

QM3: A particle of mass m and charge q is in a one-dimensional harmonic oscillator potential moving with frequency ω . In addition, it is subject to a *weak* electric field \mathcal{E} .

- Find the exact expression for the energy.
- Calculate the energy up to the first non-zero correction in non-degenerate perturbation theory. Compare your result to what you found in part a.

a) The Hamiltonian for a 1D HO is

$$\hat{H}_0 = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2} \hat{x}^2$$

Classically, the potential energy in an electric potential is given by:

$$U = qV$$

$$\text{Since } E_x = -\frac{dV}{dx} \Rightarrow V = -E_x x \quad (\text{choose const. of zero})$$

$$U = -qE_x x$$

Thus, the perturbed Hamiltonian is

$$\hat{H}' = -q\mathcal{E} \hat{x}$$

The total Hamiltonian is:

$$\hat{H} = \hat{H}_{tb} + \hat{H}'$$

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2} \hat{x}^2 - q\mathcal{E}\hat{x} \quad (1)$$

We can complete the square so that Eq. (1) resembles a harmonic oscillator.

$$ax^2 + bx = a\left(x + \frac{b}{2a}\right)^2 - \frac{b^2}{4a}$$

$$\Rightarrow \frac{m\omega^2}{2} \hat{x}^2 - q\mathcal{E}\hat{x} = \frac{m\omega^2}{2} \left(\hat{x} + \frac{(-q\mathcal{E})}{2\frac{m\omega^2}{2}} \right)^2 - \frac{q^2\mathcal{E}^2}{4\frac{m\omega^2}{2}}$$

$$\frac{m\omega^2}{2} \hat{x}^2 - q\mathcal{E}\hat{x} = \frac{m\omega^2}{2} \left(\hat{x} - \frac{q\mathcal{E}}{m\omega^2} \right)^2 - \frac{q^2\mathcal{E}^2}{2m\omega^2} \quad (2)$$

Inserting Eq. (2) into (1):

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2} \left(\hat{x} - \frac{q\mathcal{E}}{m\omega^2} \right)^2 - \frac{q^2\mathcal{E}^2}{2m\omega^2}$$

Letting $\hat{y} = \hat{x} - \frac{q\mathcal{E}}{m\omega^2}$ (note momentum doesn't change under spatial translations);

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega^2}{2} \hat{y}^2 - \frac{q^2 \mathcal{E}^2}{2m\omega^2}$$

Thus, the energy is just that of a HO shifted by $-q^2 \mathcal{E}^2 / 2m\omega^2$:

$$E_n = \hbar\omega \left(n + \frac{1}{2}\right) - \frac{q^2 \mathcal{E}^2}{2m\omega^2} \quad (3)$$

b) First order

$$\begin{aligned} E_n^{(1)} &= \langle n^{(0)} | H' | n^{(0)} \rangle \\ &= -q\mathcal{E} \langle n | \hat{x} | n \rangle \end{aligned}$$

Here $|n\rangle$ is the n^{th} eigenstate of a HO,

$$\text{Now: } \hat{x} = \sqrt{\frac{\hbar}{2m\omega}} (\hat{a}^\dagger + \hat{a})$$

where \hat{a}^\dagger, \hat{a} are the creation and annihilation operators:

$$E_n^{(1)} = -q\mathcal{E} \sqrt{\frac{\hbar}{2m\omega}} \left(\langle n | \hat{a}^\dagger | n \rangle + \langle n | \hat{a} | n \rangle \right)$$

Now, $\langle n | \hat{a}^\dagger | n \rangle = \sqrt{n+1} \langle n | n+1 \rangle = 0$

$$\langle n | \hat{a} | n \rangle = \sqrt{n} \langle n | n-1 \rangle = 0$$

Therefore

$$E_n^{(1)} = 0$$

This is expected as the correction in Eq. (b) is second order in \mathcal{E}^2 .

Second order

$$E_n^{(2)} = \sum_{m \neq n} \frac{|\langle n^{(0)} | H' | m^{(0)} \rangle|^2}{E_n^{(0)} - E_m^{(0)}}$$

Now, $E_n^{(0)} - E_m^{(0)} = \hbar\omega(n + \frac{1}{2}) - \hbar\omega(m + \frac{1}{2}) = \hbar\omega(n - m)$

Also,

$$\langle n^{(0)} | H' | m^{(0)} \rangle$$

$$= -q\mathcal{E} \langle n | \hat{x} | m \rangle$$

$$= -q\mathcal{E} \sqrt{\frac{\hbar}{2m\omega}} \langle n | \hat{a}^\dagger + \hat{a} | m \rangle$$

$$= -q\mathcal{E} \sqrt{\frac{\hbar}{2m\omega}} \left(\langle n | \hat{a}^\dagger | m \rangle + \langle n | \hat{a} | m \rangle \right)$$

$$= -q \varepsilon \sqrt{\frac{\hbar}{2m\omega}} \left(\sqrt{m+1} \langle n | m+1 \rangle + \sqrt{m} \langle n | m-1 \rangle \right)$$

$$\langle n^{(0)} | H' | n^{(0)} \rangle = -q \varepsilon \sqrt{\frac{\hbar}{2m\omega}} \left(\sqrt{m+1} \delta_{n, m+1} + \sqrt{m} \delta_{n, m-1} \right)$$

Squaring,

$$|\langle n^{(0)} | H' | n^{(0)} \rangle|^2 = \frac{q^2 \varepsilon^2 \hbar}{2m\omega} \left((m+1) \delta_{n, m+1}^2 + 2\sqrt{(m+1)m} \delta_{n, m+1} \delta_{n, m-1} + m \delta_{n, m-1}^2 \right)$$

We've used $\delta_{m,n}^2 = \delta_{m,n}$.

Thus,

$$E_n^{(2)} = \sum_{m \neq n} \frac{q^2 \varepsilon^2 \hbar}{2m\omega} \left[\frac{(m+1) \delta_{n, m+1} + \sqrt{(m+1)m} \delta_{n, m-1} \delta_{n, m+1} + m \delta_{n, m-1}}{\hbar \omega (n-m)} \right] \quad (4)$$

↑
↑
 mass index in brackets

Bad notation on my part. Don't confuse the two m's.

The middle term in the brackets disappears since it requires both $n=m-1$ and $n=m+1$.

$\delta_{n, m+1} \rightarrow$ fixes $m = n-1$

$\delta_{n, m-1} \rightarrow$ fixes $m = n+1$

Then,

$$E_n^{(2)} = \frac{q^2 \varepsilon^2}{2m\omega^2} \left[\frac{n}{(n - (n-1))} + \frac{n+1}{n - (n+1)} \right]$$

$$E_n^{(2)} = \frac{q^2 \varepsilon^2}{2m\omega^2} [n - (n+1)]$$

$$\Rightarrow E_n^{(2)} = -\frac{q^2 \varepsilon^2}{2m\omega^2}$$

This matches with Eq.(3), as it should.