence - Uniqueness Theorem

Last Time: Slope Fields

— behavior of solution curves

This time: When can we guarantee that a 1st order initial value problem has a solution? Only one solution

Theorem: (Existence + Uniqueness)

Consider $\begin{cases} \frac{dy}{dx} = f(y, t) \\ y(t_0) = y_0 \end{cases}$

- i) If $f(t, y)$ is continuous in neighborhood (t_0, y_0) then a solution exists in some rectangle $|t - t_0| < \delta$, $|y - y_0| < \epsilon$
- ii) If, in addition, $\frac{\partial f}{\partial y}(y, t)$ is continuous in a neighborhood of the point (t_0, y_0) then the solution is unique

Example: Recall that $\begin{cases} \frac{dy}{dx} = y^{1/3} \\ y(0) = 0 \end{cases}$ has two solutions $\Rightarrow y_1(t) = 0$ + $y_2(t) = (\frac{2}{3}t)^{3/2}$

Notice that we have $f(y, t) = y^{1/3}$ is continuous in a neighborhood of the point $(0, 0)$ but $\frac{\partial f}{\partial y} = \frac{1}{3}y^{-2/3}$ is not continuous in a neighborhood of $(0, 0)$

Example: Consider $\begin{cases} \frac{\partial y}{\partial t} = (y-5)\ln(1-y-5) + t \\ y(0) = 5 \end{cases}$

Here $f(y, t) = (y-5)\ln(1-y-5) + t$ is continuous in a neighborhood of $(0, 5)$

- so the theorem guarantees the existence of a solution for $t \approx 0$

However $\frac{\partial f}{\partial y} = 1 + \ln(1-y-5)$ is not continuous in a neighborhood of $(0, 5)$ so the theorem does not guarantee uniqueness

Notice that if the initial condition were different (e.g. $y(0) = 1$) then both parts of the theorem would apply.

• For some easy equations, we can carry out the iterations explicitly:

Example: Consider $\frac{\partial y}{\partial t} = 2t(1+y)$, $y(0) = 0$

Here $y_0(t) = 0$, $y_1(t) = 0 + \int_0^t 2s(1+0)ds = t^2$

$$y_2(t) = 0 + \int_0^t 2s(1+s^2)ds = t^2 + \frac{1}{2}t^4$$

$$y_3(t) = 0 + \int_0^t 2s(1+s^2+\frac{1}{2}s^4)ds = t^2 + \frac{1}{2}t^4 + \frac{1}{6}t^6$$

$$\dots y_n(t) = 0 + \int_0^t 2s(1+y_{n-1}(s))ds = t^2 + \frac{1}{2}t^4 + \dots + \frac{1}{n!}t^{2n}$$

Recall that $e^t = 1 + t + \frac{1}{2}t^2 + \dots + \frac{t^n}{n!} + \dots$

We recognize that $y_n(t) \rightarrow e^{t^2} - 1$

Notice that if $y(t) = e^{t^2} - 1$, then $\frac{\partial y}{\partial t} = 2te^{t^2} = 2t(y+1)$

so this is indeed a (the) solution

Example: Let's consider $\begin{cases} \frac{dy}{dt} = t^2 + y^2 \\ y(0) = 1 \end{cases}$

Here $y_0(t) = 1$ and for $n > 0$, $y_n(t) = 1 + \int_0^t s^2 + (y_{n-1}(s))^2 ds$

$$\text{so } y_1(t) = 1 + t + \frac{1}{3}t^3$$

$$y_2(t) = 1 + t + t^2 + \frac{2}{3}t^3 + \frac{1}{6}t^4 + \frac{2}{15}t^5 + \frac{1}{63}t^7$$

$$y_3(t) = 1 + t + t^2 + \frac{4}{3}t^3 + \frac{5}{6}t^4 + \frac{6}{15}t^5 + \frac{29}{60}t^6 + \frac{47}{315}t^7 + \dots + \frac{1}{59535}t^{15}$$

The pattern for the coefficients is not clear, so it would be hard to guess a solution

How does one show the sequence converges (in general)?

Consider the map T : functions $\rightarrow$ functions

that takes a function y to the function $T(y)(t) = y_0 + \int_0^t f(y(s), s) ds$

The Picard sequence is $y_0, y_1 = T(y_0), y_2 = T(y_1) = T(T(y_0)) = T^2(y_0)$
 $y_3 = T(y_2) = T^3(y_0) \dots$
 $y_n = T^n(y_0)$

What we show is that there is some appropriate notion of size of a function

$\|f\|$ and that T reduces distances

$$\|T(f) - T(g)\| \leq \rho \|f - g\|, \rho < 1$$

T is called a contraction

This is enough to show that $y_n = T^n(y_0)$ converges

If we call the limit y_∞ , then notice that

$$y_\infty = \lim_{n \rightarrow \infty} T^n(y_0) = \lim_{n \rightarrow \infty} (T(T^{n-1}(y_0))) = T(y_\infty)$$

so y_∞ is a "fixed point" of T but then

$$y_{\alpha}(t) = T(y_{\alpha})(t) = y_0 + \int_{t_0}^t f(y_{\alpha}(s), s) ds$$

This is just another way of writing

$$\begin{cases} \frac{dy}{dt} = f(y_{\alpha}, t) \\ y_{\alpha}(t_0) = y_0 \end{cases}$$

Historical Context: The Cycloid



$$x(\theta) = a(\theta - \sin \theta)$$

$$y(\theta) = a(1 - \cos \theta)$$

Substitutions

$$\frac{dy}{dt} = F\left(\frac{y}{t}\right) \quad y' = \sin\left(\frac{y}{t}\right) + \left(\frac{t}{y}\right)^2 + 3 + e^{\frac{y}{t}}$$

use sub: $v(t) = \frac{y}{t}$

ex:

$$\frac{dy}{dt} = \frac{t}{y} + \frac{y}{t} \quad (y \neq 0)$$

use $v = \frac{y}{t} \Rightarrow y = vt$ $\rightarrow \frac{dy}{dt} = \frac{d}{dt}(vt) = v't + v = F(v)$

substitute

$$v't + v = \frac{1}{v} + v$$

$$t\left(\frac{dv}{dt}\right) = \frac{1}{v} \quad (\text{separable equation})$$

$$\int v dv = \int \frac{1}{t} dt \quad F(v)$$

$$\frac{v^2}{2} = \ln(t) + C$$

Exact Equations

$$\Psi(x, y) = C$$

$$d\Psi = \frac{\partial \Psi}{\partial x} dx + \frac{\partial \Psi}{\partial y} dy = 0$$

$$\frac{\partial \Psi}{\partial y} \frac{dy}{dx} + \frac{\partial \Psi}{\partial x} = 0$$

$$y = y(x)$$



$$N(x, y) \frac{dy}{dx} + M(x, y) = 0$$

$$M dx + N dy = 0$$

What if $\exists$ A Ψ

$$M = \frac{\partial \Psi}{\partial x} \quad \text{and} \quad N = \frac{\partial \Psi}{\partial y}$$

$$\text{if } M = \frac{\partial \Psi}{\partial x} \quad \text{and} \quad N = \frac{\partial \Psi}{\partial y}$$

$$\Rightarrow \frac{\partial M}{\partial y} = \frac{\partial^2 \Psi}{\partial y \partial x} \quad \text{and} \quad \frac{\partial N}{\partial x} = \frac{\partial^2 \Psi}{\partial x \partial y}$$

$$\Leftrightarrow \frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

Δ the solution to the eq is

$$\Psi(x, y) = C$$

Example

$$(x^2 + 3y^2) \frac{\partial y}{\partial x} + 2xy = 0$$

$$\underbrace{2xy \partial x}_M + \underbrace{(x^2 + 3y^2) \partial y}_N = 0$$

to check if it is an exact equation

$$\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$$

$$\frac{\partial M}{\partial y} = 2x \quad \frac{\partial N}{\partial x} = 2x \quad \Rightarrow \quad \frac{\partial \psi}{\partial x} = M = 2xy$$



$$\int \frac{\partial \psi}{\partial x} dx = \int 2xy dx$$

$$\frac{\partial \psi}{\partial y} = N = x^2 + 3y^2$$

$$\psi(x, y) = x^2 y + f(y) \rightarrow \frac{\partial \psi}{\partial y} = \cancel{x^2} + f'(y) = \cancel{x^2} + 3y^2$$

$$f = y^3 + A \Rightarrow \psi(x, y) = x^2 y + y^3 + A$$

$$\text{since } \psi(x, y) = C \Rightarrow$$

$$\psi(x, y) = x^2 y + y^3 = C$$

Another example

$$2x \frac{\partial y}{\partial x} + y^2 = 0 \Rightarrow \underbrace{y^2}_{M} \partial x + \underbrace{2x}_{N} \partial y = 0$$

$$\frac{\partial M}{\partial y} = 2y$$

$$\frac{\partial N}{\partial x} = 2 \quad \neq \quad \text{not exact solution}$$

Example 3

$$\frac{\partial y}{\partial x} = \frac{2xy^2}{\sin(y) - 2x^2y} \Rightarrow (\sin(y) - 2x^2y) \partial y = 2xy^2 \partial x$$
$$\underbrace{2xy^2 \partial x}_M + \underbrace{(\sin(y) - 2x^2y) \partial y}_N = 0$$

$$\frac{\partial M}{\partial y} = 4xy$$

$$\frac{\partial N}{\partial x} = 4xy$$

cont.

Lecture 7

Bernoulli Equation

$$y' + P(t)y = Q(t)y^\alpha, \quad \alpha \neq 0, 1$$

$$\text{use } v(t) = y^{1-\alpha}$$

Example:

$$y' = \frac{y^2}{t^2} + ty$$

$$y' - ty = \frac{1}{t^2} y^2$$

$$\text{use } v = y^{1-\alpha} = y^{-1} = \frac{1}{y}$$

$$y = \frac{1}{v}$$

$$y' = -\frac{1}{v^2} v'$$

$$-\frac{1}{v^2} v' - t \frac{1}{v} = \frac{1}{t^2} \cdot \frac{1}{v^2}$$

$$v' + tv = -\frac{1}{t^2}$$

1st order equation for v

$$\rightarrow \mu(t) = e^{\frac{t^2}{2}}$$

$$(v e^{\frac{t^2}{2}})' = -e^{\frac{t^2}{2}} \left(\frac{1}{t^2} \right) = -\frac{e^{\frac{t^2}{2}}}{t^2}$$

Suppose $F(\cancel{t}, y, y', y'') = 0$ but the equation does not explicitly have y

Use $v(t) = y'$

$$y'' = v'$$

$$F(t, v, v')$$

reducible
equation
2nd $\rightarrow$ 1st order

Example: $y'' + \frac{2}{t}y' = t^2$

$$v(t) = y'$$

$$v' = y''$$

$$\begin{aligned} v' + \frac{2}{t}v &= t^2 \\ \rightarrow \mu(t) &= e^{\int \frac{2}{t} dt} = e^{2 \ln(t)} = t^2 \end{aligned}$$

$$t^2 v' + 2t v = t^4$$

$$(vt^2)' = t^4$$

$$vt^2 = \frac{t^5}{5} + A$$

$$y' = v = \frac{t^3}{5} + \frac{A}{t^2}$$

$$\int y' = \int \left(\frac{t^3}{5} + \frac{A}{t^2} \right) dt$$

$$\boxed{y = \frac{t^4}{20} - \frac{A}{t} + B}$$

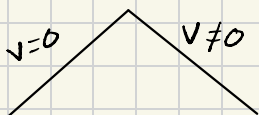
$F(\cancel{t}, y, y', y'') = 0$, use $v(y) = y'$

$$y'' = \frac{d}{dt} y' = \frac{d}{dt} v(y(t)) = \frac{dv}{dy} \cdot \frac{dy}{dt} = \frac{dv}{dy} \cdot v$$

Example: $y'' + 2yy' = 0$, Use $v(y) = y'$

$$y'' = v \frac{dv}{dy}$$

$$v \frac{dv}{dy} + 2yv = 0$$



$$y' = 0$$

↓

$$y = A$$

$$\frac{dv}{dy} + 2y = 0 \rightarrow \frac{dv}{dy} = -2y \rightarrow v = A - y^2$$

$$\frac{dy}{dt} = A - y^2 \rightarrow \frac{dy}{A - y^2} = dt$$

Slope Fields

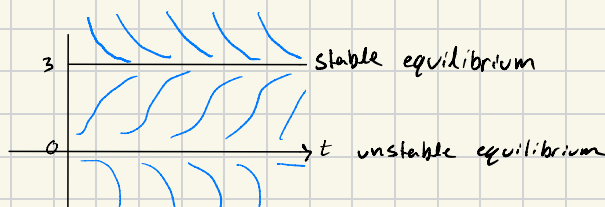
Autonomous EQs

Population Models

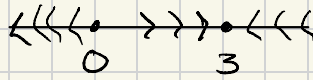
Ex: $P(t)$

$$\frac{dP}{dt} = 3P - P^2 = P(3 - P)$$

$P=0$ } critical points = Equilibrium
 $P=3$ } solutions



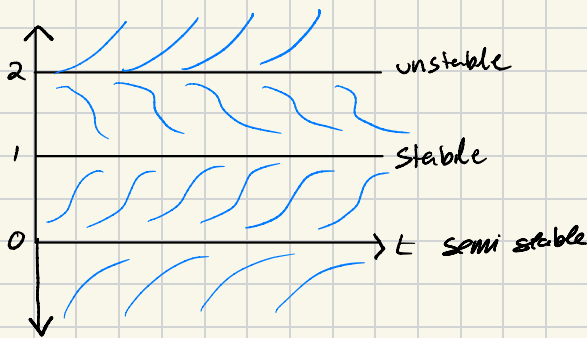
To visualize equilibrium
phase line



Ex: $\frac{dP}{dt} = P^2(P-1)(P-2)$

$P=0$
 $P=1$
 $P=2$

} Equilibrium states



Higher Order Linear ODEs

N^{th} Order Lin:

$$y^n + a_{n-1}(t)y^{n-1} + \dots + a_1(t)y' + a_0y = f(t)$$

$$y(t_0) = y_0$$

$$y'(t_0) = y_1$$

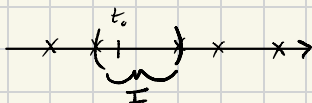
$\vdots$

$$y^{n-1}(t_0) = y_{n-1}$$

Theorem of Existence and Uniqueness for a linear IVP

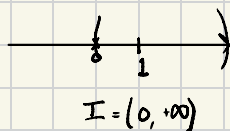
If $\exists$ an interval I around t_0 where
 $a_{n-1}, a_{n-2}, \dots, a_0, f$ are continuous

$\Rightarrow \exists$ one and only one solution to the IVP
over all of I



Example:

$$\begin{cases} y'' + \frac{1}{t}y' - \frac{9}{t^2}y = 0 \\ y(1) = 0 \\ y'(1) = 1 \end{cases}$$



Principle of Superposition For Linear Homogeneous Equations

If $y_1, y_2, \dots, y_n$ are solutions to a lin hom Eq

$\Rightarrow$ any linear combination $y = C_1 y_1 + C_2 y_2 + \dots + C_n y_n$ is also a solution to the equation

Example:

$$y'' - 3y' + 2y = 0$$

$$y_1 = e^t$$

$$y_2 = e^{2t}$$

and $y = C_1 e^t + C_2 e^{2t}$ is also a solution

Ex:

$$y'' - y = 4 \quad \leftarrow \text{linear non homogeneous}$$

$$y_1 = 4$$

$$y_2 = e^t - 4$$

} are solutions
But

$y = y_1 + y_2$ is not

Ex:

$$y'' + 2yy' = 0 \quad \leftarrow \text{not linear}$$

$$y_1 = 1$$

$$y_2 = \frac{1}{t}$$

} Are solutions
But $y = 1 + \frac{1}{t}$ is not

Ex:

Equidimensional

$$\begin{cases} y'' + \frac{1}{t} y' - \frac{9}{t^2} = 0 \\ y(1) = 0 \\ y'(1) = 1 \end{cases}$$

Try $y = t^a$

plug in

$$y' = a t^{a-1}$$

$$y'' = a(a-1) t^{a-2}$$

$$a(a-1) t^{a-2} + \frac{1}{t} a t^{a-1} - \frac{9}{t^2} t^a = 0$$

$$a(a-1) + a - 9 = 0$$

$$a^2 - 9 = 0$$

$$a = \pm 3$$

$$y_1 = t^3$$

$$y_2 = \frac{1}{t^3}$$

$$y = C_1 t^3 + \frac{C_2}{t^3}$$

$$y' = 3C_1 t^2 - \frac{3C_2}{t^4}$$

Apply initial conditions

$$y(1) = C_1 + C_2 = 1$$

$$y'(1) = 3C_1 - 3C_2 = 1$$

$$C_1 = \frac{1}{6}, C_2 = -\frac{1}{6}$$

$$y = \frac{1}{6}t^3 - \frac{1}{6} \cdot \frac{1}{t^3}$$

the solution to the IVP

To solve an n^{th} order lin-homogeneous IVP

1) Find N solutions to the equation: $y_1, y_2, \dots, y_n$

2) Combine linearly: $y = C_1 y_1 + C_2 y_2 + \dots + C_n y_n$

3) Apply ICs to find the C 's

Ex:

$$y'' + 4y = 0$$

$$y_1 = \sin(2t)$$

$$y_2 = \sin(t)\cos(t)$$

$$y = C_1 \sin(2t) + C_2 \sin(t)\cos(t)$$

$$y(0) = 0 + 0 \neq 1 \quad \leftarrow$$

does not work

because $\sin(2t) = 2\sin(t)\cos(t)$

IC's: $y(0) = 1$
 $y'(0) =$

Remember:

$$\cos(2t) = 2\cos^2(t) - 1 = \cos^2(t) - \sin^2(t) = 1 - \sin^2(t)$$

Linear Independence

Exam Review: Everything up until this past Friday

Find n solutions to the equation: $y_1, y_2, \dots, y_n(t)$

Write $\alpha_1 y_1 + \alpha_2 y_2 + \dots + \alpha_n y_n(t) = 0$

- If $\alpha_1 = \alpha_2 = \dots = \alpha_n = 0$ is the only choice $\rightarrow y_1, \dots, y_n$ are linearly independent
- If $\exists$ other choices of $\alpha \rightarrow y_1, \dots, y_n$ are linearly dependent

Example:

$$1, e^t, e^{-t}$$

$$\alpha_1 1 + \alpha_2 e^t + \alpha_3 e^{-t} = 0$$

$$\alpha_1 = \alpha_2 = \alpha_3 = 0$$

linearly independent

$$\cos(t), e^{it}, e^{-it}$$

$$\alpha_1 \cos(t) + \alpha_2 e^{it} + \alpha_3 e^{-it} = 0$$

$$\alpha_1 = 1 \quad \alpha_2 = \frac{-1}{2} \quad \alpha_3 = \frac{-1}{2}$$

linearly dependent

Note: $\cos(t) = \frac{e^{it} + e^{-it}}{2}$

$$\sin(t) = \frac{e^{it} - e^{-it}}{2i}$$

$$e^{it} = \cos(t) + i\sin(t)$$

$$e^{-it} = \cos(t) - i\sin(t)$$

$$\alpha_1 y_1 + \alpha_2 y_2 + \dots + \alpha_n y_n(t) = 0$$

$\downarrow$

$$\alpha_1 y_1' + \alpha_2 y_2' + \dots + \alpha_n y_n'(t) = 0$$

$\vdots$

$\downarrow$

$$\alpha_1 y_1^{(n-1)} + \alpha_2 y_2^{(n-1)} + \dots + \alpha_n y_n^{(n-1)}(t) = 0$$

linearly independent: $\begin{pmatrix} y_1 & y_2 & \dots & y_n \\ y_1' & y_2' & \dots & y_n' \\ \vdots & \vdots & \ddots & \vdots \\ y_1^{(n-1)} & y_2^{(n-1)} & \dots & y_n^{(n-1)} \end{pmatrix} \begin{pmatrix} \alpha_1 \\ \alpha_2 \\ \vdots \\ \alpha_n \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 0 \end{pmatrix}$

$\nearrow \det \begin{pmatrix} y_1 & \dots & y_n \\ \vdots & \ddots & \vdots \\ y_1^{(n-1)} & \dots & y_n^{(n-1)} \end{pmatrix} \neq 0$

Wronskian iff $\exists$ only one solution $(\alpha_1 = \alpha_2 = \dots = \alpha_n = 0)$

$y \cdot \alpha = 0 \quad w(y_1, \dots, y_n) \neq 0 \rightarrow \alpha_1 = \alpha_2 = \dots = \alpha_n = 0 \rightarrow y_1, \dots, y_n \text{ L.I.}$

$w = 0 \longrightarrow y_1, \dots, y_n \text{ L.D.}$

Why do we care about linear independence:

N^{th} ord. Lin hom IVP

1) Find N solutions to $\mathcal{L}: y_1, y_2, \dots, y_n$

2) $y = c_1 y_1 + c_2 y_2 + \dots + c_n y_n$

3) apply initial conditions

$y(t_0) = c_1 y_1(t_0) + c_2 y_2(t_0) + \dots + c_n y_n(t_0) = a_0$

$y'(t_0) = c_1 y_1'(t_0) + \dots + c_n y_n'(t_0) = a_1$

$y^{(n-1)}(t_0) = c_1 y_1^{(n-1)}(t_0) + \dots + c_n y_n^{(n-1)}(t_0) = a_{n-1}$

initial conditions

$$\begin{pmatrix} y_1(t_0) & \dots & y_n(t_0) \\ \vdots & & \vdots \\ y_1^{(n-1)} & \dots & y_n^{(n-1)} \end{pmatrix} \begin{pmatrix} c_1 \\ \vdots \\ c_n \end{pmatrix} = \begin{pmatrix} a_0 \\ \vdots \\ a_{n-1} \end{pmatrix}$$

Example:

$$y'' + \frac{1}{t}y' - \frac{q}{t^2}y = 0$$

$$y_1 = t^3 \quad y_2 = \frac{1}{t^3}$$

$$w \begin{vmatrix} t^3 & t^{-3} \\ 3t^2 & -3t^{-4} \end{vmatrix} = \frac{-3}{t} - \frac{3}{t} = \frac{-6}{t} \neq 0$$

$$y^{(n)} + a_{n-1}(t)y^{(n-1)} + \dots + a_0(t)y = 0$$

n solutions $y_1, y_2, \dots, y_n$

Calculate $w(y_1, \dots, y_n)$

Abel's Theorem:

w satisfies

$$w' + a_{n-1}w = 0 \quad \begin{matrix} \nearrow w(t) = w(t_0) e^{\int_{t_0}^t a_{n-1}(s) ds} \\ \searrow w(t) = 0 \end{matrix}$$

Method of Undetermined Coefficients

- If the LHS is constant coefficients and linear and the RHS is an exp, cos, sin, polynomial, or a product of these, try a y_p of the same form as the RHS plus all possible derivatives, each with a coefficient, multiply by t until you have no duplication with y_c , and find the coeff.

Ex: $y'' + y = e^t$

$y_p = Ae^t$

$y_p'' = Ae^t$

$Ae^t + Ae^t = e^t$

$A + A = 1$

$A = \frac{1}{2}$

Ex: $y'' + y = t^2$

$y_p = At^2 + Bt + C$

$y_p' = 2At + B$

$2At + B + At^2 + Bt + C = t^2$

$A = 1$

$2A + B = 0 \rightarrow B = -2$

$B + C = 0 \rightarrow C = 2$

Ex $y'' + y' + y = \cos(2t)$

$y_p = A\cos(2t) + B\sin(2t)$

$y_p' = -2A\sin(2t) + 2B\cos(2t)$

$y_p'' = -4A\cos(2t) - 4B\sin(2t)$

$-4Ac - 4Bs - 2As + 2Bc + Ac + Bs = C$

$\cos(2t): -4A + 2B + A = 1 \quad -3A + 2B = 1$

$\sin(2t): -4B - 2A + B = 0 \quad -2A - 3B = 0$

$A = -\frac{3}{13}, B = \frac{2}{13}$

$y_p = -\frac{3}{13}\cos(2t) + \frac{2}{13}\sin(2t)$

Ex: $y'' - y = e^t$

$y_p = Ae^t$

$y_p'' = Ae^t$

$Ae^t - Ae^t = e^t$
 $0 = e^t$
problem

$y_c = y'' - y = 0$

$y_c = e^{rt}$

$r^2 - 1 = 0$

$r = \pm 1$

$y_c = C_1 e^t + C_2 e^{-t}$

$y_p = Ae^t t$

$y_p' = Ae^t + Ate^t$

$y_p'' = 2Ae^t + Ate^t$

$2Ae^t + Ate^t - Ae^t t = e^t$

$2Ae^t = e^t$

$2A = 1$

$A = \frac{1}{2}$

Ex:

$y'' - 2y' + y = e^t$

Let's find y_c first:

$y'' - 2y' + y = 0$

$e^{rt} \rightarrow r^2 - 2r + 1 = 0$

$(r-1)^2 = 0$

$r = 1$ twice

$y_c = C_1 e^t + C_2 t e^t$

$y_p = Ae^t$

duplicates

$y_p = Ae^t t$

duplicates

no duplication

$y_p = Ae^t t^2$

Example: $y''' - y' = \cos(2t) + t^2$

$y_c: y''' - y' = 0$

$e^{rt} \rightarrow r^3 - r = 0$

$r(r^2 - 1) = 0$

$r = 0, 1, -1$

$y_c = C_1 + C_2 e^t + C_3 e^{-t}$

$y_p = (A \cos(2t) + B \sin(2t)) + (Ct^2 + Dt + E)t$

similar

must multiply by t

Example: ... = $t^2 \cos(3t)$

$$y_p = At^2 \cos(3t) + Bt^2 \sin(3t) + Ct \cos(3t) + Dt \sin(3t) + E \cos(3t) + F \sin(3t)$$

↑
then
add
all
possible derivatives

Ex: ... = $\ln(t)$

$$y_p = A \ln(t) + \frac{B}{t} + \frac{C}{t^2} + \frac{D}{t^3} + \dots$$

this method does not work for $\ln$.

Annihilator Method

- LHS is const coeff
- RHS is Exp, sin, cos, polynomial, or product of these

$$Ly = f(t)$$

1. Find y_c

$$Ly = 0$$

$$y_c = e^{rt} \rightarrow r = r_1, \dots, r_n$$

$$y_c = c_1 e^{r_1 t} + \dots + c_n e^{r_n t}$$

2. Find an annihilator for $f(t)$ st.

$$\tilde{L}f = 0$$

3. Apply $\tilde{L}$ to the original eq.

$$\tilde{L}Ly = \tilde{L}f = 0$$

$$\underbrace{\tilde{L}L}_{\text{const coeffs}} y = 0$$

4. Use $y = e^{rt} \rightarrow r = r_1, \dots, r_n, r_{n+1}, \dots$
 $e^{r_1 t}, \dots, e^{r_n t}, e^{r_{n+1} t}, \dots$

get rid of the ones already in y_c

$$y_p = A e^{r_{n+1} t} + \dots + K e^{r_k t}$$

and find $A, \dots, K$

Example: $y'' - 4y = 4e^{2t}$

$y_c: y'' - 4y = 0$

$y = e^{rt} \rightarrow r^2 - 4 = 0$
 $r = \pm 2$

$y_c = C_1 e^{2t} + C_2 e^{-2t}$

$y_p: \tilde{L} e^{2t} = 0$

$(D-2)e^{2t} = 2e^{2t} - 2e^{2t} = 0$

↑
derivative
operator

$\tilde{L} = (D-2)$

plug into original eq:

$y'' - 4y = 4e^{2t}$
 $(D^2 - 4)y = 4e^{2t}$ > rewrote

$\hat{L}$

$(D-2)(D^2-4)y = (D-2)4e^{2t} = 0$

$(D-2)(D^2-4)y = 0$

$y = e^{rt}$

$(r-2)(r^2-4)e^{rt} = 0$

$r = 2, \pm 2$

e^{2t}, te^{2t}, e^{-2t}

already have in y_c

$y_p = Ate^{2t} \rightarrow$ plug into original equation

$A=1 \rightarrow y_p = te^{2t}$

$$y = y_c + y_p$$

$$y = C_1 e^{2t} + C_2 e^{-2t} + t e^{2t}$$

Example: $y''' + y'' = 3te^t + 2t^2$

$$y_c = y''' + y'' = 0$$

$$y = e^{rt} \rightarrow r^4 + r^3 = 0$$

$$r^3(r+1) = 0$$

$$y = C_1 + C_2 t + C_3 \cos(t) + C_4 \sin(t)$$

$$y_p: 3te^t$$

$$(D-1)^2 te^t = 0$$

$$(D^4 + D^3)y = 3te^t$$

$$(D-1)^3 (D^4 + D^3)y = (D-1)^3 3te^t = 0$$

$$y = e^{rt} \rightarrow (r-1)^3 (r^4 + r^3) = 0$$

$$(r-1)^3 r^3 (r+1) = 0$$

$$r = 1, 1, 1, 0, 0, \pm i$$

$$e^t, te^t, \cancel{x}, \cancel{t}, \cos(t), \sin(t)$$

$$y_{p1} = A + Bte^t$$

$$y_{p2}: 2t^2$$

$$D^3(2t^2) = 0$$

$$(D^3)(D^4 + D^3)y = (D^3)2t^2 = 0$$

$$y = e^{rt} \rightarrow r^3(r^4 + r^3) = 0$$

$$r^6(r^2+1)$$

$$r = 0, 0, 0, 0, 0, \pm i$$

$$\cancel{x}, \cancel{t}, t^2, t^3, t^4, \cos(t), \sin(t)$$

already in y_c

$$y_{p2} = (t^2 + Dt^3 + Et^4)$$

$$y_p = y_{p1} + y_{p2} \rightarrow \text{find coeff}$$

$$y = y_p + y_c$$

Variation of Parameters

Ex: $y'' + y = \ln(t)$

$y_c: y'' + y = 0$

$y_c = C_1 \cos(t) + C_2 \sin(t)$

$y_p = \mu_1(t) \cos(t) + \mu_2(t) \sin(t)$

$y_p = \mu_1'(t) \cos - \mu_1(t) \sin + \mu_2'(t) \sin + \mu_2 \cos$

$\mu_1' C + \mu_2' S = 0 \quad \leftarrow \text{add this constraint (don't want anything higher than } \mu')$

$y_p'' = -\mu_1' S - \mu_1 \cos + \mu_2' \cos - \mu_2 \sin$

plug into original equation $\rightarrow$

$-\mu_1' \sin - \cancel{\mu_1 \cos} + \mu_2' \cos - \cancel{\mu_2 \sin} + \cancel{\mu_1 \cos} + \cancel{\mu_2 \sin} = \ln(t)$

$-\mu_1' \sin + \mu_2' \cos = \ln(t) \rightarrow \text{put next to constraint:}$

$$\begin{cases} \mu_1' \cos + \mu_2' \sin = 0 \\ -\mu_1' \sin + \mu_2' \cos = \ln(t) \end{cases}$$

sub $\hookrightarrow \begin{cases} \mu_2' = \frac{-\cos}{\sin} \mu_1' \\ -\mu_1' \sin - \frac{\cos^2}{\sin} \mu_1' = \ln(t) \end{cases}$

$$\begin{cases} \mu_1' = -\sin \ln(t) \\ \mu_2' = \cos \ln(t) \end{cases}$$

$\mu_1 = \int -\sin \ln(t) dt$

$\mu_2 = \int \cos \ln(t) dt$

plug back into y_p

$y_p = -\cos(t) \int \sin(z) \ln(z) dz + \sin(t) \int \cos(z) \ln(z) dz$

$y = y_c + y_p$

$\mu_1 = - \int \frac{y_a f}{\omega(y_1, y_2)} dt$

$\mu_2 = \int \frac{y_b f}{\omega(y_1, y_2)} dt$

Laplace Transform

(Note: is a linear function
because $L(f) + L(g) = L(f+g)$)

Given $f(t)$,

$$L(f) = \int_0^{\infty} e^{-st} f(t) dt$$

Ex:

$$L(1): \int_0^{\infty} e^{-st} dt = \left. \frac{-1}{s} e^{-st} \right|_0^{\infty} = \frac{1}{s} \quad s > 0$$

$$L(t): \int_0^{\infty} e^{-st} t dt = \left. \frac{-1}{s} e^{-st} t \right|_0^{\infty} + \frac{1}{s} \int_0^{\infty} e^{-st} 1 dt = \left. \frac{-1}{s^2} e^{-st} \right|_0^{\infty} = \frac{1}{s^2}$$

$$L(e^{at}): \int_0^{\infty} e^{-st} e^{at} dt = \int_0^{\infty} e^{-(s-a)t} dt = \left. \frac{-1}{s-a} e^{-(s-a)t} \right|_0^{\infty} = \frac{1}{s-a} \quad s > a$$

$$L(e^{i\omega t}) = \frac{1}{s - i\omega} = \frac{s + i\omega}{s^2 + \omega^2} = \frac{s}{s^2 + \omega^2} + i \frac{\omega}{s^2 + \omega^2}$$

$$L(\cos(\omega t) + i \sin(\omega t)) = L(\cos(\omega t)) + i L(\sin(\omega t))$$

$$\begin{aligned} L(f') &= \int_0^{\infty} e^{-st} f'(t) dt = \\ &= \left. e^{-st} f(t) \right|_0^{\infty} - \int_0^{\infty} (-s) e^{-st} f(t) dt \\ &= -f(0) + sL(f) \end{aligned}$$

$$\begin{aligned} L(f'') &= sL(f') - f'(0) \\ &= s^2 L(f) - sf(0) - f'(0) \end{aligned}$$

I.V.P.

$$Y = L(y)$$

$$y' + 3y = 0$$

$$y(0) = 2$$

$$L(y') + 3L(y) = 0$$

$$sY - y(0) + 3Y = 0$$

$$(s+3)Y - 2 = 0$$

$$Y = \frac{2}{s+3}$$

$$y = 2e^{-3t}$$

$$\begin{cases} y' + 3y = t \\ y(0) = 2 \end{cases}$$

$$sY - y(0) + 3Y = \frac{1}{s^2}$$

$$(s+3)Y - 2 = \frac{1}{s^2}$$

$$(s+3)Y = \frac{1}{s^2} + 2$$

$$Y = \frac{1}{s^2(s+3)} + \frac{2}{s+3}$$

$$\begin{aligned} \frac{1}{s^2(s+3)} &= \frac{A}{s+3} + \frac{B}{s} + \frac{C}{s^2} = \frac{As^2 + B(s+3s) + C(s+3)}{(s+3)s^2} \\ &= \frac{3C + (3B+C)s + (A+B)s^2}{(s+3)s^2} \end{aligned}$$

$$3C = 1$$

$$C = \frac{1}{3}$$

$$3B + C = 0$$

$$B = -\frac{1}{9}$$

$$A + B = 0$$

$$A = \frac{1}{9}$$

$$Y = \frac{1}{9} \cdot \frac{1}{s+3} - \frac{1}{9} \cdot \frac{1}{s} + \frac{1}{3} \cdot \frac{1}{s^2} + \frac{2}{s+3}$$

$$= \frac{19}{9} \cdot \frac{1}{s+3} - \frac{1}{9} \cdot \frac{1}{s} + \frac{1}{3} \cdot \frac{1}{s^2}$$

$$y = \frac{19}{9} e^{-3t} - \frac{1}{9} + \frac{1}{3} t$$

$$y'' - y' - 2y = 2e^t$$

$$y(0) = 3$$

$$y'(0) = 4$$

$$s^2 Y - sy(0) - y'(0) - [sY - y(0)] - 2Y = \frac{2}{s-1}$$

$$(s^2 - s - 2)Y - s \cdot 3 - 4 + 3 = \frac{2}{s-1}$$

$$(s-2)(s+1)Y = \frac{2}{s-1} + 3s + 1$$

$$Y = \frac{2}{(s-1)(s-2)(s+1)} + \frac{3s}{(s-2)(s+1)} + \frac{1}{(s-2)(s+1)}$$

$$Y = \frac{-1}{s-1} + \frac{3}{s-2} + \frac{1}{s+1}$$

$$y = -e^t + 3e^{2t} + e^{-t}$$

Definitions

$$H(t-a) = \begin{cases} 0 & t < a \\ 1 & t \geq a \end{cases}$$

Heaviside

$$\mathcal{L}(H(t-a)) = \frac{1}{s} e^{-as}$$

$$\mathcal{L}(f(t)H(t-a)) = \frac{1}{s} e^{-as} f(t+a)$$

Watch Monday Lecture

Free Damped Oscillator

$$(F_{\text{external}} = 0) \quad (\gamma > 0)$$

$$Mx'' + \gamma x' + Kx = 0$$

$$x = e^{rt}$$

$$Mr^2 + \gamma r + K$$

$$r = \frac{-\gamma \pm \sqrt{\gamma^2 - 4MK}}{2M}$$

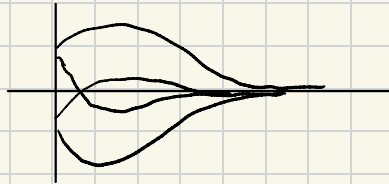
$$M, K, \gamma > 0$$

If $\gamma^2 > 4MK$ (over damped)

$$R_1 \neq R_2 < 0$$

$$x = A e^{\left(\frac{-\gamma}{2M} + \frac{\sqrt{\gamma^2 - 4MK}}{2M}\right)t} + B e^{\left(\frac{-\gamma}{2M} - \frac{\sqrt{\gamma^2 - 4MK}}{2M}\right)t}$$

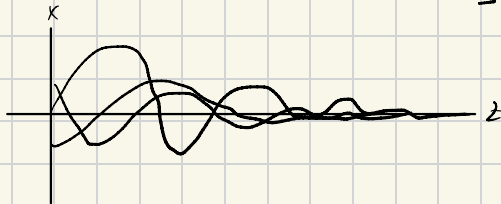
$$= e^{\frac{-\gamma}{2M}t} \left(A e^{\frac{\sqrt{\gamma^2 - 4MK}}{2M}t} + B e^{-\frac{\sqrt{\gamma^2 - 4MK}}{2M}t} \right)$$



If $\gamma^2 < 4MK$ (underdamped)

$$r = \frac{-\gamma}{2M} \pm i \frac{\sqrt{4MK - \gamma^2}}{2M}$$

$$x(t) = e^{\frac{-\gamma}{2M}t} \left[A \cos\left(\frac{\sqrt{4MK - \gamma^2}}{2M}t\right) + B \sin\left(\frac{\sqrt{4MK - \gamma^2}}{2M}t\right) \right]$$

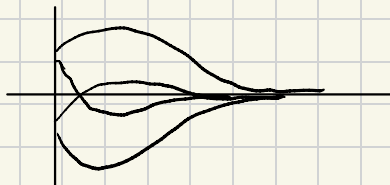


Critically Damped
and overdamped
look similar to the eye,
but have different numbers.

If $\gamma^2 = 4MK$ (critically damped)

$$R_1 = R_2$$

$$x(t) = A e^{\frac{-\gamma}{2M}t} + B t e^{\frac{-\gamma}{2M}t}$$



Forced Damped Oscillator.

$$Mx'' + \gamma x' + Kx = F \cos(\omega t)$$

← sinusoidal force

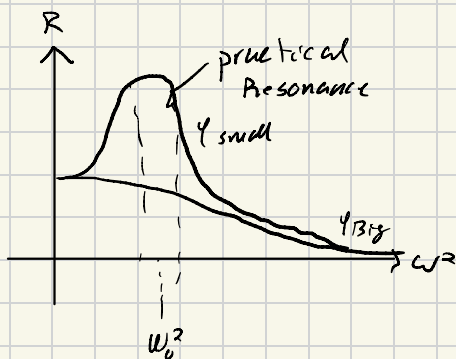
$$x_c: Mx'' + \gamma x' + Kx = 0$$

$$x_c = e^{\frac{-\gamma}{2m}t} [Ae^{\sqrt{\dots}t} + Be^{-\sqrt{\dots}t}]$$

$$\left. \begin{aligned} x_p &= C \cos(\omega t) + D \sin(\omega t) \\ x_p' &= \dots \\ x_p'' &= \dots \end{aligned} \right\} \begin{array}{l} \text{plug into} \\ \text{original} \end{array} \quad \left. \begin{aligned} &= R \cos(\omega t - \delta) ; R = \frac{\frac{F}{m}}{\sqrt{(\omega_0^2 - \omega^2)^2 + \frac{\gamma^2 \omega^2}{m^2}}} \\ &\delta = \tan^{-1}\left(\frac{\gamma}{\omega_0^2 - \omega^2}\right) \end{aligned} \right\}$$

$$C = \frac{\frac{F}{m}(\omega_0^2 - \omega^2)}{(\omega_0^2 - \omega^2)^2 + \frac{\gamma^2 \omega^2}{m^2}} ; D = \frac{\frac{\gamma}{m} \omega \frac{F}{m}}{\dots}$$

$$x = x_c + x_p = \underbrace{e^{\frac{-\gamma}{2m}t} [A \dots + B \dots]}_{\text{Transient } x_{tr}} + \underbrace{R \cos(\omega t - \delta)}_{\text{steady state } x_{ss}}$$



Complexification

$$Mx'' + Kx = F \cos(\omega t) \rightarrow Mx'' + Kx = \text{Re}(F e^{i\omega t}) \rightarrow \text{Re}(M\tilde{x}'') + \text{Re}(K\tilde{x}) = \text{Re}(F e^{i\omega t})$$

$\tilde{x}$ complex

$$x = \text{Re}(\tilde{x})$$

$$\tilde{x} = A e^{i\omega t} ; \tilde{x}'' = -\omega^2 A e^{i\omega t}$$

Matrix Multiplication

$$AB = C$$

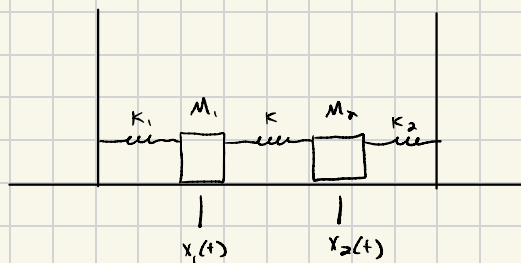
$\begin{matrix} & \swarrow & & \searrow \\ M \times N & N \times P & & M \times P \end{matrix}$

Row-by-column

$$i \quad \left[\begin{array}{|c|} \hline - \\ \hline \end{array} \right] \quad \left[\begin{array}{|c|} \hline j \\ \hline \end{array} \right] = \left[\begin{array}{|c|} \hline ij \\ \hline \end{array} \right]$$

$$\begin{bmatrix} 1 & 4 \\ 2 & 5 \\ 3 & 6 \end{bmatrix} \begin{bmatrix} 1 & 2 \\ 3 & 4 \end{bmatrix} = \begin{bmatrix} 13 & 18 \\ 17 & 24 \\ 21 & 30 \end{bmatrix}$$

Examples



$$\underline{x}(t) = \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix}$$

$$A = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -\frac{(K+K_1)}{M_1} & \frac{K}{M_1} & 0 & 0 \\ \frac{K}{M_2} & -\frac{(K+K_2)}{M_2} & 0 & 0 \end{pmatrix}$$

$$g = \begin{pmatrix} 0 \\ 0 \\ g_1 \\ g_2 \end{pmatrix}$$

$$\underline{x}' = A(t)\underline{x} + g$$

$$\begin{cases} M_1 x_1'' = -K_1 x_1 - K x_1 + K x_2 + g_1(t) \\ M_2 x_2'' = K x_1 - K x_2 - K_2 x_2 + g_2(t) \end{cases}$$

$$\text{let } \begin{matrix} x_3 = x_1' & x_3' = x_1'' \\ x_4 = x_2' & x_4' = x_2'' \end{matrix}$$

$$x_1' = x_3$$

$$x_2' = x_4$$

$$x_3' = -\frac{(K+K_1)x_1}{M_1} + \frac{Kx_2}{M_1} + \frac{g_1(t)}{M_1}$$

$$x_4' = \frac{K}{M_2}x_1 - \frac{(K+K_2)x_2}{M_2} + \frac{g_2(t)}{M_2}$$

Ex 5

$$y'' + t^2 y' + y y' = \cos(t)$$

$$y_1 = y \quad y_1' = y_2$$

$$y_2 = y' = y_1' \rightarrow y_2' = y_3$$

$$y_3 = y'' = y_2' \quad y_3' = -t^2 y_3 - y_1 y_2 + \cos(t)$$

$$y(t) = \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix} \quad F = \begin{pmatrix} y_2 \\ y_3 \\ -t^2 y_3 - y_1 y_2 + \cos(t) \end{pmatrix} \quad y' = F(t, y)$$

Theorem of Existence & Uniqueness

For $\begin{cases} y' = F(t, y) \\ y(t_0) = y_0 \end{cases}$ • If $\exists$ a region around (t_0, y_0)
 $\forall F(t, y)$ is continuous $\rightarrow$ a solution exists
around (t_0, y_0)

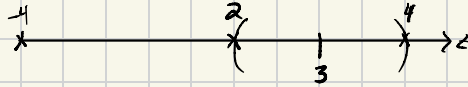
• And if

$\frac{\partial F}{\partial y_1}, \frac{\partial F}{\partial y_2}, \dots, \frac{\partial F}{\partial y_n}, \frac{\partial F}{\partial y_1}, \dots, \frac{\partial F}{\partial y_n}$ are continuous $\rightarrow$ sol
is unique

Ex1 $y_1' = t^2 y_1 + \frac{1}{t-2} y_2 + \cos(t)$

$y_2' = \frac{t}{t+4} y_1 + \sin(t) y_2 + \frac{1}{t-4}$

$\begin{cases} y_1(3)=0 \\ y_2(3)=1 \end{cases}$



$I = (2, 4)$

Lin Hom Syst

$y' = A(t)y$

suppose we have n -sols $v_1, v_2, \dots, v_n$

$v_1 = \begin{pmatrix} 1 \\ t \\ t^2 \end{pmatrix} \quad v_2 = \begin{pmatrix} 1 \\ t^2-1 \\ t^2+1 \end{pmatrix}$

$\text{Det} \begin{pmatrix} v_1 & v_2 & \dots & v_n \end{pmatrix} = w(t)$
Wronskian

if $w(t) \neq 0$, $v_1, v_2, \dots, v_n$ are lin ind

if $w(t) = 0$, $v_1, v_2, \dots, v_n$ are lin dep

Abel's Theorem

$\frac{dw}{dt} = T_r(A) \cdot w(t)$

$\nearrow w(t) = w(t_0) e^{\int_{t_0}^t T_r(A) dt} \rightarrow v_1, v_2, \dots, v_n \text{ lin ind}$
 $\searrow w = 0 \rightarrow v_1, v_2, \dots, v_n \text{ lin dep}$

$$y' = Ay + f$$

$$y(t_0) = y_0$$

$$M(t) = \begin{cases} \phi'(t) = A\phi \\ \phi(0) = I \end{cases}$$

$$\begin{cases} y' = ay \rightarrow y = e^{at} \\ y(0) = 1 \end{cases}$$

$$e^{at} = 1 + at + \frac{(at)^2}{2!} + \frac{(at)^3}{3!} + \dots$$

$$e^{At} = I + At + \frac{A^2 t^2}{2!} + \dots = I + \sum_{k=1}^{\infty} \frac{A^k t^k}{k!}$$

$$e^{At} \Big|_{t=0} = I, (e^{At})' = 0 + \sum_{k=1}^{\infty} \frac{A^k k t^{k-1}}{k!} =$$

$$A + \sum_{k=2}^{\infty} \frac{A^k k t^{k-1}}{(k-1)!} = A + A \sum_{k=2}^{\infty} \frac{A^{k-1} t^{k-1}}{(k-1)!} =$$

$$A \left(I + \sum_{k=1}^{\infty} \frac{A^k t^k}{k!} \right) = A e^{At}$$

$A_{n \times n}$ is Diagonalizable

If it has N lin. ind. eigenvectors

$\varphi_1, \varphi_2, \varphi_3, \dots, \varphi_n \rightarrow$ eigenvectors

$\lambda_1, \lambda_2, \lambda_3, \dots, \lambda_n \rightarrow$ eigenvalues

$$(\varphi_1, \varphi_2, \dots, \varphi_n) = U \quad \Lambda = \begin{pmatrix} \lambda_1 & 0 & \dots \\ 0 & \lambda_2 & \dots \\ \vdots & \vdots & \ddots \\ 0 & \dots & \lambda_n \end{pmatrix}$$

$$AU = U\Lambda$$

$$AUU^{-1} = U\Lambda U^{-1}$$

* < $A = U\Lambda U^{-1} \rightarrow$ given A is diagonalizable

$$U^{-1}AU = \Lambda$$

$$A^2 = A A = U \cancel{\Lambda} U^{-1} = U \Lambda^2 U^{-1}$$

$$\begin{aligned} \phi(t) = e^{At} &= I + \sum_{k=1}^{\infty} \frac{A^k t^k}{k!} = I + \sum_{k=1}^{\infty} U \Lambda^k U^{-1} \frac{t^k}{k!} = U \left[I + \sum_{k=1}^{\infty} \frac{\Lambda^k t^k}{k!} \right] U^{-1} = \\ &= U \begin{pmatrix} e^{\lambda_1 t} & 0 & \dots \\ 0 & e^{\lambda_2 t} & \dots \\ \vdots & \vdots & e^{\lambda_n t} \end{pmatrix} U^{-1} \\ &= e^{At} = \phi \end{aligned}$$

Ex1)

$$y' = A_{\text{new}} y \quad A = \begin{pmatrix} 1 & 1 & 0 \\ 1 & 0 & 3 \end{pmatrix}$$

$$\lambda_1 = -1$$

$$\lambda_2 = \lambda_3 = 3$$

$$\begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}$$

$$\begin{pmatrix} 1 \\ 1 \\ 0 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$U = \begin{pmatrix} 1 & 1 & 0 \\ -1 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$U^{-1} = \frac{1}{2} \begin{pmatrix} 1 & -1 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 2 \end{pmatrix}$$

$$\Lambda = \begin{pmatrix} -1 & 0 & 0 \\ 0 & 3 & 0 \\ 0 & 0 & 3 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 1 & 0 \\ -1 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{-t} & 0 & 0 \\ 0 & e^{3t} & 0 \\ 0 & 0 & e^{3t} \end{pmatrix} \frac{1}{2} \begin{pmatrix} 1 & -1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 2 \end{pmatrix}$$

$$= \begin{pmatrix} e^{-t} & e^{3t} & 0 \\ e^{-t} & e^{3t} & 0 \\ 0 & 0 & e^{3t} \end{pmatrix} \frac{1}{2} \begin{pmatrix} 1 & -1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 2 \end{pmatrix}$$

$$= \frac{1}{2} \begin{pmatrix} e^{-t} + e^{3t} & -e^{-t} + e^{3t} & 0 \\ -e^{-t} + e^{3t} & e^{-t} + e^{3t} & 0 \\ 0 & 0 & 2e^{3t} \end{pmatrix} = \phi(t)$$

$$y(t) = \phi(t) y_0$$

Putzer's Algorithm

Another way of finding the exp. of a matrix

e^{At} is the fundamental solution matrix of $\frac{\partial \vec{v}}{\partial t} = A\vec{v}$ which $= I_{n \times n}$ at $t=0$

$$e^{At} = I + At + \frac{A^2 t^2}{2!} + \dots = \sum_{k=0}^{\infty} \frac{A^k t^k}{k!}$$

If $\vec{v}_1, \dots, \vec{v}_n$ is a complete set of eigenvectors for A w/ $A\vec{v}_j = \lambda_j \vec{v}_j$

and $U = [\vec{v}_1 \dots \vec{v}_n]$ then $A = U \begin{bmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_n \end{bmatrix} U^{-1}$

$$e^{At} = U \begin{bmatrix} e^{\lambda_1 t} & & 0 \\ & \ddots & \\ 0 & & e^{\lambda_n t} \end{bmatrix} U^{-1}$$

Cayley-Hamilton Theorem

If A is an $n \times n$ matrix

and $\lambda_1, \dots, \lambda_n$ is its eigenvalues

and we define

$$B_0 = I$$

$$B_1 = A - \lambda_1 I$$

$$B_2 = (A - \lambda_2 I) B_1$$

$$\vdots$$
$$B_n = (A - \lambda_n I) B_{n-1}$$

$$\frac{1}{2} \int_0^2 (2-x) dx = \frac{1}{2} \left[2x - \frac{x^2}{2} \right]_0^2 = \frac{1}{2} [4 - 2]$$

①

$$A_n = \frac{a}{L} \int_L f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$

$$A_n = \int_0^2 (2-x) \cos\left(\frac{n\pi x}{2}\right) dx$$

$$\text{let } u = 2-x \quad v = \frac{2}{n\pi} \sin\left(\frac{n\pi x}{2}\right)$$

$$du = -dx \quad dv = \cos\left(\frac{n\pi x}{2}\right) dx$$

$$= (2-x) \left(\frac{2}{n\pi} \right) + \int \frac{2}{n\pi} \sin\left(\frac{n\pi x}{2}\right) dx$$

$\uparrow$
 $\sin\left(\frac{n\pi x}{2}\right)$

$$\frac{2}{n\pi} \left[(2-x) \sin\left(\frac{n\pi x}{2}\right) - \left(\frac{2}{n\pi}\right) \cos\left(\frac{n\pi x}{2}\right) \right]_0^2$$

$$\frac{2}{n\pi} \left[0 - \frac{2}{n\pi} \cos(n\pi) - 0 + \frac{2}{n\pi} \right]$$

$$\frac{2}{n\pi} \left[\frac{2}{n\pi} - \frac{2}{n\pi} \cos(n\pi) \right]$$

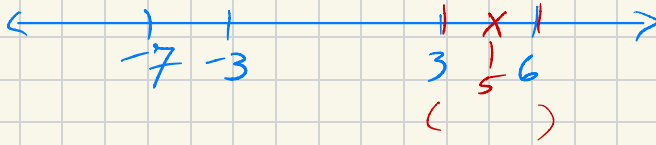
$$\left(\frac{2}{n\pi} \right)^2 [1 - \cos(n\pi)]$$

$$A_n = \frac{4[1 - \cos(n\pi)]}{(n\pi)^2}$$

$$\begin{vmatrix} 4-\lambda & 3 \\ 2 & 3-\lambda \end{vmatrix} = 0$$

$$(4-\lambda)(3-\lambda) - 6 = 0$$

12



$$r^2 + r - 6 = 0$$

$$(r$$

$$r = -3, 2$$

$$y_p = At^2 + Bt + C$$

$$y_p' = 2At + B$$

$$y_p'' = 2A$$

$$y_p = -2t^2 + t + 1$$

$$y_h = C_1 e^{-3t} + C_2 e^{2t}$$

$$2A + 2At + B - 6At^2 - 6Bt - 6C = 12t^2 - 10t - 9$$

$$-6A = 12$$

$$A = -2$$

$$B = 1$$

$$C = 1$$

$$y = y_p + y_h$$

$$y = C_1 e^{-3t} + C_2 e^{2t} - 2t^2 + t + 1$$

$$1 = C_1 + C_2 + 1$$

$$C_1 = -C_2$$

$$0 = C_1 e^{-3} - C_1 e^2 - 2 + 1 + 1$$

$$C_1 e^{-3} = C_1 e^2$$

$$C_1 = 0 \quad C_2 = 0$$

$$(D^2 + \lambda) y$$

$$D = \pm \sqrt{\lambda} i$$

$$\underline{\lambda > 0}$$

$$y = C_1 \cos(\sqrt{\lambda} x)$$

$$y' = -C_1 \sqrt{\lambda} \sin(\sqrt{\lambda} x) + C_2 \sqrt{\lambda} \cos(\sqrt{\lambda} x)$$

$$r^2 + 1 = 0$$

$$r = \pm i$$

$$y_h = C_1 \cos(x) + C_2 \sin(x)$$

$$a_0 = C_1$$

$$a = -a_0$$