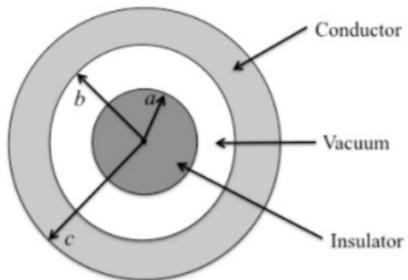


EM2: An insulating sphere of radius a carries a charge density of

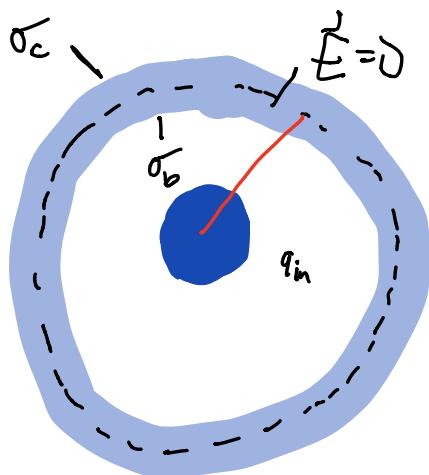
$$\rho = \frac{k}{r}$$

in the region $r \leq a$, where k is a positive constant and r is the distance from the center of the sphere. The sphere is surrounded by a thick, concentric conducting metal shell with an inner radius b and an outer radius c . The conducting shell carries no net charge.

- Determine the surface charge density at the inner ($r = b$) and outer surfaces ($r = c$) of the conductor.
- Determine the electric field \vec{E} in all four regions (i) $r < a$, (ii) $a < r < b$, (iii) $b < r < c$, and (iv) $r > c$.
- Determine the electric potential V at the center of the sphere using infinity ($r = \infty$) as a reference point.
- If the outer shell is grounded, what would the potential at the center of the sphere be using the same reference point as in part (c).



a) Let q_b be the induced charge on inner surface, $E = 0$ inside conductors.



For $b < r < c$, by Gauss' law then:

$$Q_{\text{enclosed}} = 0$$

$$\Rightarrow q_{\text{in}} + q_b = 0$$

$q_{\text{in}} \rightarrow$ charge of insulator

$q_b \rightarrow$ charge of inner surface of conductor

$$\begin{aligned}
 q_{in} &= \int d^3r \rho(r) \\
 &= \int_{S^2} d\Omega \int_0^a dr r^2 \rho(r) \\
 &= 4\pi \int_0^a kr r^2 dr \\
 &= 4\pi k \left[\frac{r^2}{2} \right]_0^a \\
 q_{in} &= 2\pi k a^2 \Rightarrow q_b = -2\pi k a^2
 \end{aligned}$$

Thus, $\sigma_b = \frac{-2\pi k a^2}{4\pi b^2}$

$$\boxed{\sigma_b = -\frac{k a^2}{2 b^2}}$$

Since the conducting shell has no net charge and (ideal) conductors only have charges on their surfaces, then:

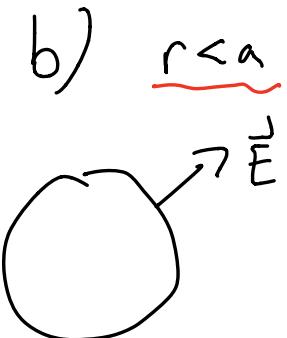
$$q_b + q_c = 0$$

$q_c \rightarrow$ charge of outer surface

Then,

$$q_c = -q_b = 2\pi k a^2$$

$$\Rightarrow \boxed{\sigma_c = \frac{ka^2}{2b^2}}$$



$$\oint \vec{E} \cdot d\vec{a} = Q_{\text{enc}}/\epsilon_0$$

$$E(4\pi r^2) = \frac{1}{\epsilon_0} \int d\sigma \int_0^r dr' r'^2 \left(\frac{\kappa}{r'}\right)$$

$$E(4\pi r^2) = \frac{4\pi k}{\epsilon_0} \left[\frac{r'^3}{3} \right]_0^r$$

$$Er^2 = \frac{k}{2\epsilon_0} r^2$$

$$\Rightarrow \boxed{\vec{E} = \frac{k}{2\epsilon_0} \hat{r}}$$

$a < r < b$

$$\oint \vec{E} \cdot d\vec{a} = q_{\text{in}}/\epsilon_0$$

$$E(4\pi r^2) = \frac{2\pi k a^2}{\epsilon_0}$$

$$\vec{E} = \frac{k}{2\epsilon_0} \frac{a^2}{r^2} \hat{r}$$

$b < r < c$

$$\vec{E} \approx 0$$

(inside conductor)

$c \leq r$

$$\oint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} (q_{in} + q_b + q_c)$$

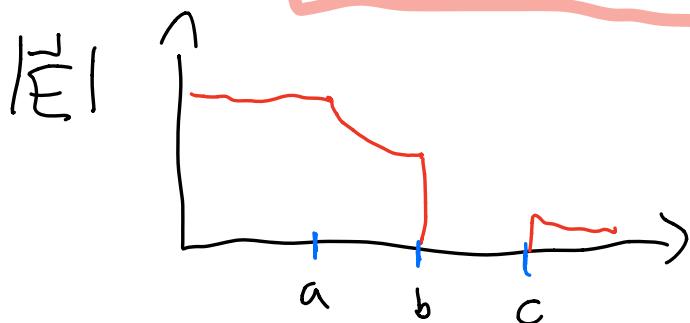
b

$$E(4\pi r^2) = \frac{q_{in}}{\epsilon_0}$$

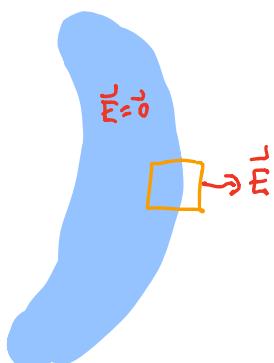
≈ 0 , no
net charge for
the conductor

$$E(4\pi r^2) = \frac{2\pi k a^2}{\epsilon_0}$$

$$\vec{E} = \frac{k}{2\epsilon_0} \frac{a^2}{r^2} \hat{r}$$



Note : drawing a small box at the outer surface of the conductor



$$E \cdot A = \frac{\sigma A}{\epsilon_0} \Rightarrow \sigma = \frac{ka^2}{2b^2}$$

consistent with our result from (a),

$$c) V_0 = - \int_{\infty}^0 \vec{dl} \cdot \vec{E} = - \int_{\infty}^0 dr E(r)$$

$$V_0 = - \int_{\infty}^C dr \frac{k}{2\epsilon_0} \frac{a^2}{r^2} - \int_b^a dr \frac{k}{2\epsilon_0} \frac{a^2}{r^2} - \int_a^0 dr \frac{k}{2\epsilon_0}$$

$$V_0 = \frac{k}{2\epsilon_0} \left(\left[\frac{a^2}{r} \right]_0^C + \left[\frac{a^2}{r} \right]_b^a - \left[r \right]_a^0 \right)$$

$$\approx \frac{k}{2\epsilon_0} \left(\frac{a^2}{C} - 0 + \frac{a^2}{a} - \frac{a^2}{b} - 0 + a \right)$$

$$V_0 = \frac{k a^2}{2\epsilon_0} \left(\frac{1}{C} - \frac{1}{b} + \frac{2}{a} \right)$$