Modern Spring 2006

Problem A1: A highly energetic photon (wavelength λ) collides with an electron at rest. After the collision the angle θ between the incident photon direction and the final direction is measured. Calculate the change in the wavelength, $\Delta \lambda = \lambda' - \lambda$ in terms of θ , where λ' is the scattered photon wavelength. What is particularly remarkable about this result?



The 4-momentum of the incident photon and electron are:

$$p_{\gamma}^{\mu} = (E_{\gamma}, c\boldsymbol{p}_{\gamma}) = \left(\frac{hc}{\lambda}, c\boldsymbol{p}_{\gamma}\right)$$
$$p_{e}^{\mu} = (m_{e}c^{2}, \mathbf{0})$$

Since the photon is massless $|c\mathbf{p}_{\gamma}| = E_{\gamma} = \frac{hc}{\lambda}$. The 4-momentum of the outgoing photon and electron are:

$$p_{\gamma}^{\prime \mu} = (E_{\gamma}^{\prime}, c\boldsymbol{p}_{\gamma}^{\prime}) = \left(\frac{hc}{\lambda^{\prime}}, c\boldsymbol{p}_{\gamma}^{\prime}\right)$$
$$p_{e}^{\prime \mu} = (m, c\boldsymbol{p}_{e}^{\prime})$$

As before, $|c p'_{\gamma}| = E'_{\gamma} = \frac{hc}{\lambda'}$. By 4-momentum conservation:

$$p_{\gamma}^{\mu} + p_{e}^{\mu} = p_{\gamma}^{\prime \mu} + p_{e}^{\prime \mu}$$

$$p_{\gamma}^{\mu} - p_{\gamma}^{\prime \mu} = p_{e}^{\prime \mu} - p_{e}^{\mu}$$
(1)

Square both sides of Eq. (1):

$$p_{\gamma}^{2} - 2p_{\gamma} \cdot p_{\gamma}' + {p_{\gamma}'}^{2} = p_{e}^{2} - 2p_{e} \cdot p_{e}' + {p_{e}'}^{2}$$

(Use on mass shell condition)
$$0 - 2p_{\gamma} \cdot p_{\gamma}' + 0 = m_{e}^{2}c^{4} - 2p_{\gamma} \cdot p_{\gamma}' + m_{e}^{2}c^{4}$$
$$-2\left(E_{\gamma}E_{\gamma}' - c\boldsymbol{p}_{\gamma} \cdot c\boldsymbol{p}_{\gamma}'\right) = 2m_{e}^{2}c^{4} - 2\left(E_{e}E_{e}' - c\boldsymbol{p}_{e} \cdot c\boldsymbol{p}_{e}'\right)$$

$$-2(E_{\gamma}E_{\gamma}'-E_{\gamma}E_{\gamma}'\cos\theta) = 2m_e^2c^4 - 2(E_eE_e'-0)$$

Here we used
$$c\mathbf{p}_{\gamma} \cdot c\mathbf{p'}_{\gamma} = |c\mathbf{p}_{\gamma}| |c\mathbf{p'}_{\gamma}| \cos\theta = E_{\gamma}E'_{\gamma}\cos\theta$$

 $-2E_{\gamma}E'_{\gamma}(1-\cos\theta) = 2m_e^2c^4 - 2(m_ec^2E'_e-0)$
 $(\div -2 \text{ and use } E = \frac{hc}{\lambda} \text{ for } \gamma's)$
 $\frac{hc}{\lambda}\frac{hc}{\lambda'}(1-\cos\theta) = -m_e^2c^4 + m_ec^2E'_e$
 $\frac{(hc)^2}{\lambda\lambda'}(1-\cos\theta) = -m_e^2c^4 + m_ec^2E'_e$
(2)

Now, from conservation of energy:

$$m_e c^2 + E_{\gamma} = E'_e + E'_{\gamma}$$

$$E'_e = m_e c^2 + E_{\gamma} - E'_{\gamma}$$

$$\rightarrow E'_e = m_e c^2 + \frac{hc}{\lambda} - \frac{hc}{\lambda'}$$
(3)

Insert Eq. (3) into (2):

$$\frac{(hc)^2}{\lambda\lambda'}(1-\cos\theta) = -m_e^2c^4 + m_ec^2\left(m_ec^2 + \frac{hc}{\lambda} - \frac{hc}{\lambda'}\right)$$
$$\frac{(hc)^2}{\lambda\lambda'}(1-\cos\theta) = -m_e^2c^4 + m_e^2c^4 + m_ec^2\left(\frac{hc}{\lambda} - \frac{hc}{\lambda'}\right)$$
$$\frac{(hc)^2}{\lambda\lambda'}(1-\cos\theta) = m_ec^2\left(\frac{hc}{\lambda} - \frac{hc}{\lambda'}\right)$$
$$\frac{(hc)^2}{\lambda\lambda'}(1-\cos\theta) = m_ec^3h\left(\frac{1}{\lambda} - \frac{1}{\lambda'}\right)$$
$$\frac{(hc)^2}{\lambda\lambda'}(1-\cos\theta) = m_ec^3h\left(\frac{\lambda'-\lambda}{\lambda\lambda'}\right)$$
$$h(1-\cos\theta) = m_ec(\lambda'-\lambda)$$

$$\rightarrow \quad \Delta \lambda = \frac{h}{m_e c} (1 - \cos \theta)$$

Since $(1 - \cos \theta) \ge 0$ the photon always increases its wavelength, hence always loses energy. The maximum energy loss occurs at $\cos \theta = -1$, backward scattering. Finally, $\Delta \lambda \sim \frac{1}{m_e}$ and thus in general the more massive the object the less change in wavelength you'll see in Compton scattering.

The Compton effect shows light cannot be purely explained by waves.