



12th ALEMI workshop

Phase-field model for mixed mode transformations and its perspective for steel

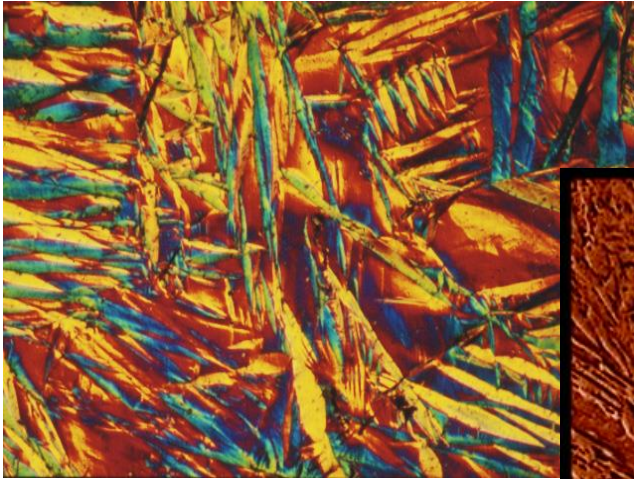
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Interdisciplinary Centre for Advanced Materials Simulation (ICAMS),
Ruhr-Universität Bochum, Germany

Delft, 25th of June 2013

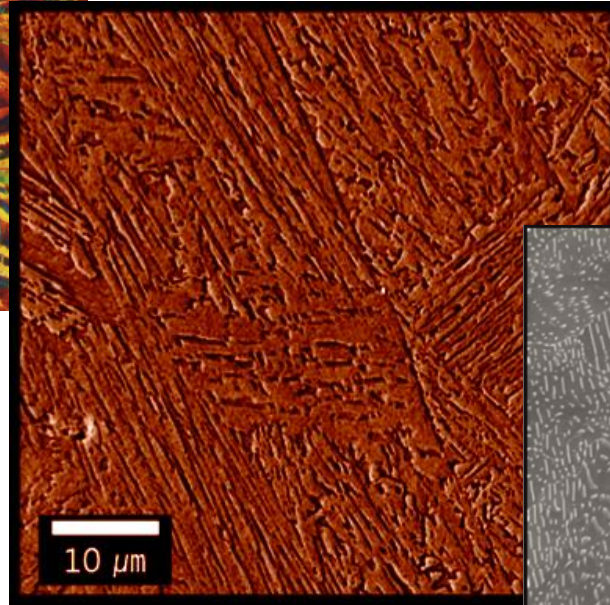
Steel Microstructures

Martensite



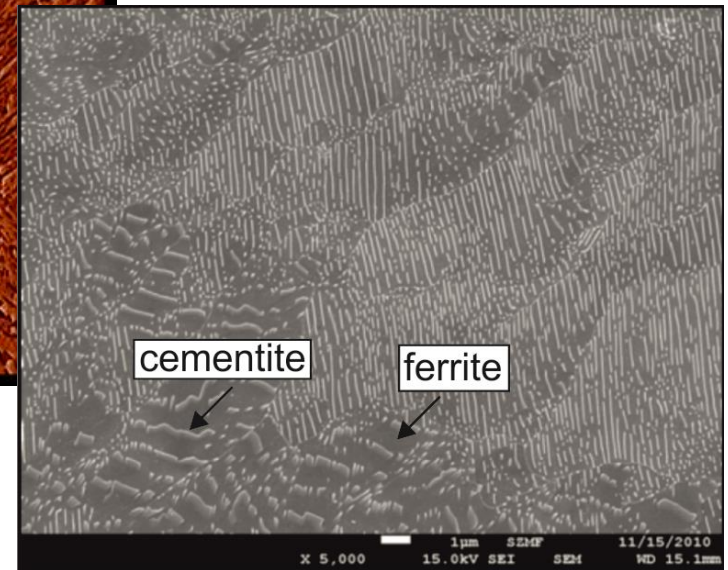
from lecture material
of H.K.D.H.Bhadeshia

Bainite



Junhak Pak, Dong Woo Suh,
and H. K. D. H. Bhadeshia,
Scripta Materialia 66 (2012)

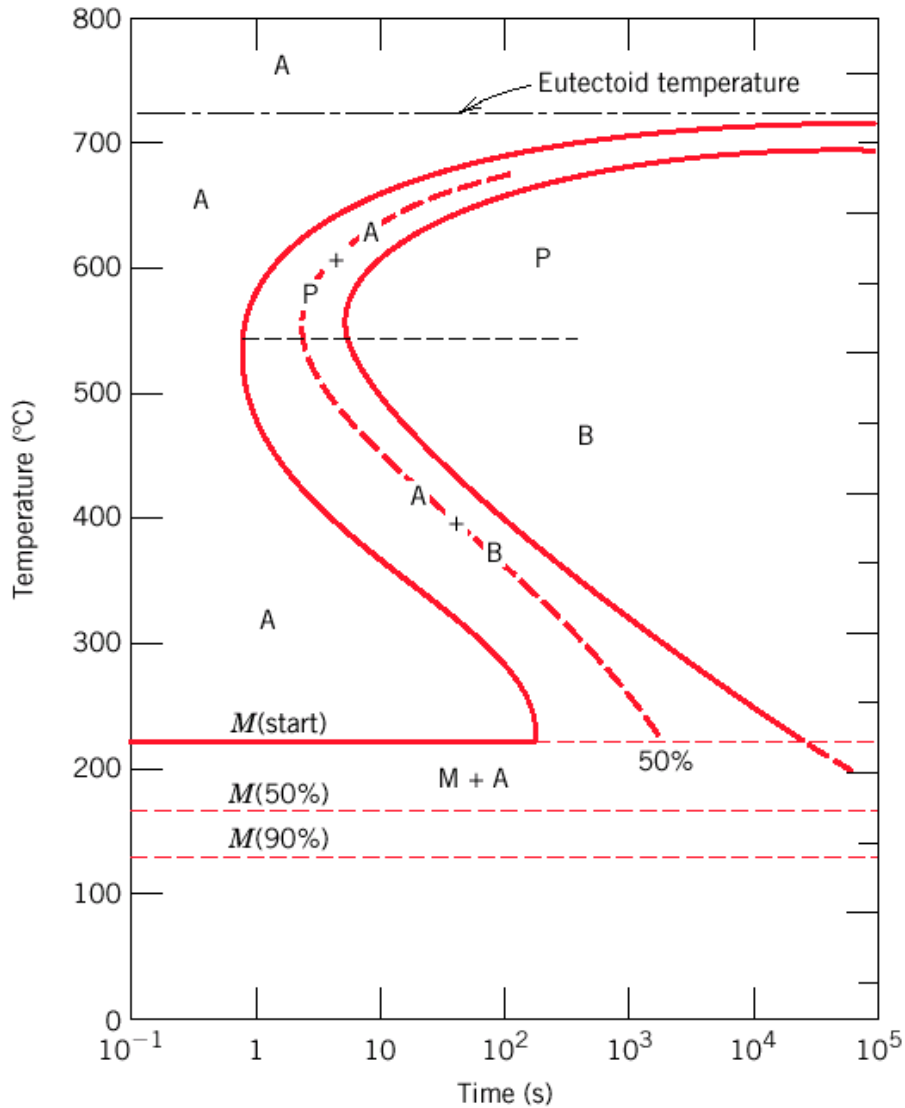
Pearlite



28Mn6 steel Benteler Stahl/Rohr “as-cast”
microstructure

Steel forms many phases and microstructures
and every microstructure possesses different properties!

Fe-C TTT Diagram



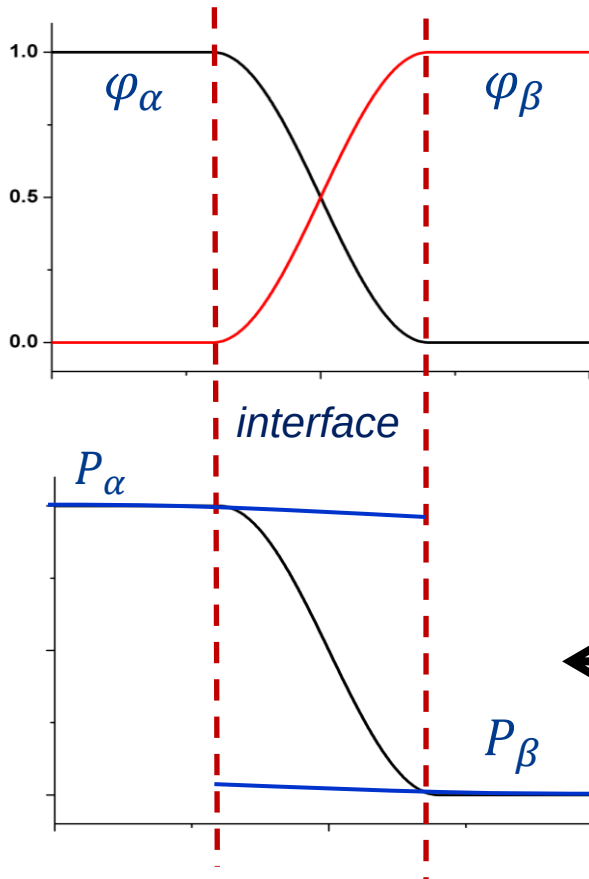
The key to steel properties are highly non-equilibrium phase transformations (e.g. martensitic, bainitic, ...)

Outline

1. Standard multi phase field model
2. Phase field model for non-equilibrium transformations
3. Modeling phase transformations in steel
4. New developments
5. Summary

Phase field, properties and free energy

$$F = \int_{\Omega} \overset{\text{interfacial}}{f^{GB}} + \overset{\text{chemical}}{f^{CH}} + \overset{\text{elastic}}{f^{EL}} + \overset{\text{disl. dens.}}{f^{DD}} + \dots$$



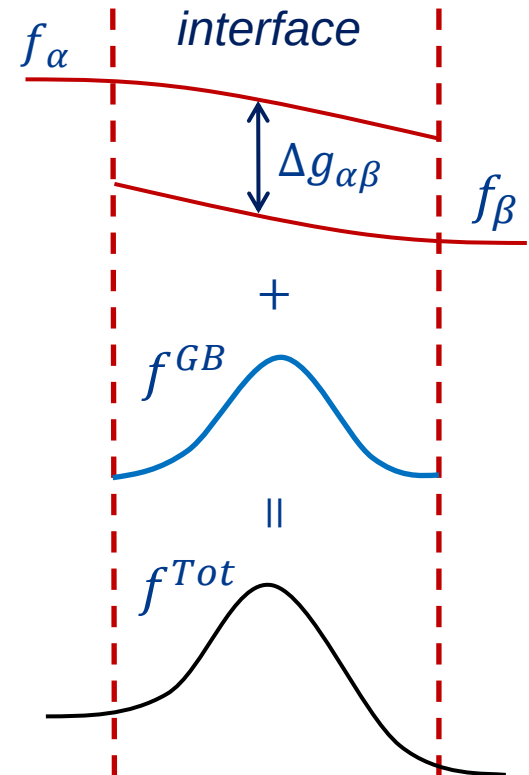
Phase field

$$\varphi_\alpha + \varphi_\beta \equiv 1$$

Properties

$$P = \varphi_\alpha * P_\alpha + \varphi_\beta * P_\beta$$

Energy →



Free energy densities and evolution equations

Interfacial energy density

$$f^{GB} = \sum_{\alpha, \beta=1}^N \frac{4\sigma_{\alpha\beta}}{\eta_{\alpha\beta}} \left\{ \frac{\eta_{\alpha\beta}^2}{\pi^2} |\nabla \phi_\alpha| \cdot |\nabla \phi_\beta| + \phi_\alpha \phi_\beta \right\}$$

Phase field equation

$$\dot{\phi}_\alpha + \vec{u}_\alpha \cdot \nabla \phi_\alpha = - \sum_{\beta=1}^N \frac{\mu_{\alpha\beta}}{N} \left(\frac{\delta F}{\delta \phi_\alpha} - \frac{\delta F}{\delta \phi_\beta} \right)$$

Chemical energy density

$$f^{CH} = \sum_{\alpha=1}^N \phi_\alpha f_\alpha(\bar{c}_\alpha) + \bar{\mu} \left(\bar{c} - \sum_{\alpha=1}^N \phi_\alpha \bar{c}_\alpha \right)$$

Advection-diffusion equation

$$\dot{\bar{c}} + \sum_{\alpha=1}^N \vec{u}_\alpha \cdot \nabla (\phi_\alpha \bar{c}_\alpha) = \nabla \cdot \left(\sum_{\alpha=1}^N M_\alpha \phi_\alpha \nabla \frac{\delta F}{\delta \bar{c}_\alpha} \right) + \nabla \cdot \vec{j}_{at}$$

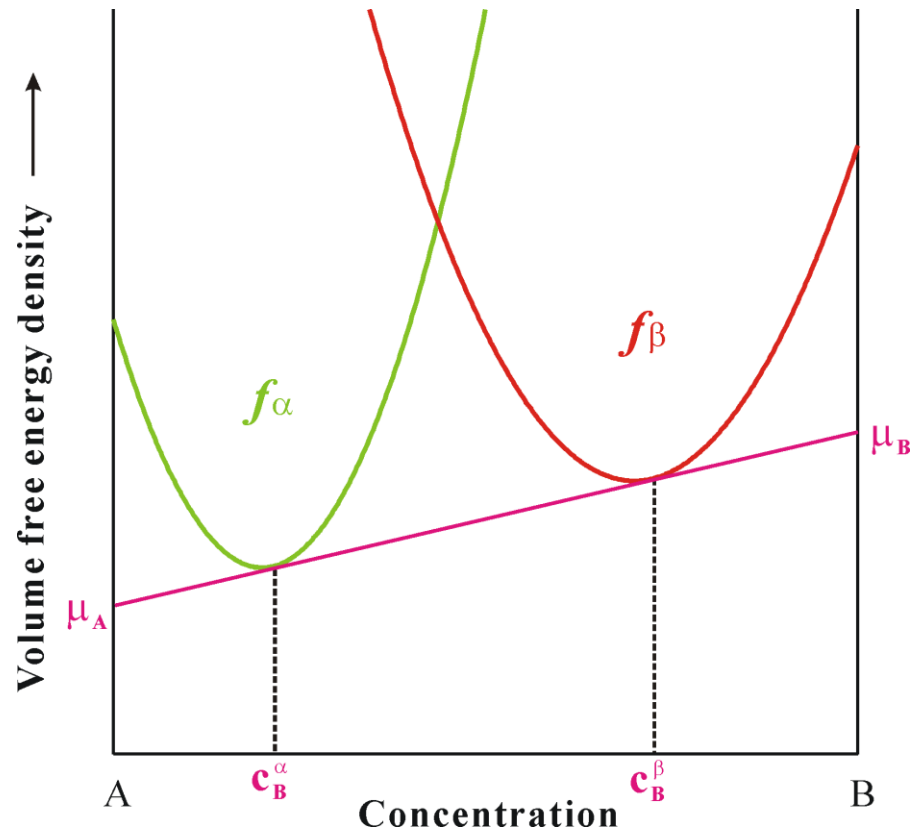
Elastic energy density

$$f^{EL} = \frac{1}{2} \sum_{\alpha=1}^N \phi_\alpha \left(\varepsilon_\alpha^{ij} - \varepsilon_\alpha^{*ij} \right) C_\alpha^{ijkl} \left(\varepsilon_\alpha^{kl} - \varepsilon_\alpha^{*kl} \right)$$

Elastic equilibrium equation

$$0^i = \nabla^j \sigma^{ij} = \nabla^j \frac{\delta F}{\delta \varepsilon^{ij}};$$

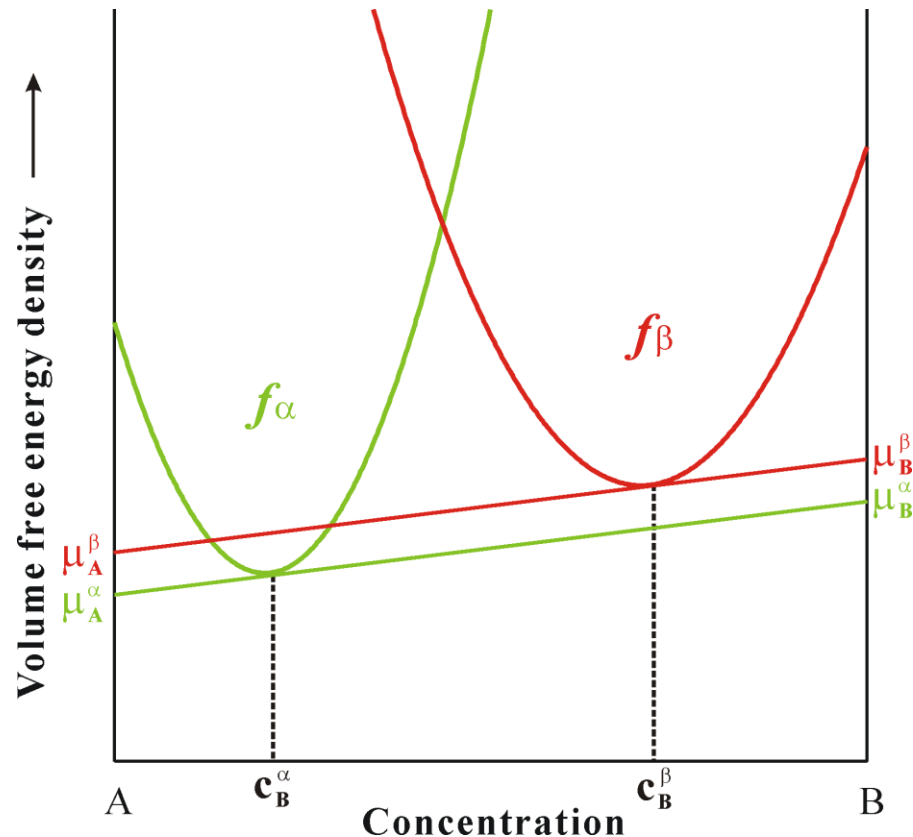
Current approaches: Local equilibrium



- Double tangent constraint
- Equilibrium concentrations
- Equal chemical potential



Current approaches: Quasi equilibrium



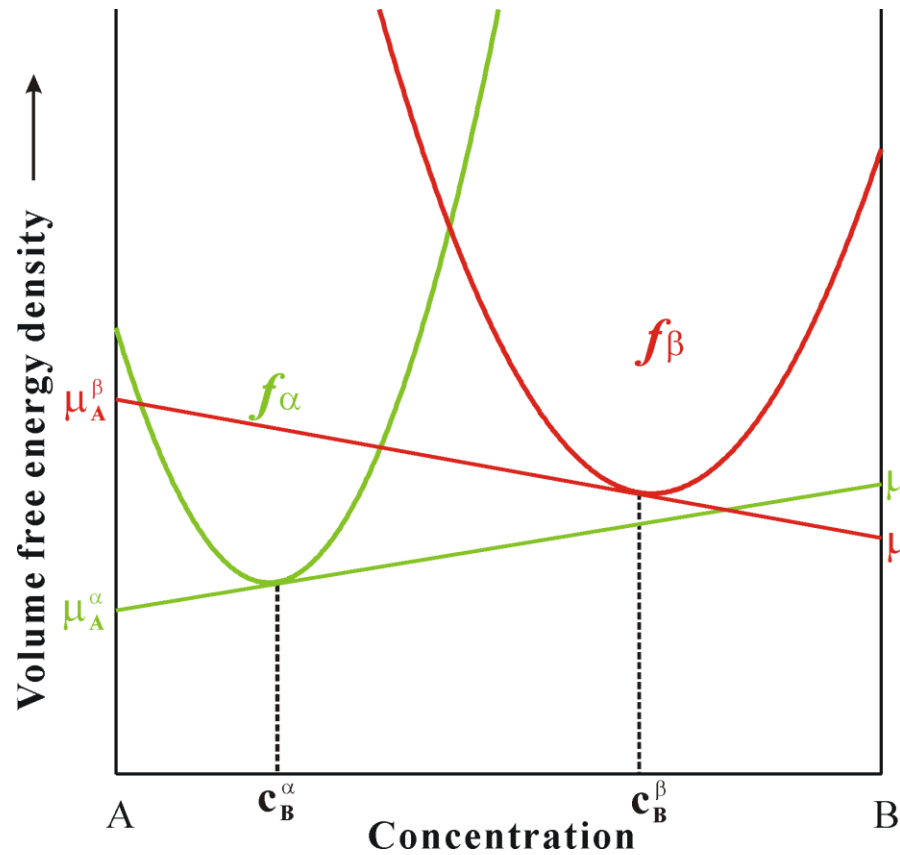
- Parallel tangent constraint
- Equal diffusion potential

$$\mu_\alpha = \frac{\partial f_\alpha(T, c_\alpha)}{\partial c_\alpha} = \mu_\beta = \frac{\partial f_\beta(T, c_\beta)}{\partial c_\beta} \equiv \mu$$

$$c(x, t) = c_\alpha(x, t)\phi(x, t) + c_\beta(x, t)[1 - \phi(x, t)]$$

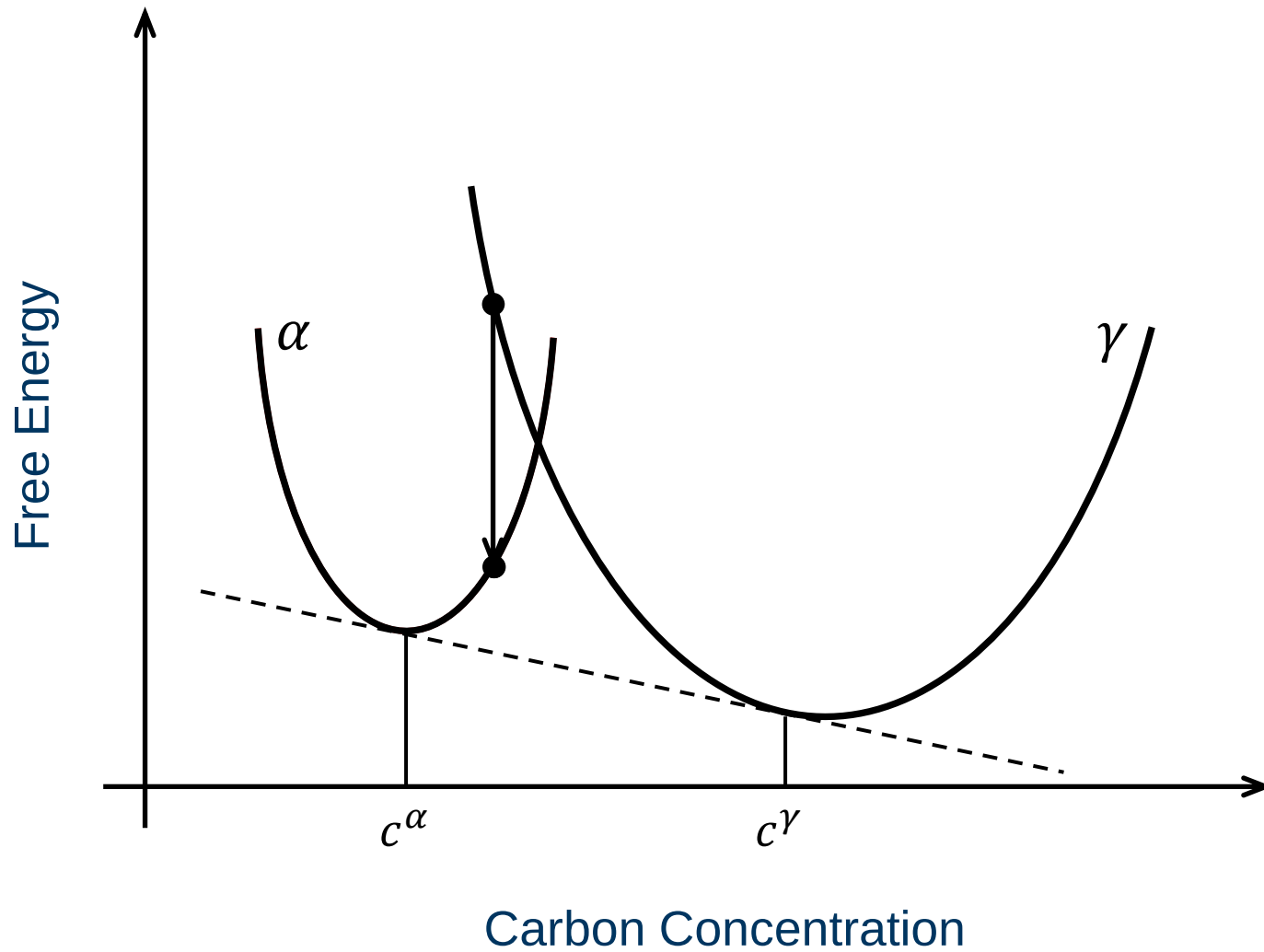
$$\dot{c} = \nabla D_\alpha \phi \nabla c_\alpha + \nabla D_\beta (1 - \phi) \nabla c_\beta$$

The general case



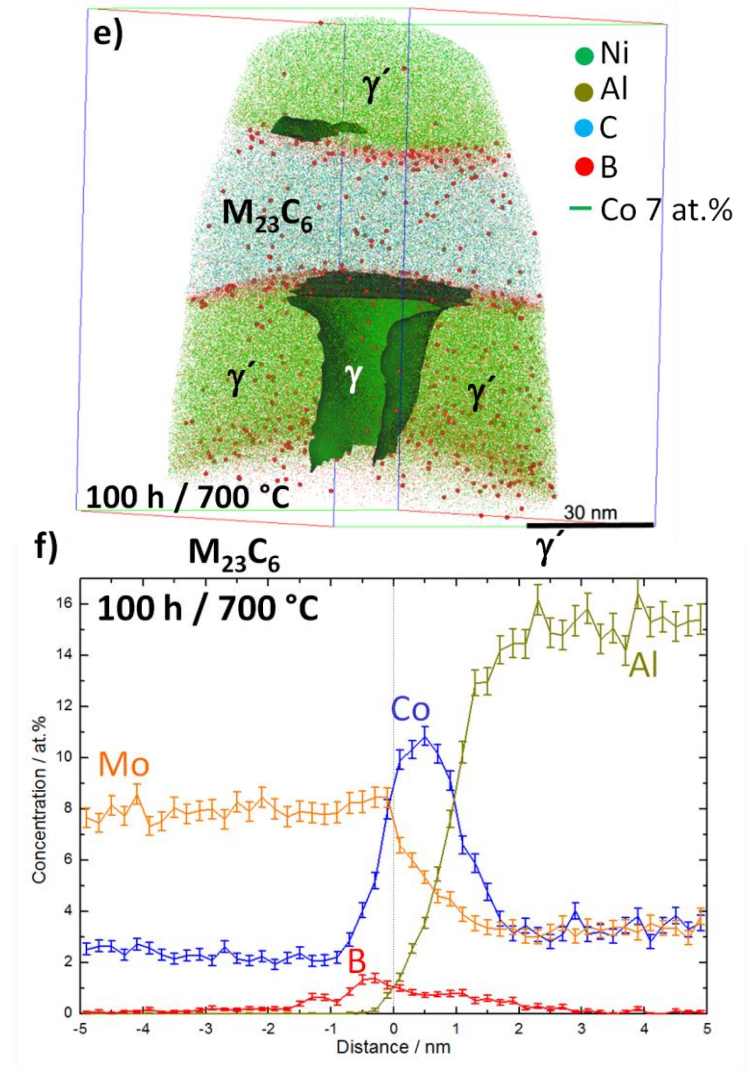
$$\mu_\alpha = \frac{\partial f_\alpha}{\partial c_\alpha} \neq \mu_\beta = \frac{\partial f_\beta}{\partial c_\beta} \neq \mu$$

Effect of stress



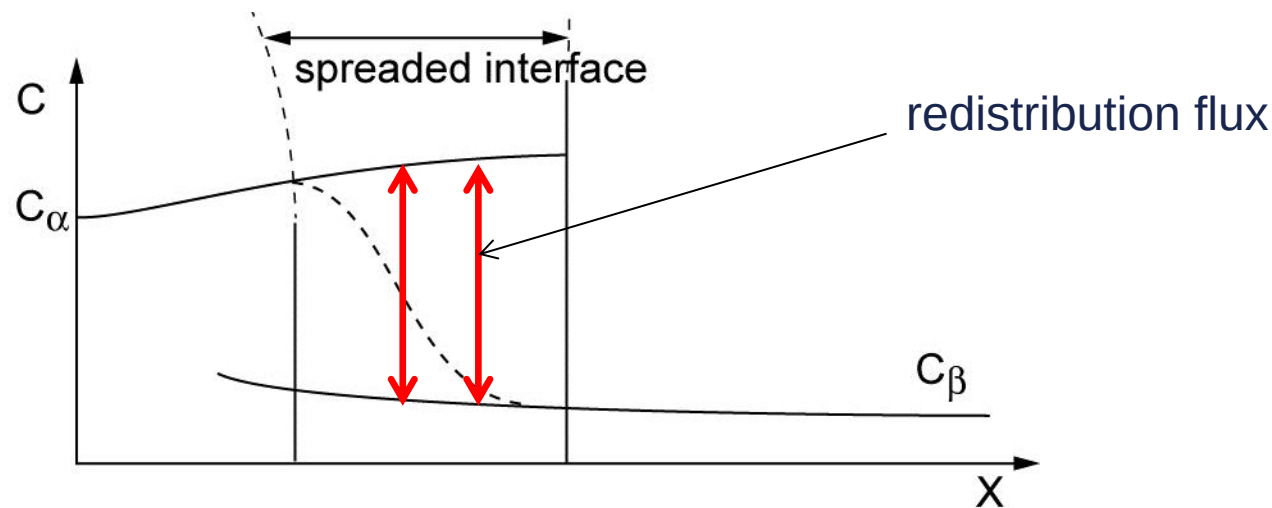
Coarse graining of the diffusion equation

$$\dot{c} = \vec{\nabla} \left(M(c, \phi) \frac{\partial^2 f}{\partial c^2} \vec{\nabla} c \right)$$



D. Tytko et. al, Acta Mat 2012

Coarse graining of the diffusion equation



$$\dot{c}_\alpha(\vec{x}) = -4\pi^2 M \int_{|\vec{k}| < k_0} k^2 \mu(\vec{k}) e^{-2\pi i \vec{k} \vec{x}} d\vec{k} + P(\mu_\beta - \mu_\alpha); \quad \vec{x} \in \Omega_\alpha$$

$$\dot{c}_\beta(\vec{x}) = -4\pi^2 M \int_{|\vec{k}| < k_0} k^2 \mu(\vec{k}) e^{-2\pi i \vec{k} \vec{x}} d\vec{k} + P(\mu_\alpha - \mu_\beta); \quad \vec{x} \in \Omega_\beta.$$

$$P = -\frac{1}{(\mu_\alpha - \mu_\beta)} 4\pi M \int_{|\vec{k}| < k_0} k^2 \mu(\vec{k}) e^{-2\pi i \vec{k} \vec{x}} d\vec{k}.$$

Diffusion equation in one phase

$$\phi_{\alpha} \dot{c}_{\alpha} = \vec{\nabla} \left(\phi_{\alpha} D_{\alpha} \vec{\nabla} c_{\alpha} \right) + P_{\alpha\beta} \phi_{\alpha} \phi_{\beta} (\tilde{\mu}_{\beta} - \tilde{\mu}_{\alpha}) + \phi_{\alpha} \dot{\phi}_{\alpha} (c_{\alpha} - c_{\beta})$$

diffusion within one phase

redistribute solute if there is a
chemical potential difference
between the phases

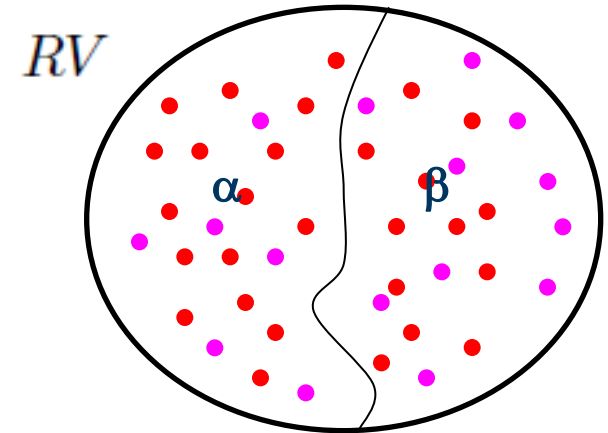
phase concentrations change
if there is a phase change

*Steinbach, Plapp, Zhang. Acta Mat 2012;
Zhang, Steinbach. Acta Mat 2012*

Redistribution within one reference volume (RV)

- No flux over the boundary
- Phase fractions are fixed inside the RV

P: inverse interface resistivity

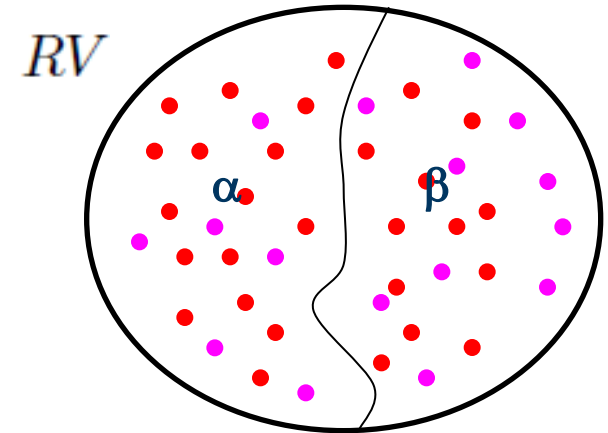


$$\phi_{\alpha} \dot{c}_{\alpha} = \underbrace{\vec{\nabla} \left(\phi_{\alpha} D_{\alpha} \vec{\nabla} c_{\alpha} \right)}_{\substack{|| \\ 0}} + \underbrace{P_{\alpha\beta} \phi_{\alpha} \phi_{\beta} (\tilde{\mu}_{\beta} - \tilde{\mu}_{\alpha})}_{\text{red oval}} + \underbrace{\phi_{\alpha} \dot{\phi}_{\alpha} (c_{\alpha} - c_{\beta})}_{\substack{|| \\ 0}}$$

Redistribution within one reference volume (RV)

- No flux over the boundary
- Phase fractions can adjust with finite rate

P: inverse interface resistivity



$$\phi_{\alpha} \dot{c}_{\alpha} = \vec{\nabla} \left(\phi_{\alpha} D_{\alpha} \vec{\nabla} c_{\alpha} \right) + P_{\alpha\beta} \phi_{\alpha} \phi_{\beta} (\tilde{\mu}_{\beta} - \tilde{\mu}_{\alpha}) + \phi_{\alpha} \dot{\phi}_{\alpha} (c_{\alpha} - c_{\beta})$$

$\parallel$
 0

The relaxation model for c and Φ

The phase-field equation is modified according to the definition of driving force and phase-field mobility K

$$\dot{\phi}_{\alpha} = K \left\{ \sigma_{\alpha\beta} \left[\nabla^2 \phi_{\alpha} + \frac{\pi^2}{\eta^2} \left(\phi_{\alpha} - \frac{1}{2} \right) \right] + \frac{\pi^2}{8\eta} [f_{\beta} - f_{\alpha} - (\phi_{\alpha} \tilde{\mu}_{\alpha} + \phi_{\beta} \tilde{\mu}_{\beta})(c_{\beta} - c_{\alpha})] \right\}$$

$$\Delta G_{\alpha\beta} = \Delta G_{\alpha\beta}^{\phi} + \Delta G_{\alpha\beta}^{Diff}$$

$$K = \frac{8P\eta\mu_{\alpha\beta}}{8P\eta + \pi^2\mu_{\alpha\beta}(c_{\beta} - c_{\alpha})^2}$$

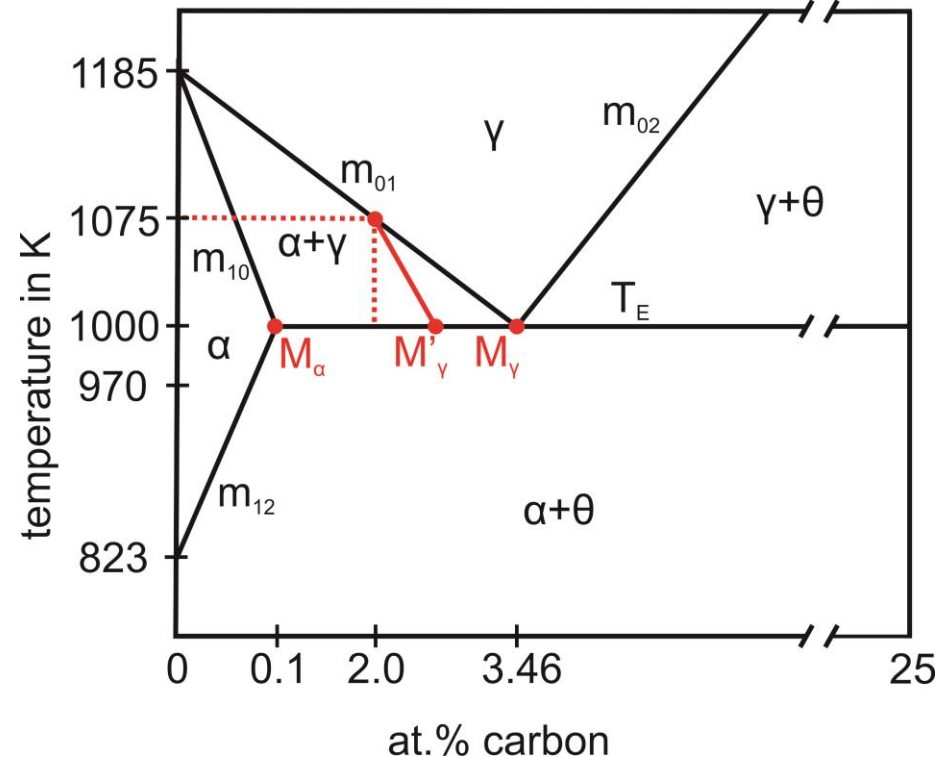
$P \rightarrow \infty$ The conventional model
(two phases with same chemical potential) $K \rightarrow \mu_{\alpha\beta}$

$P \rightarrow 0$ Transformation can not proceed $K \rightarrow 0$

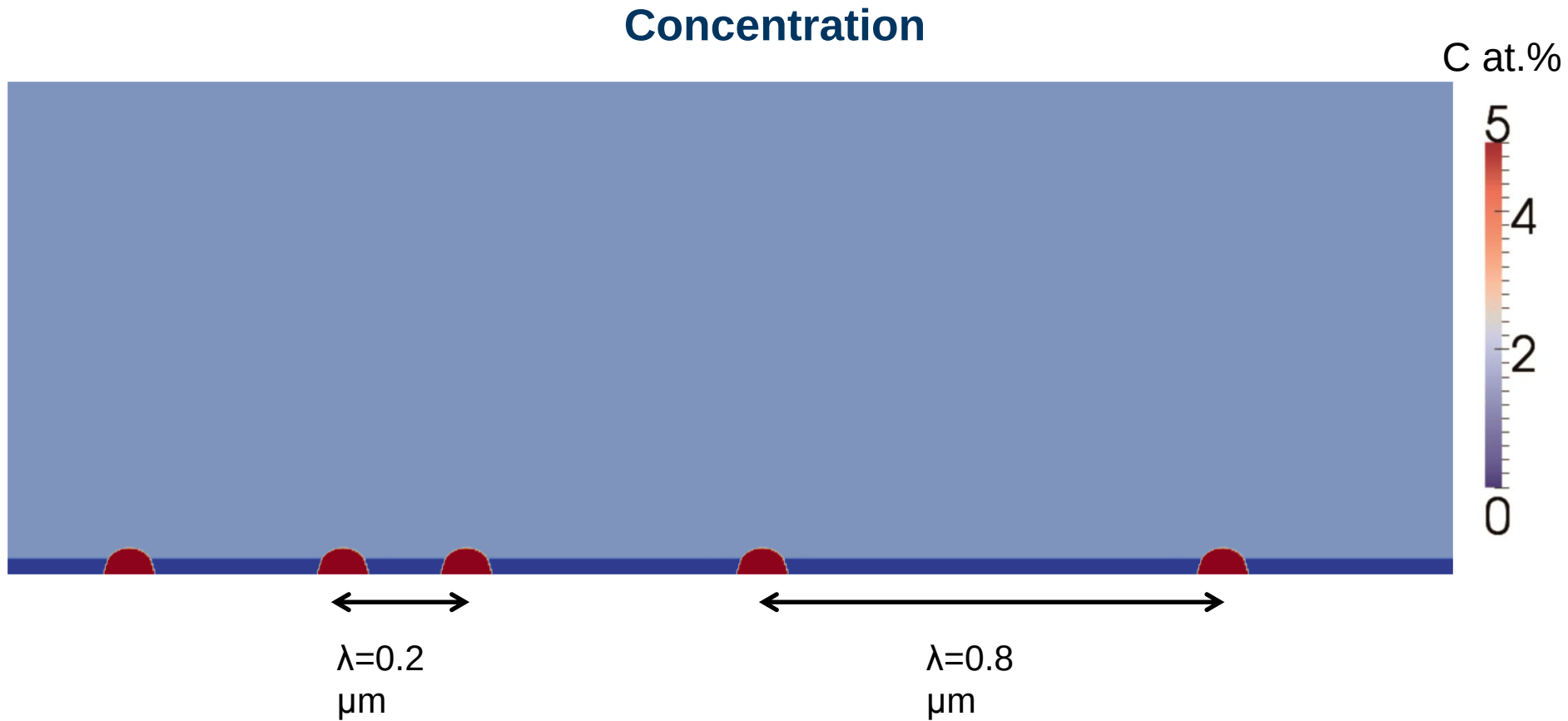
Ferrite growth simulation

Example simulation:

- steel with 2 at-% C (hypoeutectoid)
- simulation starts at 1075 K in a box of austenite and one ferrite grain
- cooling with different cooling rates to eutectoid temperature



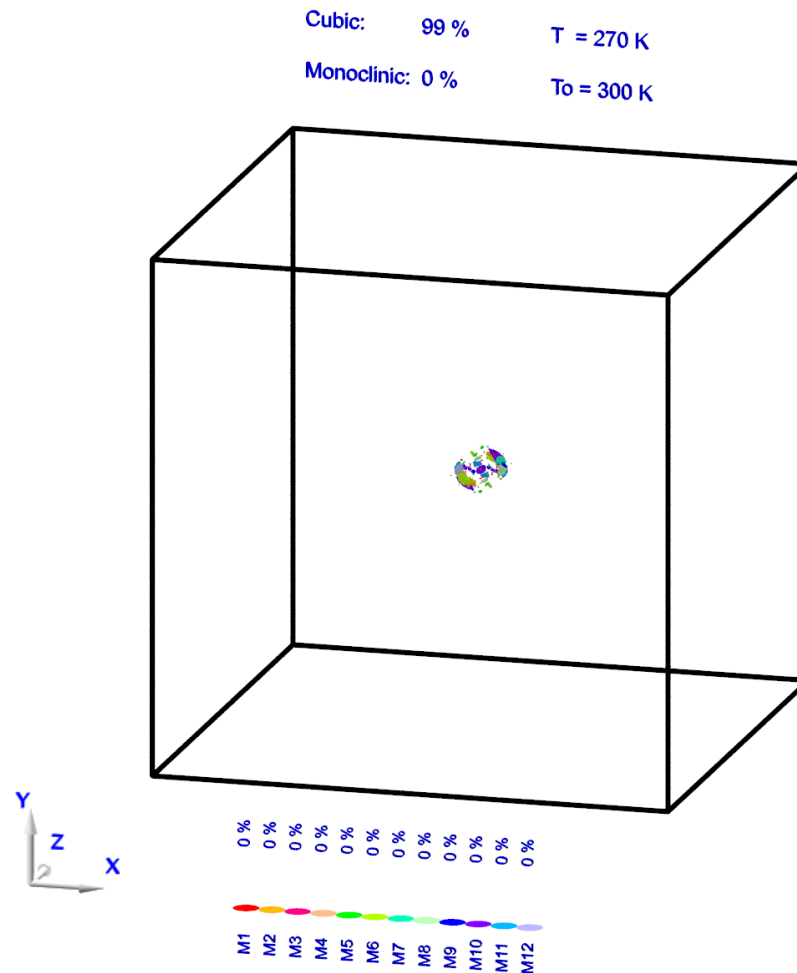
Pearlite growth simulation



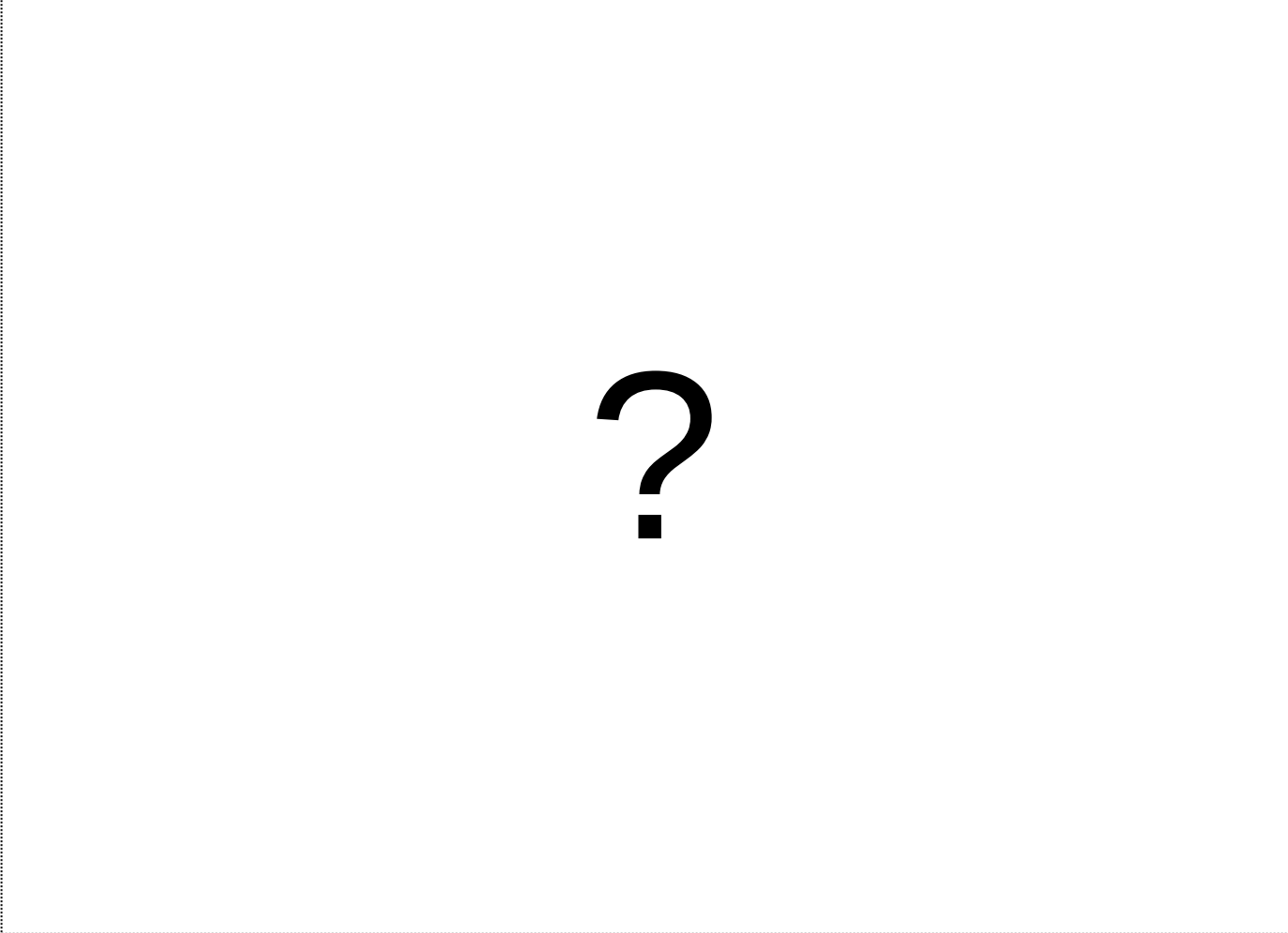
Martensite Simulation (SMA)

NiTi SMA

System size
512x512x512
grid points
(**1x1x1** μm)



Martensite Simulation (Steel)



?

Modeling transformations in steel

The necessary ingredients for successful modeling of solid-solid transformations in steel:

- **Phase field model for strongly non-equilibrium transformations**
- **Gibbs free energy dependence on local stress**

It allows diffusion and transformations modeling under stress and in presence of defects

- **Incorporation of plastic relaxation mechanism**

Is an absolute necessity for handling high strains associated with martensite and bainite formation or severe deformation during processing (recrystallization)

- **Investigation of the effect of dislocation density on equilibrium concentration of elements**

Is a necessary ingredient for simulations of realistic concentration distributions in real structures

- **Tools and algorithms development**

Any good model before being used should be implemented

Large deformations handling

Deformation decomposition:

$$f^{El} = \frac{1}{2} \underbrace{\left[\varepsilon_i^{tot} - \left(\varepsilon_i^* - \bar{\varepsilon}_i - \varepsilon_i^{rot} + \varepsilon_i^{acc} \right) \right]}_{\varepsilon_i^{elast}} C_{ij} \left[\varepsilon_j^{tot} - \left(\varepsilon_j^* - \bar{\varepsilon}_j - \varepsilon_j^{rot} + \varepsilon_j^{acc} \right) \right]$$

Diagram illustrating the decomposition of total strain into components for the elastic potential function f^{El} :

- ε_i^{tot} : Total strain
- ε_i^* : Eigenstrain
- $\bar{\varepsilon}_i$: Homogeneous expansion
- ε_i^{rot} : Rotation
- ε_i^{acc} : Accumulated strain

Local deformation accumulation:

strain accumulated
in one time step

$$\bar{\varepsilon}_{new}^{acc} = \bar{\varepsilon}_{old}^{acc} + \ln(1 + \Delta\varepsilon^{acc})$$

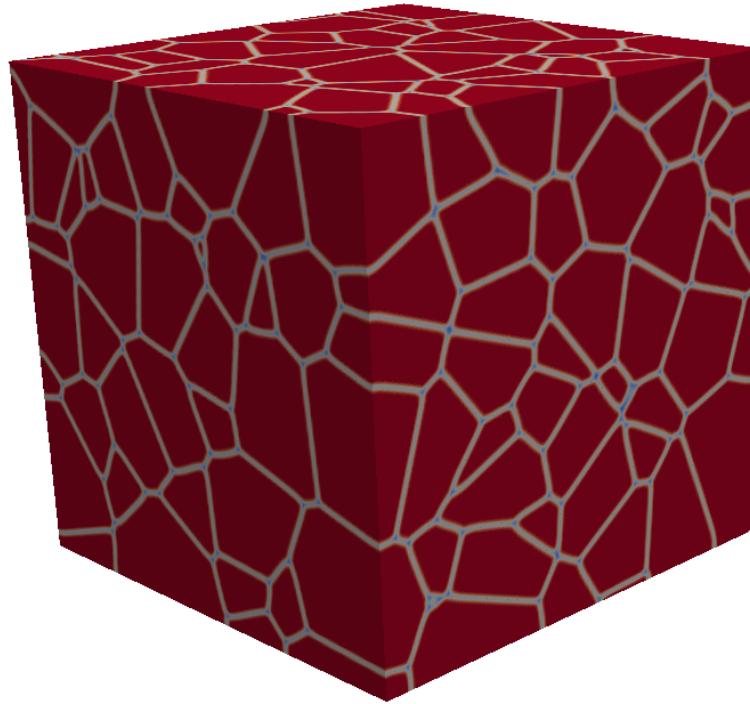
Explicit rotation of sensitive tensors and properties:

C_{ij} , ε_i^* , ε_i^{acc} , ε_i^{plast} , slip systems, anisotropic interface properties, etc.

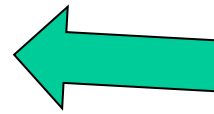
"Hot rolling"



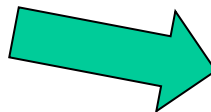
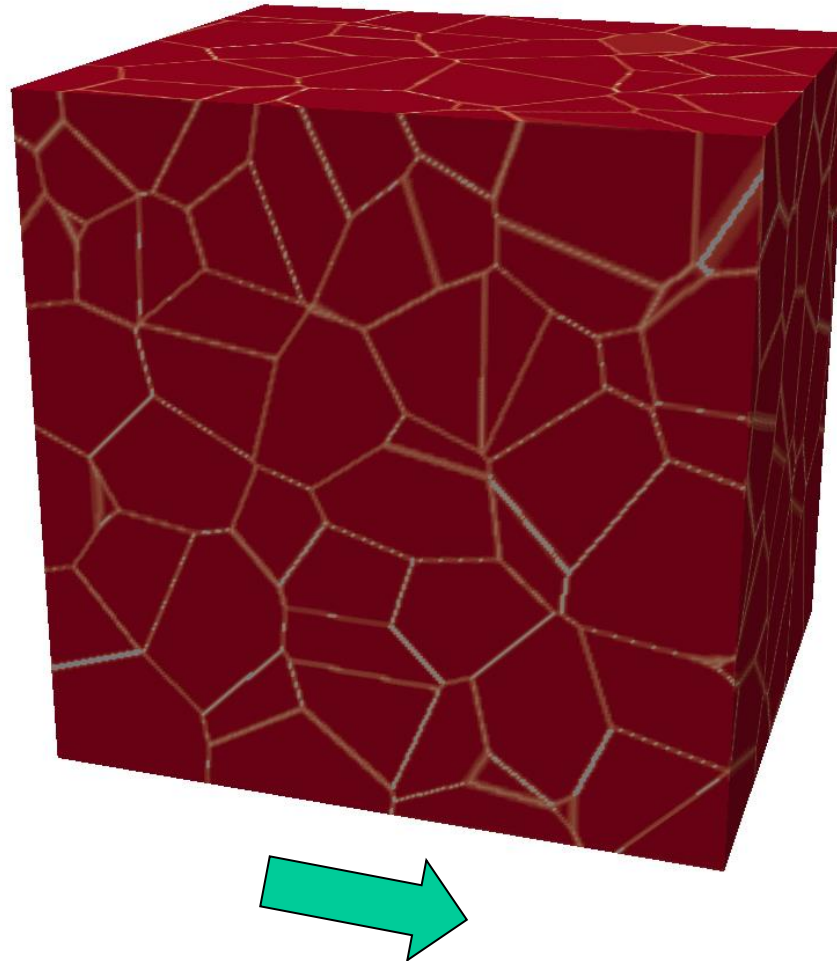
Load



"Slip"



Shear



Summary

- New phase field model suitable for simulations of highly non-equilibrium transformations is presented.
- The main features of the model are yet to be explored but it already shows great potential for simulations of transformations in steel
- In order to model steel transformation incorporation of elastic and plastic effects on diffusion, and evolution of phase and grain boundaries in one simulation is required
- There is a need for new algorithms and new models in order to improve current modeling techniques

Acknowledgements

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Efim Borukhovich	<i>Li-ion batteries, large deformation modeling</i>
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Adam Gießmann	<i>Pearlite and bainite formation</i>
Alexander Monas	<i>Light alloys solidification</i>
Guanxing Du	<i>Martensite formation in steel</i>
Philipp Engels	<i>Crystal plasticity model for phase field</i>

Postdocs:

Reza Darvishi Kamachali	<i>Grain growth</i>
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Dmitry Medvedev	<i>Fluid flow coupling, diffusion modeling</i>

www.OpenPhase.de

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Thank you for your attention!