Phase field modeling of ferrite growth in Fe-C-Mn steels

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Motivation

Martensite banding in dual phase steels

[Grange, Metallurgical Transactions 2 (1971) 417]



[B. Krebs, A. Hazotte, M. Gouné]

- Mn Microsegregation
- After thermomechanical process of dual phase steels, martensite banding is likely to appear
- This morphology is detrimental to mechanical properties

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• Role of ferrite *growth* on the banding or is *nucleation* sufficient to explain everything ?

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- Mn Microsegregation
- After thermomechanical process of dual phase steels, martensite banding is likely to appear
- This morphology is detrimental to mechanical properties

- Phase field could be well suited for that purpose
- ... provided that kinetics is properly predicted

Motivation

Different Growth Modes ($D_{Mn} \ll D_C$)

[Coates, Metallurgical Transactions 3 (1972) 1203]

Orthoequilibrium

- Local Equilibrium with Partitioning (LEP)
- Local Equilibrium with Non Partitioning (LENP)
- Paraequilibrium
- Transitions between the different regimes

Is that straightforward to account for with phase field ?

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Mixing *Kim, Kim, Suzuki* and *Warren, Kobayashi, Craig Carter* [Kim et al., Phys. Rev. E 60 (1999) 7186], [Warren et al., Acta Mater. 51 (2003) 6035]

$$\mathcal{F} = \int_{\Omega} \left(f_0(\phi, \mathbf{c}_i^{\alpha}, \mathbf{c}_i^{\gamma}) + \frac{\epsilon^2}{2} |\nabla \phi|^2 + \chi \left[1 - h(\phi)\right] |\nabla \psi| \right) \mathrm{d}\Omega$$

Homogeneous free energy

 $f_0(\phi, \boldsymbol{c}_i^{\alpha}, \boldsymbol{c}_i^{\gamma}) = h(\phi) f^{\gamma}(\boldsymbol{c}_i^{\gamma}) + [1 - h(\phi)] f^{\alpha}(\boldsymbol{c}_i^{\alpha}) + W g(\phi)$

- Arbitrary $f^{\gamma}(c_i^{\gamma})$ and $f^{\alpha}(c_i^{\alpha})$ (e.g. Calphad)
- No extra contribution of the concentration profile within the interface to the interfacial energy thanks to auxilliary c^φ_i
- Thin interface asymptotics analysis for determining M_{ϕ}

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Governing equations

$$\mathcal{F} = \int_{\Omega} \left(f_0(\phi, \boldsymbol{c}_i^{\alpha}, \boldsymbol{c}_j^{\gamma}) + \frac{\epsilon^2}{2} |\nabla \phi|^2 \right) \, \mathrm{d}\Omega$$

$$\partial_t \phi = -M_\phi \left(\frac{\delta \mathcal{F}}{\delta \phi} \right)$$

 $\partial_t c_i = \nabla \cdot \left(L_{ij} \nabla \frac{\delta \mathcal{F}}{\delta c_j} \right)$

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ight) \end{aligned}$$

supplied with:

$$c_i = h(\phi) c_i^{\alpha} + [1 - h(\phi)] c_i^{\gamma}$$

$$\frac{\partial f^{\alpha}}{\partial \boldsymbol{c}_{i}^{\alpha}} = \frac{\partial f^{\gamma}}{\partial \boldsymbol{c}_{i}^{\gamma}}$$

similar to homogenization schemes

[Ammar et al., Eur. J. Comp. Mech. (2010)]

Talk K. Ammar (thursday, Benoît XII)

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AIEMI 2010 - Avignon

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Linearized phase diagram



Asymptotics analysis (over-simplified) Poster Y. Le Bouar

- local frame
- expansions of the fields $(\phi, c_i^{\phi} \dots)$
- Order 0 (conventional analysis)

$$\phi(x) = rac{1}{2} (1 - \tanh \xi)$$
 $c^{lpha(0)} = c^{lpha}_{
m eq}$
 $c^{\gamma(0)} = c^{\gamma}_{
m eq}$



Order 1

$$\frac{1}{M_{\phi}} = \frac{\epsilon^{3}}{\sigma\sqrt{2W}} \left\{ \frac{1}{D_{\mathrm{C}}^{\gamma}} \zeta_{\mathrm{C}} + \frac{1}{D_{\mathrm{Mn}}^{\gamma}} \zeta_{\mathrm{Mn}} \right\}$$

Thuillier's experiments

[Thuillier et al., Scripta Materialia 55 (2006) 1071]

• Fe – 0.08%C – 1.7%Mn; γ grain size = 20 μ m; $T = 700^{\circ}$ C



- Measured kinetics is faster than orthoequilibrium (NPLE)
- Measured kinetics stops below paraequilibrium fraction
- Good agreement between phase field and experiments

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Nature of local equilibrium at the interface in the growth of ferrite from alloyed austenite [M. Hillert, Scripta Mater. 46 (2002) 447]

Thuillier's experiments

[Thuillier et al., Scripta Materialia 55 (2006) 1071]

• Fe – 0.08%C – 1.7%Mn; γ grain size = 20 μ m; $T = 700^{\circ}$ C



Influence of the interface thickness

- Weak for the first growth regime intermediate between LENP and paraequilibium (at least in 1D)
- Strong for the transition to LENP

Analysis of the spike

Recasting the diffusion equation in a steady moving frame

$$-\mathbf{v} \,\mathrm{d}_{\mathbf{x}} \mathbf{c} = \mathrm{d}_{\mathbf{x}} ig(L \,\mathrm{d}_{\mathbf{x}} \partial_{\gamma} f_{\gamma} ig)$$

$$c^{\gamma}(r) = \exp(-A(r)) \left[\frac{c^{\alpha*}}{k} \exp(A(-\infty)) + \int_{-\infty}^{r} \operatorname{Pe} B(u) c^{\alpha*} \exp(A(u)) du \right]$$

where $r = 2x/\delta$, $\operatorname{Pe} = v\delta/(2D^{\alpha})$

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Analysis of the spike



Analysis of the spike



Analysis of the spike



Analysis of the spike



What can we conclude ?

- This is a generic behaviour (whatever δ) depending on Pe for a given interpolation of $L \partial_{\gamma\gamma} f_{\gamma}$ (the B(u) effect)
- The position z_{α} of the *plateau* is set by a critical value Pe_t $\Longrightarrow z_{\alpha} = K^2 \, \delta / (4D^{\alpha} \operatorname{Pe}_t)$
- large δ mimics slugish trans-interface Mn diffusion
 - Realistic transition for $\delta = 44 \text{ nm}$
 - Realistic transition for $\delta = 12 \text{ nm}$ in Fe-C-Ni

Consistent with [A. Zurob et al., Acta Mater. 56 (2008) 2203]

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Interface versus grain size



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Conclusion

9h33 in TGV this morning

- We must try to extract some interface diffusivities from our calculations
- We were not able to reproduce the trends observed in the decarburization experiments
- I'd like to play with strange interpolations of the *B*(*u*) within the interface for that purpose
- and to play with cross terms in the Onsager mobility matrix

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