

**QM4:** A particle with positive energy ( $E > 0$ ) encounters a Dirac delta-function potential barrier

$$V(x) = \alpha \delta(x)$$

where  $\alpha$  is positive real constant with the units of energy $\times$ length.

- Solve the Schrödinger equation for the regions  $x < 0$  and  $x > 0$ . If you make any assumptions, you must verify your answer to receive full credit.
- State and apply the appropriate boundary conditions for the potential interface at  $x = 0$ . Obtain expressions relating the different coefficients.
- Using the results from part b), determine the transmission coefficient  $T$  for a particle incident and transmitted through this potential barrier. Discuss the limiting cases for  $T$  as the energy  $E$  approaches zero ( $E \rightarrow 0$ ) and  $E$  approaches infinity ( $E \rightarrow \infty$ ).

a) For  $x > 0$ ,  $V(x) = 0$  and :

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi = E \psi$$

$$\frac{d^2 \psi}{dx^2} = -\frac{2mE}{\hbar^2} \psi$$

Let  $k = \sqrt{\frac{2mE}{\hbar^2}}$ . Then,

$$\frac{d^2 \psi}{dx^2} = -k^2 \psi \quad (1)$$

Eq. (1) has the solution

$$\psi_+(x) = c_1 e^{ikx} + c_2 e^{-ikx} \quad (2)$$

The logic for  $x < 0$  is identical

$$\psi_-(x) = d_1 e^{ikx} + d_2 e^{-ikx} \quad (3)$$

b) We require continuity

$$\psi_+(0) = \psi_-(0) \quad (4)$$

Inserting Eq. (2) and (3) into (4):

$$c_1 + c_2 = d_1 + d_2 \quad (6)$$

Now,

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + \alpha \delta(x) \psi = E \psi$$

Integrate both sides from  $-\epsilon$  to  $\epsilon$

$$-\frac{\hbar^2}{2m} \int_{-\epsilon}^{\epsilon} dx \frac{d^2 \psi}{dx^2} + \alpha \int_{-\epsilon}^{\epsilon} \delta(x) \psi(x) = E \int_{-\epsilon}^{\epsilon} dx \psi(x)$$

$$-\frac{\hbar^2}{2m} \left[ \frac{d\psi}{dx} \right]_{-\epsilon}^{\epsilon} + \alpha \psi(0) = E \int_{-\epsilon}^{\epsilon} dx \psi(x)$$

Letting  $\epsilon \rightarrow 0$ ,

$$-\frac{\hbar^2}{2m} \left[ \frac{d\psi_+}{dx} \Big|_{x=0} - \frac{d\psi_-}{dx} \Big|_{x=0} \right] + q\psi(0) = 0$$

$$\frac{d\psi_+}{dx} \Big|_{x=0} - \frac{d\psi_-}{dx} \Big|_{x=0} = \frac{2m\alpha}{\hbar^2} \psi(0)$$

From Eq. (2) and (3) :

$$i\hbar(c_1 - c_2) - i\hbar(d_1 - d_2) = \frac{2m\alpha}{\hbar^2} (c_1 + c_2) \quad (7)$$

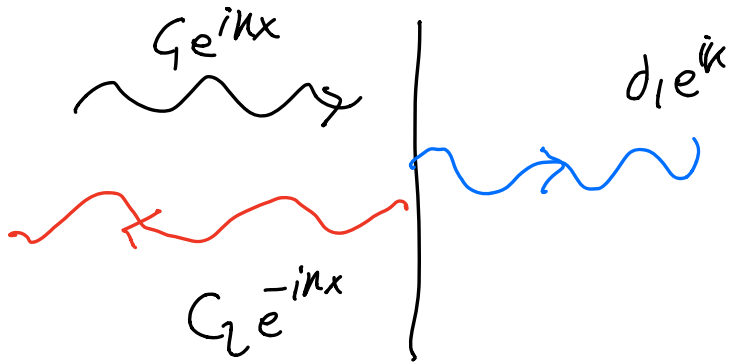
Note : Since  $\psi_+(0) = \psi_-(0)$  we could have used  $d_1 + d_2$  on the RHS.

C) Let's take the incident particle moving in the positive  $x$  direction from  $-\infty$ .  
Then  $d_2 = 0$  and Eq. (6) and (7) become:

$$c_1 + c_2 = d_2 \quad (8)$$

$$i\hbar(c_1 - c_2) + i\hbar d_1 = \frac{2m\alpha}{\hbar^2} (c_1 + c_2) \quad (9)$$

Since  $T = \left| \frac{d_1}{c_1} \right|^2$  then we use  $c_2 = d_2 - c_1$



$$ik(c_1 + c_2 - d_1) - ikd_1 = \frac{2m\alpha}{\hbar^2} (c_1 + d_1 - c_2)$$

$$2ikc_1 - 2ikd_1 = \frac{2m\alpha}{\hbar^2} d_1$$

$$\left(-\frac{m\alpha}{\hbar^2} = ik\right) d_1 = -ikc_1$$

$$\frac{d_1}{c_1} = \frac{ik}{\frac{m\alpha}{\hbar^2} + ik} \quad \begin{matrix} 1/ik \\ 1/ik \end{matrix}$$

$$\frac{d_1}{c_1} = \frac{1}{1 - i \frac{m\alpha}{\hbar^2 k}}$$

$$\Rightarrow T = \left| \frac{d_1}{c_1} \right|^2 = \frac{1}{1 + \frac{m^2 \alpha^2}{\hbar^4 k^2}}$$

Using  $k^2 = \frac{2mE}{\hbar^2}$ ,

$$T = \frac{1}{1 + \frac{m^2 \alpha^2}{\hbar^4 \left( \frac{2mE}{\hbar^2} \right)}}$$

$$T = \frac{1}{1 + \frac{m \alpha^2}{2\hbar^2 E}}$$

(10)

In Eq. (10) as  $E \rightarrow 0$   $T \rightarrow 0$ .  
With little energy probability of getting past  
barrier is small.

As  $E \rightarrow \infty$ ,  $T \rightarrow 1$ . Lots  
of energy guarantees particle will make  
it past the delta barrier.