

Fully Microscopic Description of Fission with Three Degrees of Freedom

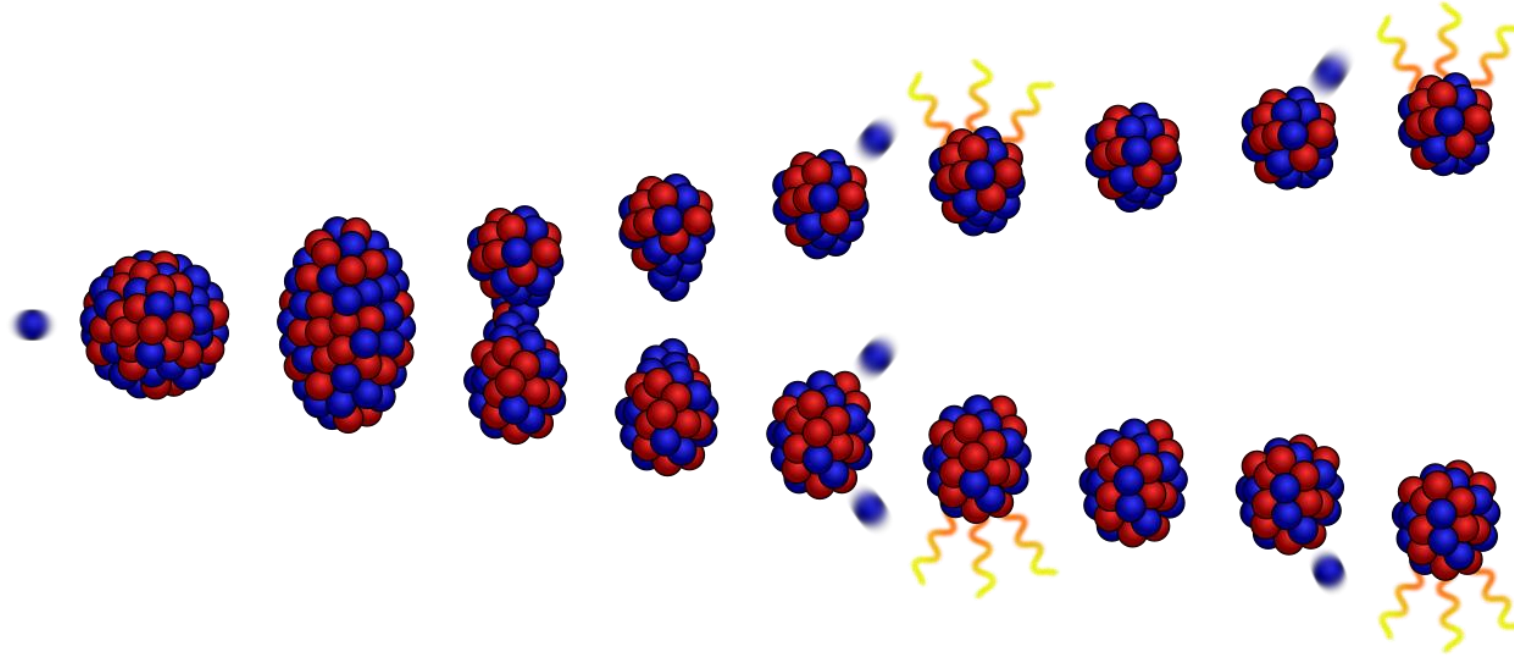
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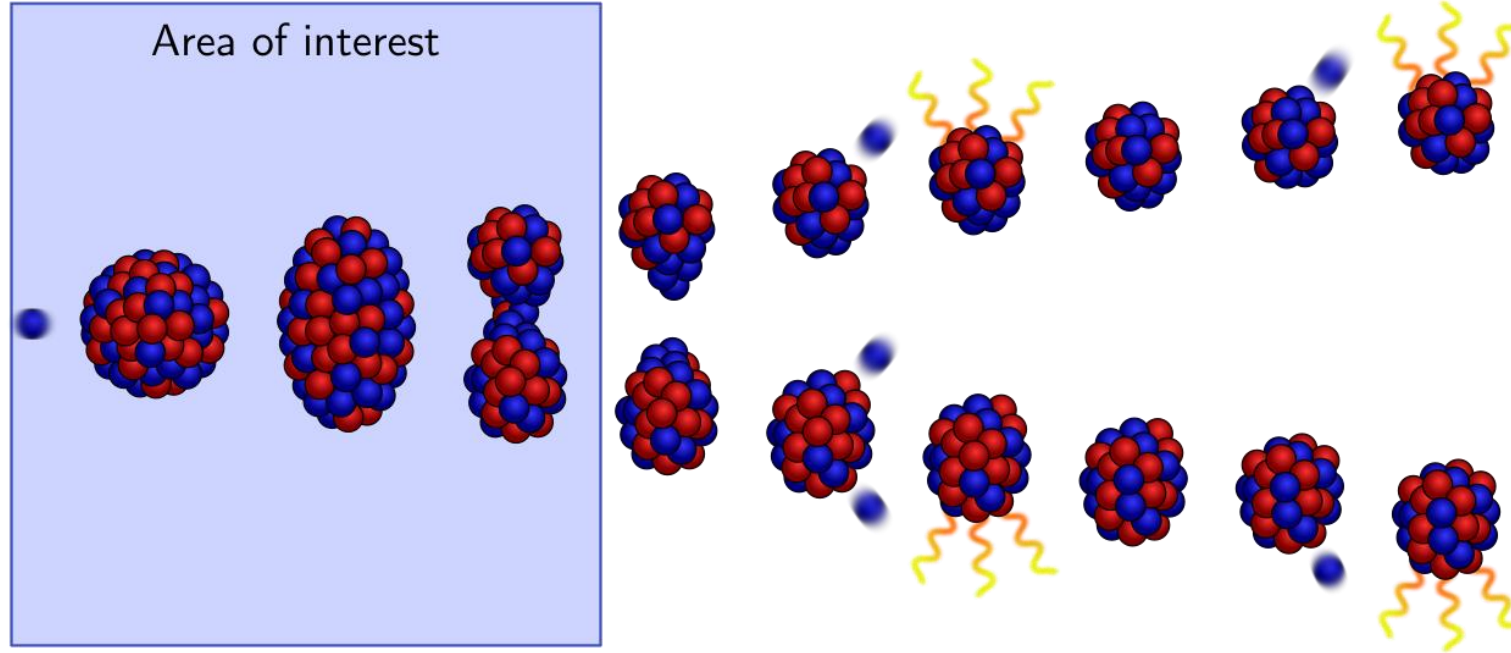
Marc Verriere, Nicolas Schunck



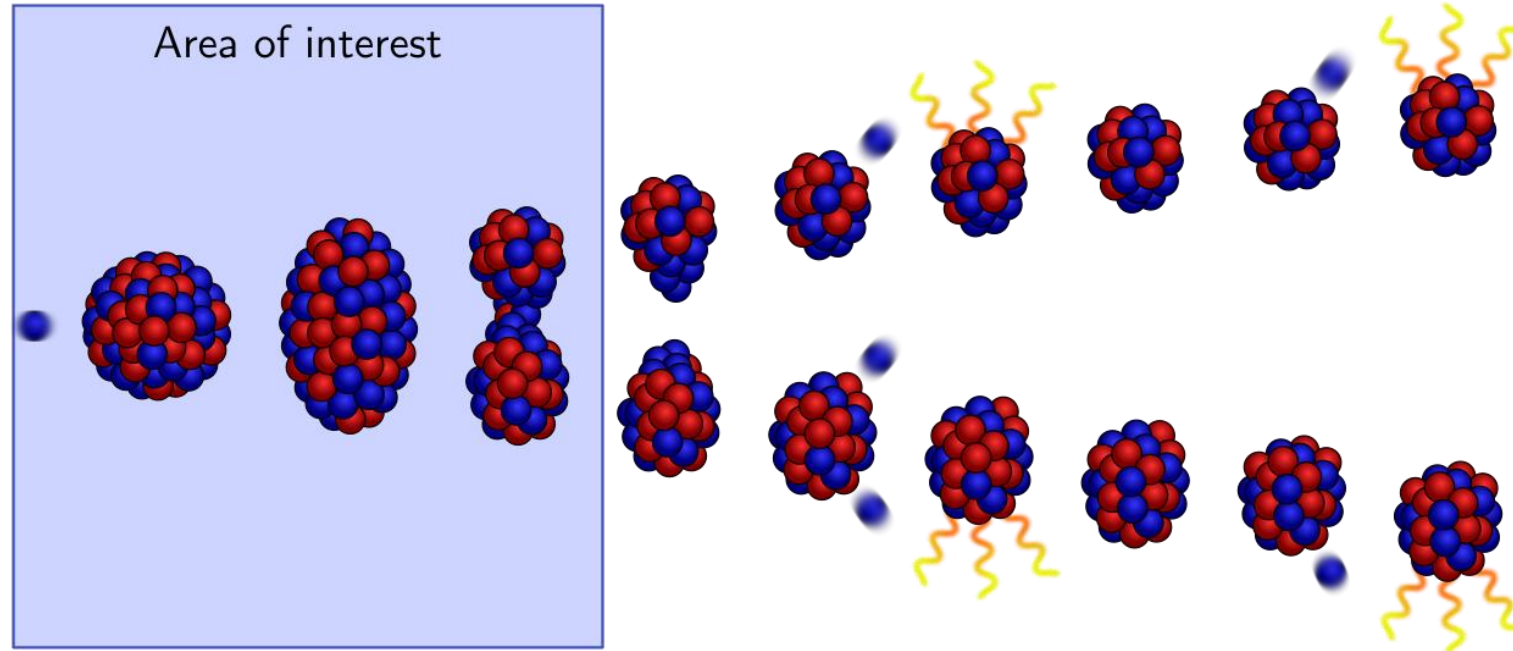
We describe the fission process using a fully microscopic approach.



We simulate the time-evolution of the deformations leading to the formation of the fragments.



Our theoretical framework leads to Schrodinger-like equation.

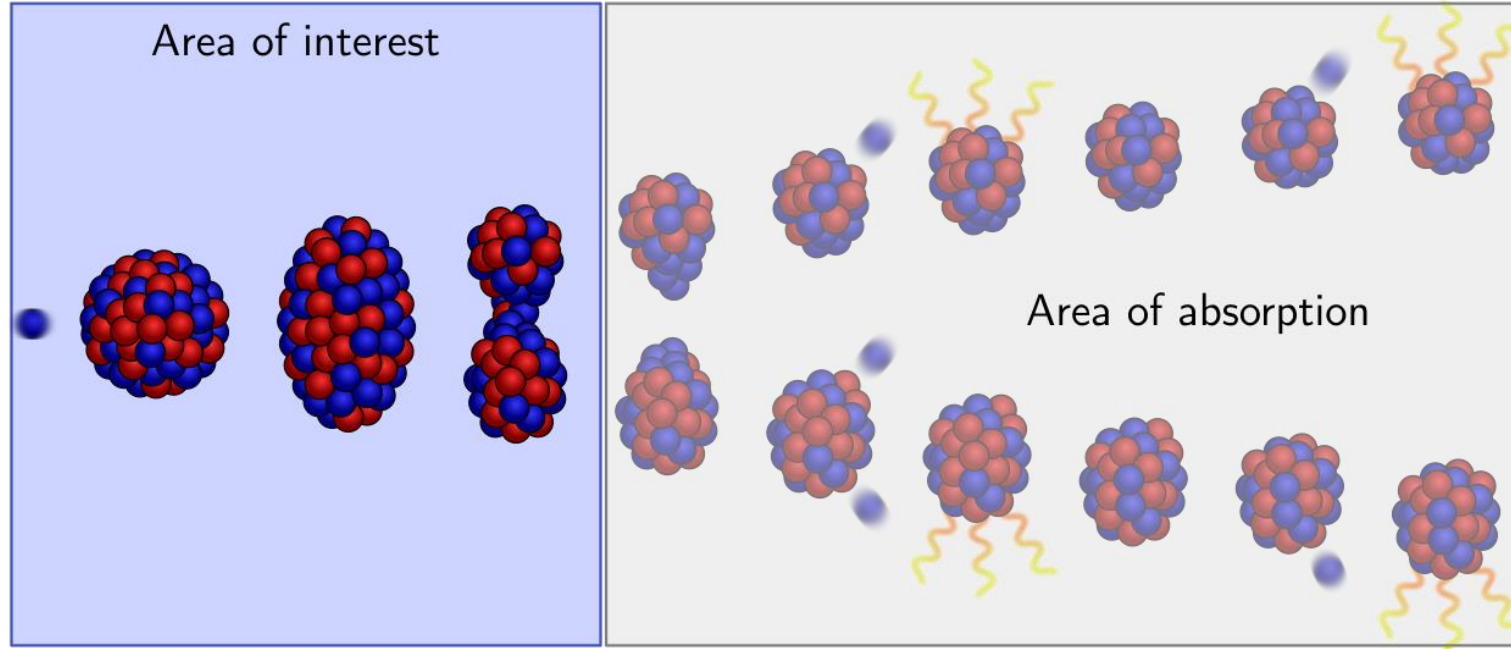


$$i\hbar \frac{\partial g(\mathbf{q}, t)}{\partial t} = \left[-\frac{\hbar^2}{2} \gamma^{-\frac{1}{2}}(\mathbf{q}) \nabla \cdot \gamma^{\frac{1}{2}}(\mathbf{q}) B(\mathbf{q}) \nabla + V(\mathbf{q}) \right] g(\mathbf{q}, t)$$

Equation describing the fission dynamics at a microscopic level with quantum effects:

- local **complex-valued** diffusion equation with real-valued \mathbf{q} -dependent coefficients $\gamma^{\pm\frac{1}{2}}$, B and V .
- \mathbf{q} : collective degrees of freedom describing the nucleus' deformation (2-D \rightarrow **3-D**).
- $g(\mathbf{q}, t)$: nucleus' probability amplitude to be in \mathbf{q} at time t (complex-valued).

We add an absorption term to limit our description to the area of interest.



$$i\hbar \frac{\partial g(\mathbf{q}, t)}{\partial t} = \left[-\frac{\hbar^2}{2} \gamma^{-\frac{1}{2}}(\mathbf{q}) \nabla \cdot \gamma^{\frac{1}{2}}(\mathbf{q}) B(\mathbf{q}) \nabla + V(\mathbf{q}) - iA(\mathbf{q}) \right] g(\mathbf{q}, t)$$

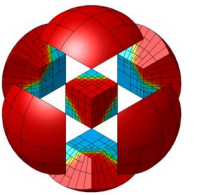
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We developed a new MFEM-based solver able to tackle three collective degrees of freedom.

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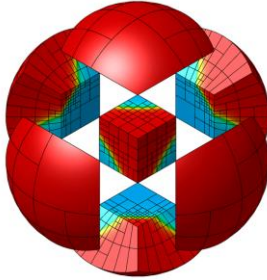
1. We discretize the collective space using the Finite Element Method with MFEM.



2. We developed a high-order (≥ 10) numerical time-discretization scheme.

3. We revamped our approach to predict the fission fragment properties.

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$$i\hbar \frac{\partial g(\mathbf{q}, t)}{\partial t} = \left[-\frac{\hbar^2}{2} \gamma^{-\frac{1}{2}}(\mathbf{q}) \nabla \cdot \gamma^{\frac{1}{2}}(\mathbf{q}) B(\mathbf{q}) \nabla + V(\mathbf{q}) - iA(\mathbf{q}) \right] g(\mathbf{q}, t)$$

Partition space

Collective space (infinite size)

Q_{21}	Q_{22}	Q_{25}	Q_{26}	Q_{37}	Q_{38}	Q_{41}	Q_{42}
Q_{20}	Q_{23}	Q_{24}	Q_{27}	Q_{36}	Q_{39}	Q_{40}	Q_{43}
Q_{19}	Q_{18}	Q_{29}	Q_{28}	Q_{35}	Q_{34}	Q_{45}	Q_{44}
Q_{16}	Q_{17}	Q_{30}	Q_{31}	Q_{32}	Q_{33}	Q_{46}	Q_{47}
Q_{15}	Q_{12}	Q_{11}	Q_{10}	Q_{53}	Q_{52}	Q_{51}	Q_{48}
Q_{14}	Q_{13}	Q_8	Q_9	Q_{54}	Q_{55}	Q_{50}	Q_{49}
Q_1	Q_2	Q_7	Q_6	Q_{57}	Q_{56}	Q_{61}	Q_{62}
Q_0	Q_3	Q_4	Q_5	Q_{58}	Q_{59}	Q_{60}	Q_{63}

Define basis functions

The $g(\mathbf{q}, t)$ scalar field is smooth:

- H1-conformal real basis,
- piecewise p -degree polynomials.
- Basis function noted φ_i

Project the equation

Orthogonal projection on the basis functions using the scalar product:

$$(\phi, \psi) = \int \phi^*(\mathbf{q}) \psi(\mathbf{q}) \gamma^{\frac{1}{2}}(\mathbf{q}) d\mathbf{q}$$

$$i\hbar M \frac{\partial G(t)}{\partial t} = (D + V - iA) G(t)$$

$$M_{ij} = \left(\varphi_i, \gamma^{\frac{1}{2}} \varphi_j \right),$$

$$V_{ij} = \left(\varphi_i, \gamma^{\frac{1}{2}} V \varphi_j \right),$$

$$A_{ij} = \left(\varphi_i, \gamma^{\frac{1}{2}} A \varphi_j \right),$$

$$D_{ij} = (\varphi_i, -\nabla \cdot \Lambda \nabla \varphi_j) \\ = (\nabla \varphi_i, \Lambda \nabla \varphi_j) + \cancel{\partial B}$$

→ **Sparse & partially assembled!**

2. We developed a high-order time-discretization scheme.

$$i\hbar M \frac{\partial G(t)}{\partial t} = (K - iA)G(t)$$

First-order linear differential equation:

$$G(t) = U(t) G_0$$

$$U(t) = \exp\left(-\frac{it}{\hbar} M^{-1}(K - iA)\right)$$

- Inverting M is **expensive**,
- M^{-1} and $U(t)$ are **not sparse**,
- Inverting M is **not stable**.

→ **Only compute the action of $U(t)$.**

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Step 1: Taylor-expand $U(t)$

$$U(t) \approx \sum_{n=0}^N \frac{1}{n!} \left(-\frac{it}{\hbar} M^{-1}(K - iA)\right)^n$$

→ Only valid around $t = 0$.

Step 2: Discretize time

$$t = \delta t + \dots + \delta t$$
$$U(t) = U(\delta t) \dots U(\delta t)$$

$$U(\delta t) \approx \sum_{n=0}^N \frac{1}{n!} \left(-\frac{i\delta t}{\hbar} M^{-1}(K - iA)\right)^n$$

Step 3: Evaluate action of $U(t)$

$$U(\delta t)G \approx \sum_{n=0}^N \frac{1}{n!} F^n G$$

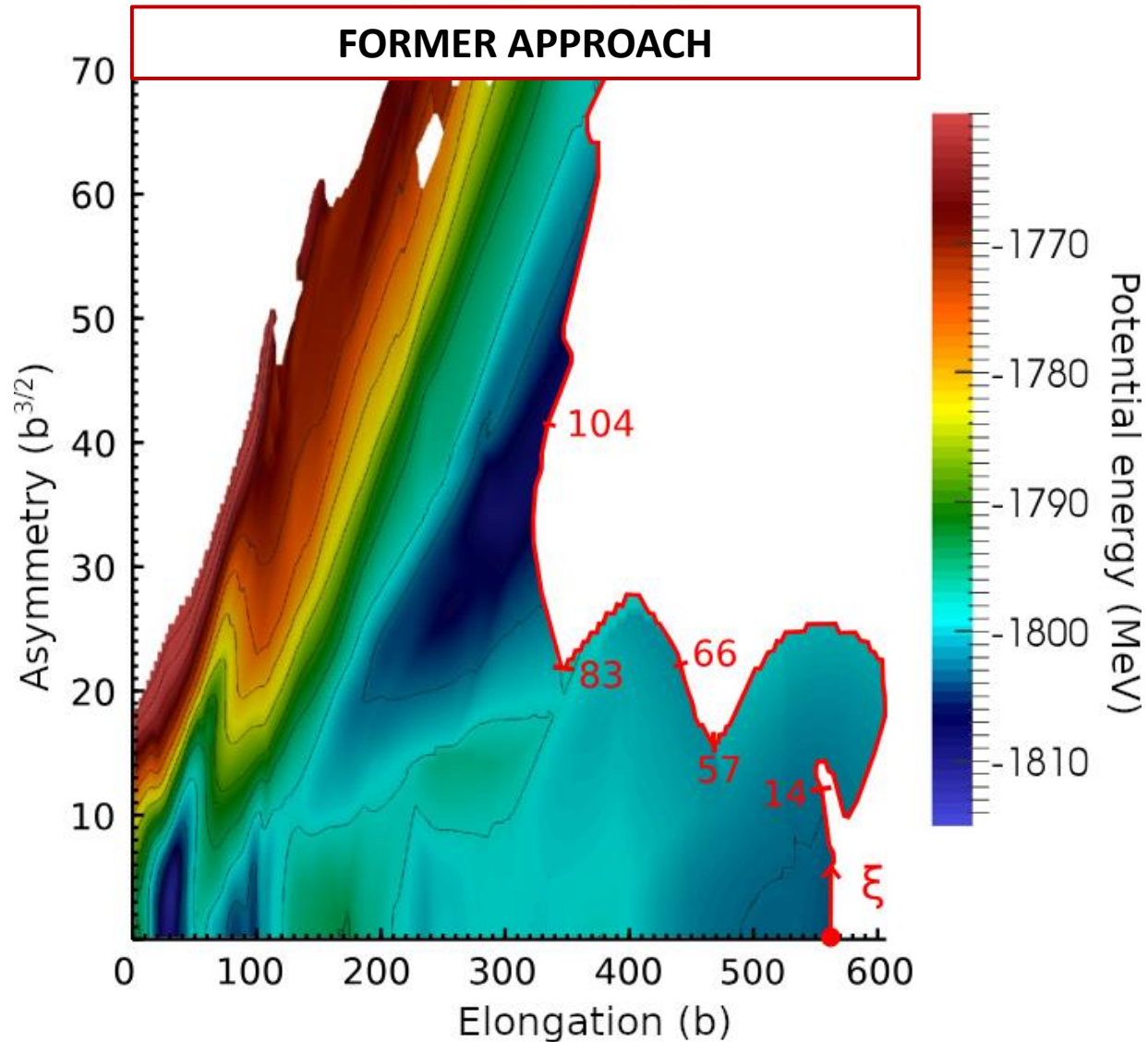
$$F = -\frac{i\delta t}{\hbar} M^{-1}(K - iA)$$

Determine $F^n G$ by recurrence:

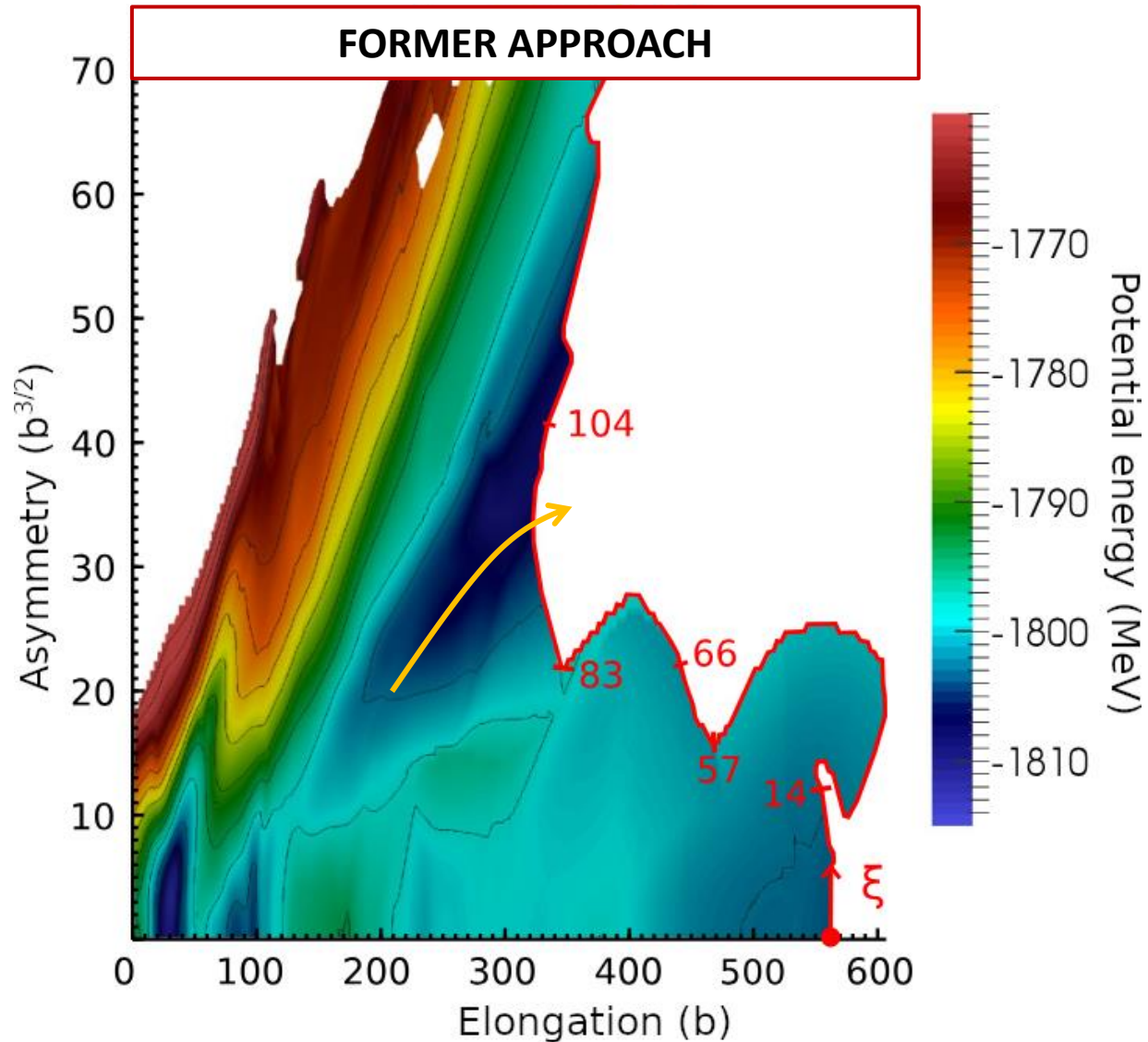
$$F^0 G = G$$

$$M(F^{n+1}G) = -\frac{i\delta t}{\hbar} (K - iA) (F^n G)$$

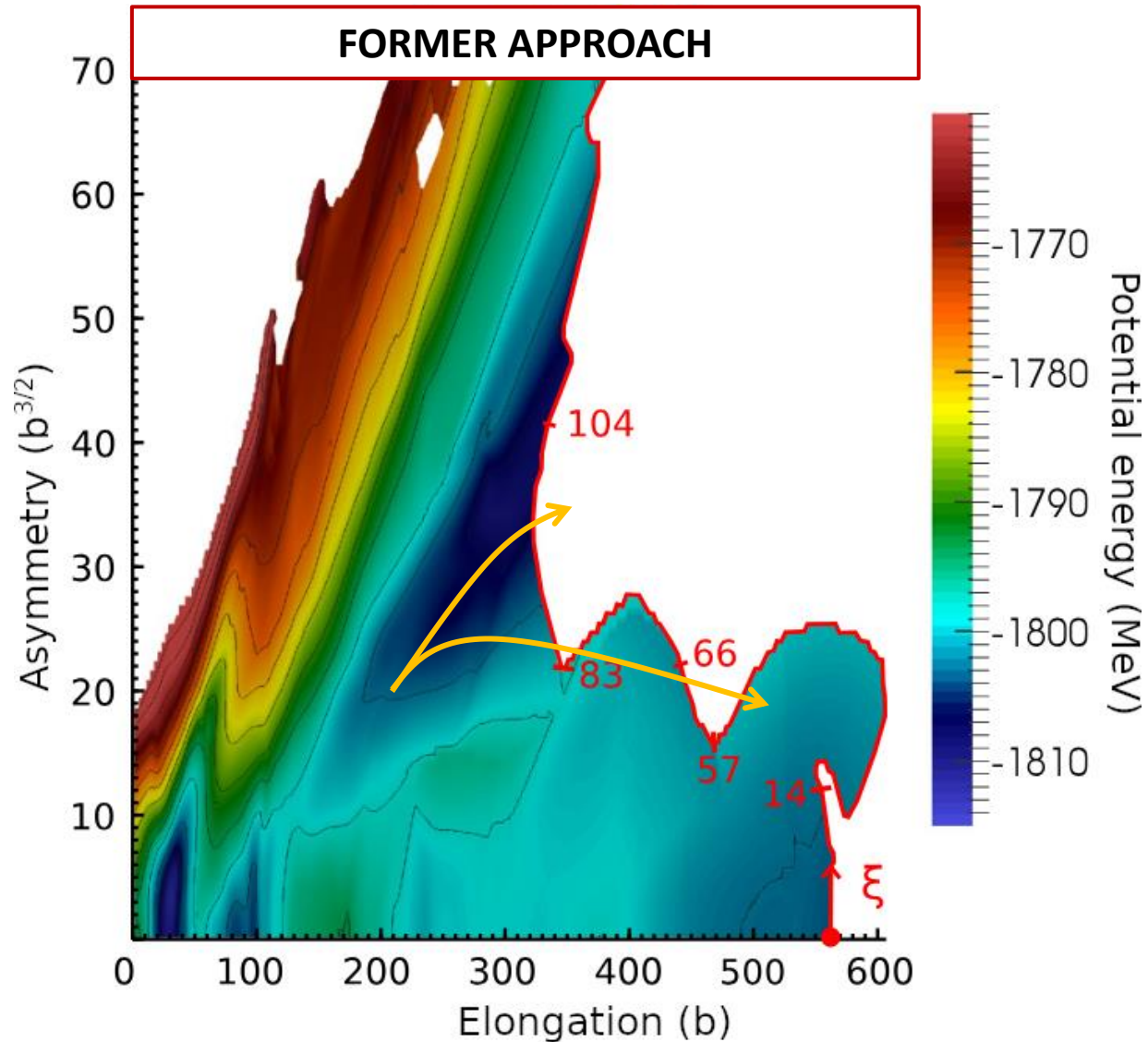
3. We revamped our approach to predict the fragment properties.



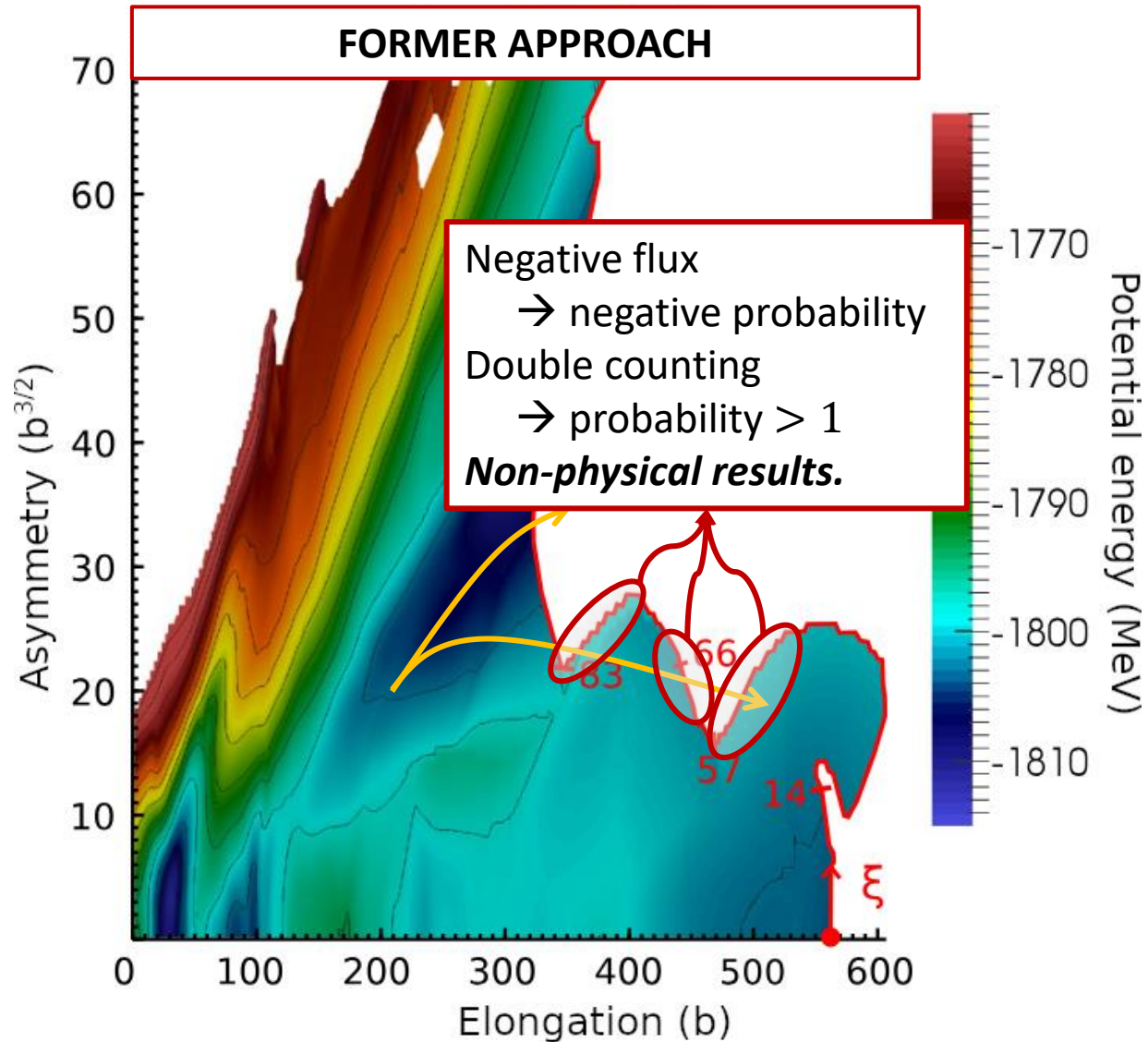
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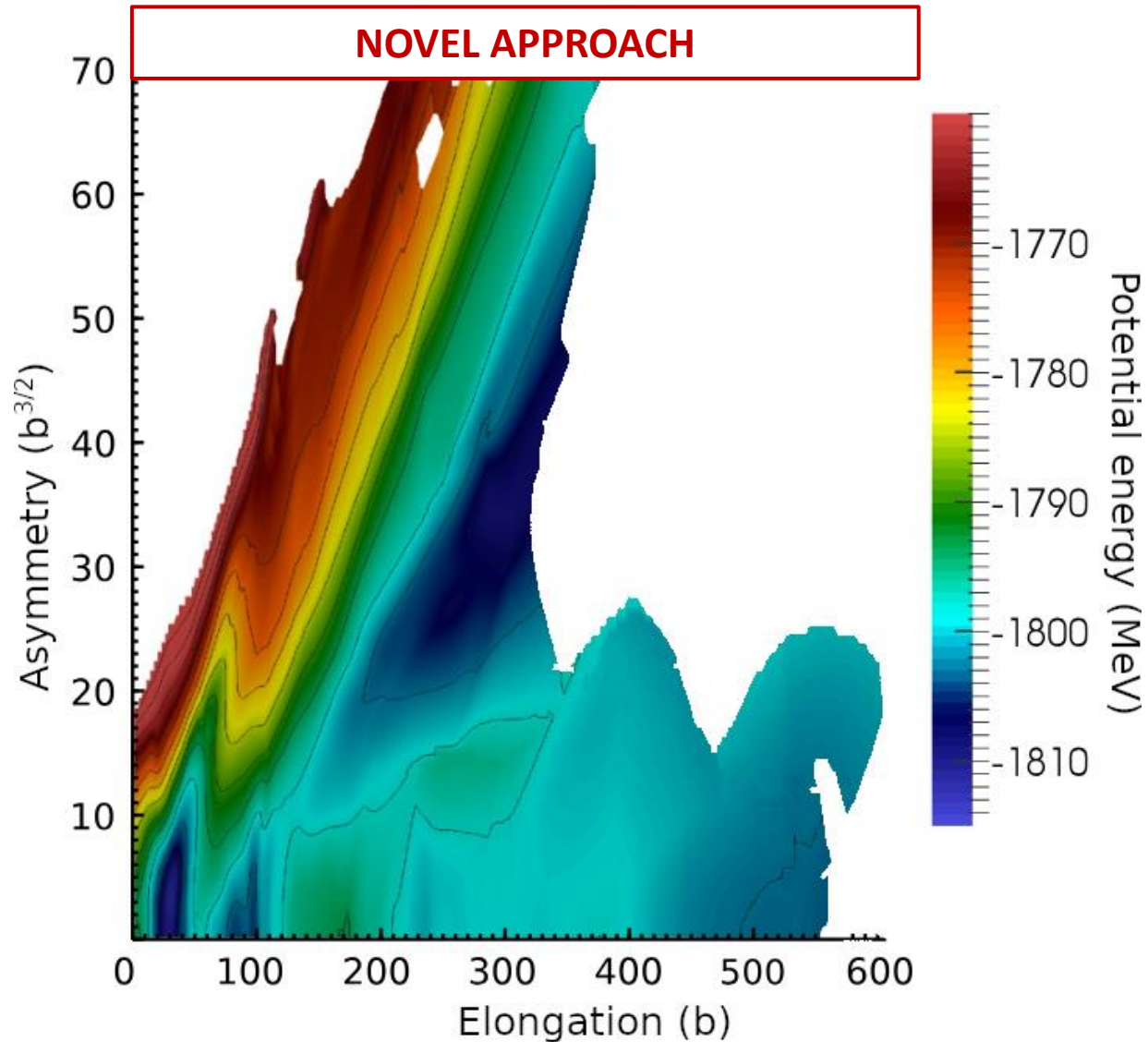
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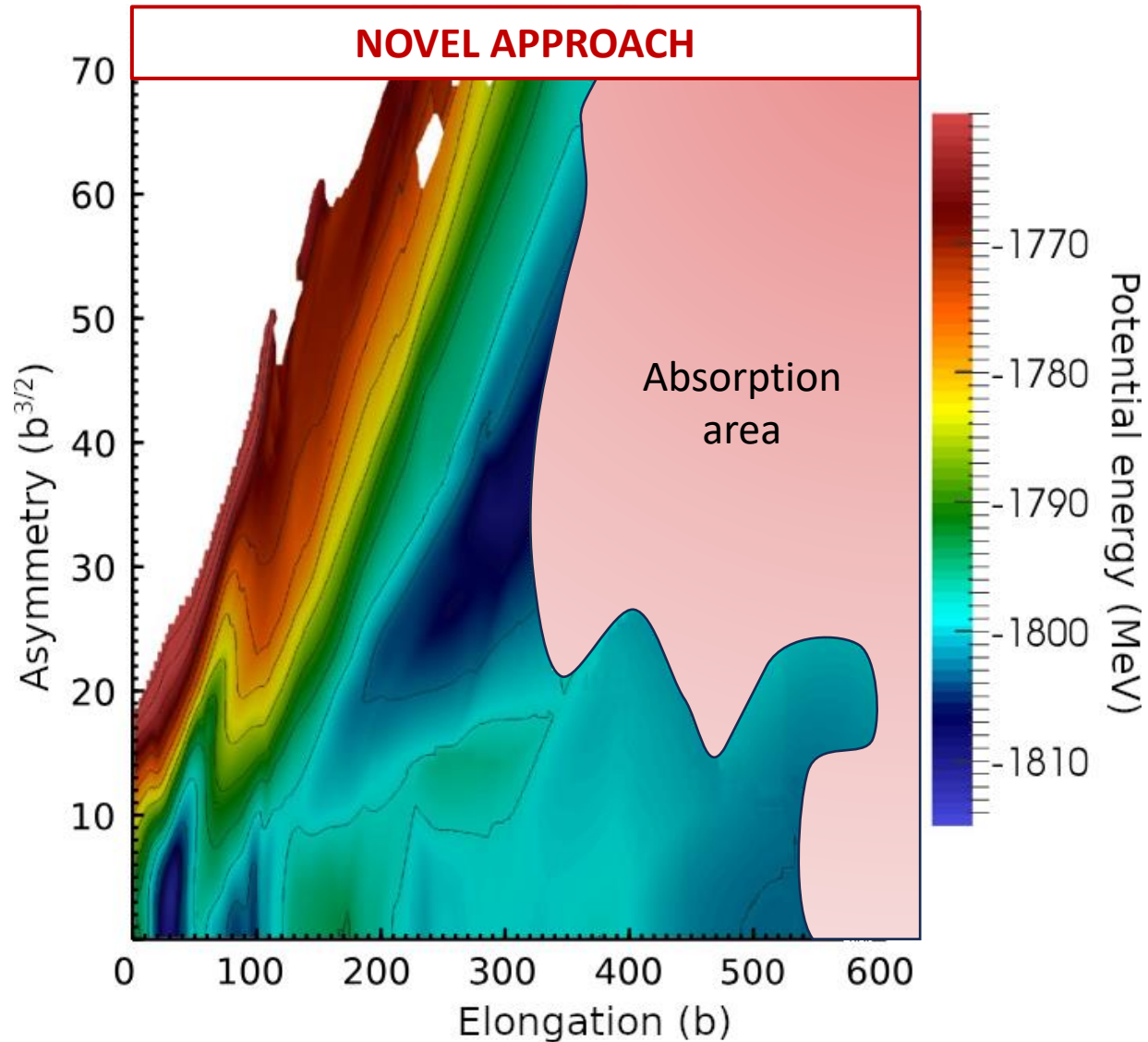
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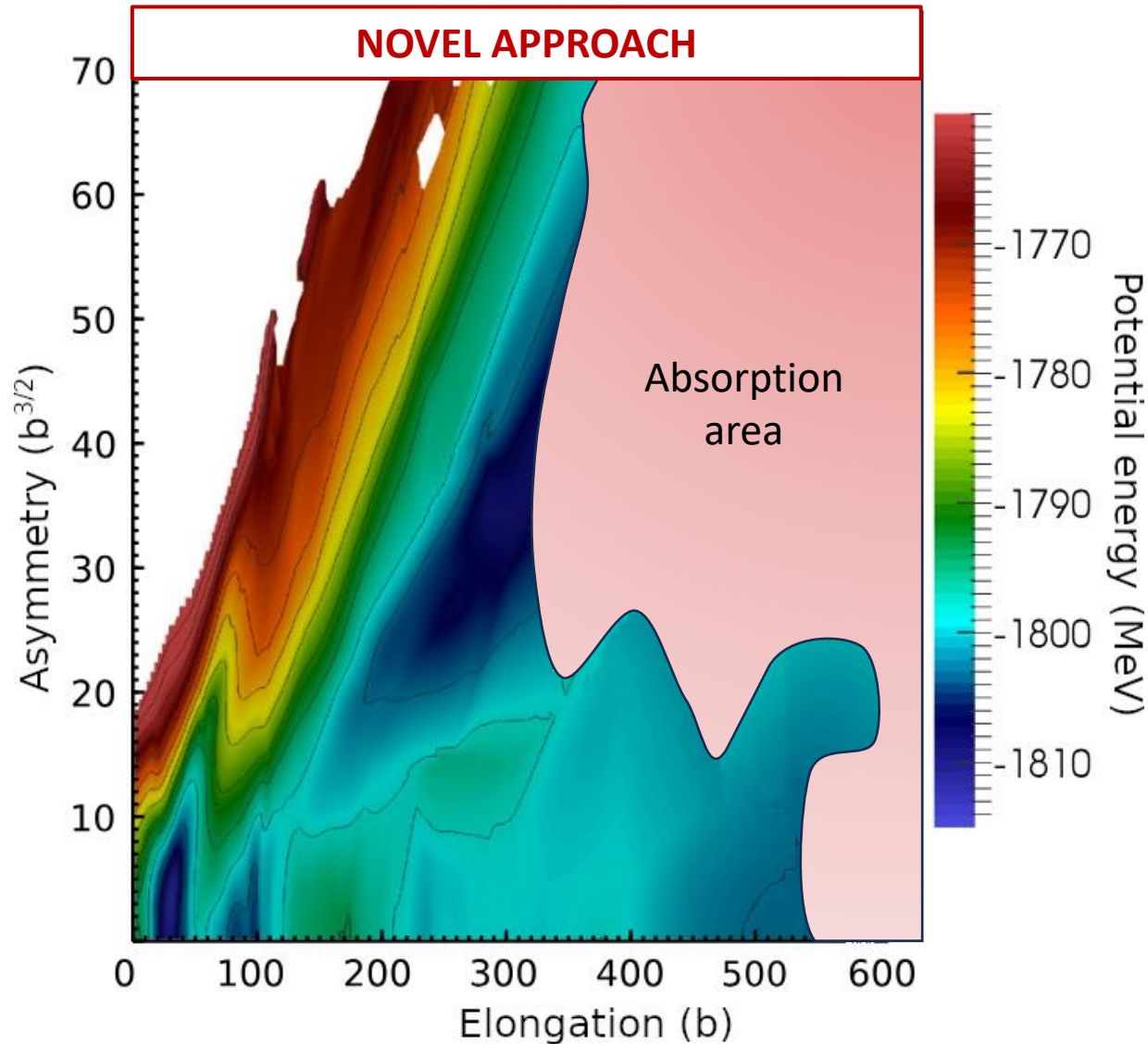
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Probability $\rho_{\text{abs}}(\mathbf{q})$ to be absorbed in \mathbf{q} :

$$\rho_{\text{abs}}(\mathbf{q}) = \left(1 - e^{-\frac{\delta t}{\hbar} A(\mathbf{q})}\right) |g(\mathbf{q}, t)|^2 \gamma^{\frac{1}{2}}(\mathbf{q})$$

Probability to measure a fragment with N particles at a given \mathbf{q} :

$$P(X | \mathbf{q}) \approx \int_{X-\frac{1}{2}}^{X+\frac{1}{2}} e^{-\frac{1}{2} \frac{(x - \bar{X}(\mathbf{q}))^2}{\sigma_X^2}} \frac{dx}{\sqrt{2\pi\sigma_X^2}}$$

→ Probability of a fragment with X particles:

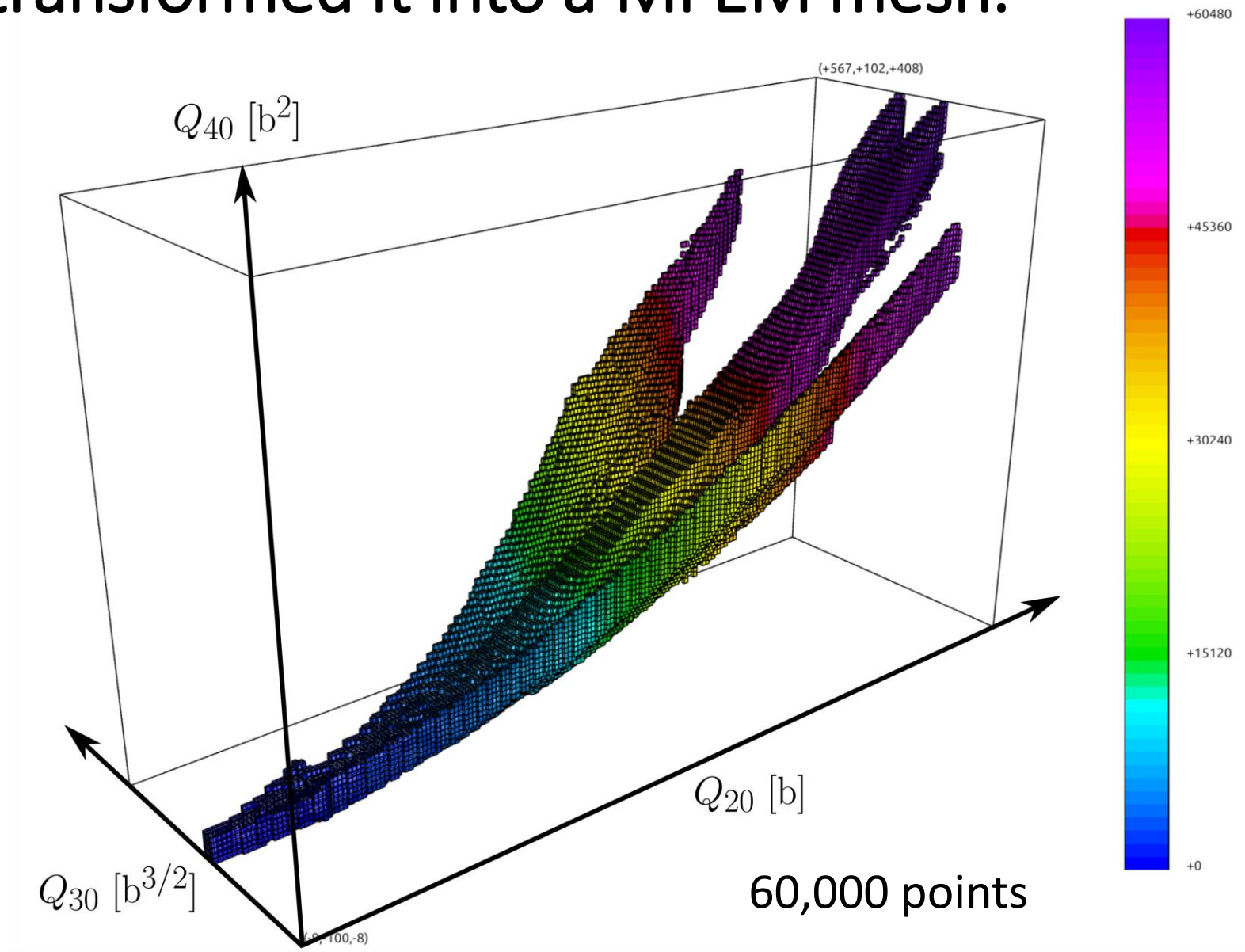
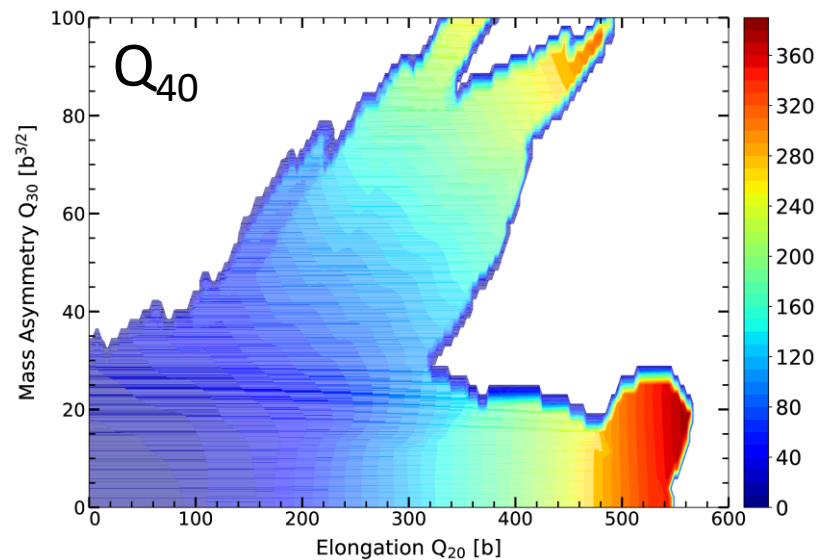
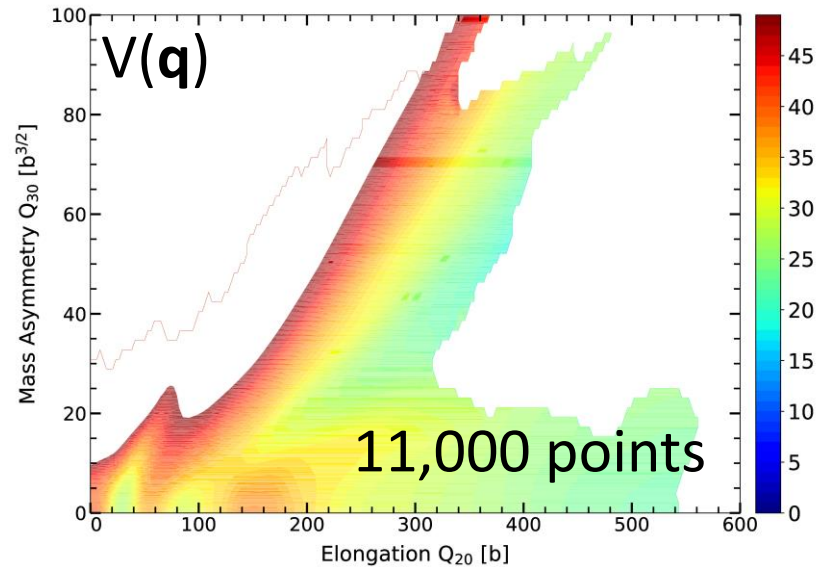
$$Y(X) \approx \int P(X | \mathbf{q}) \rho_{\text{abs}}(\mathbf{q}) d\mathbf{q}$$

$$X = A_f, Z_f$$

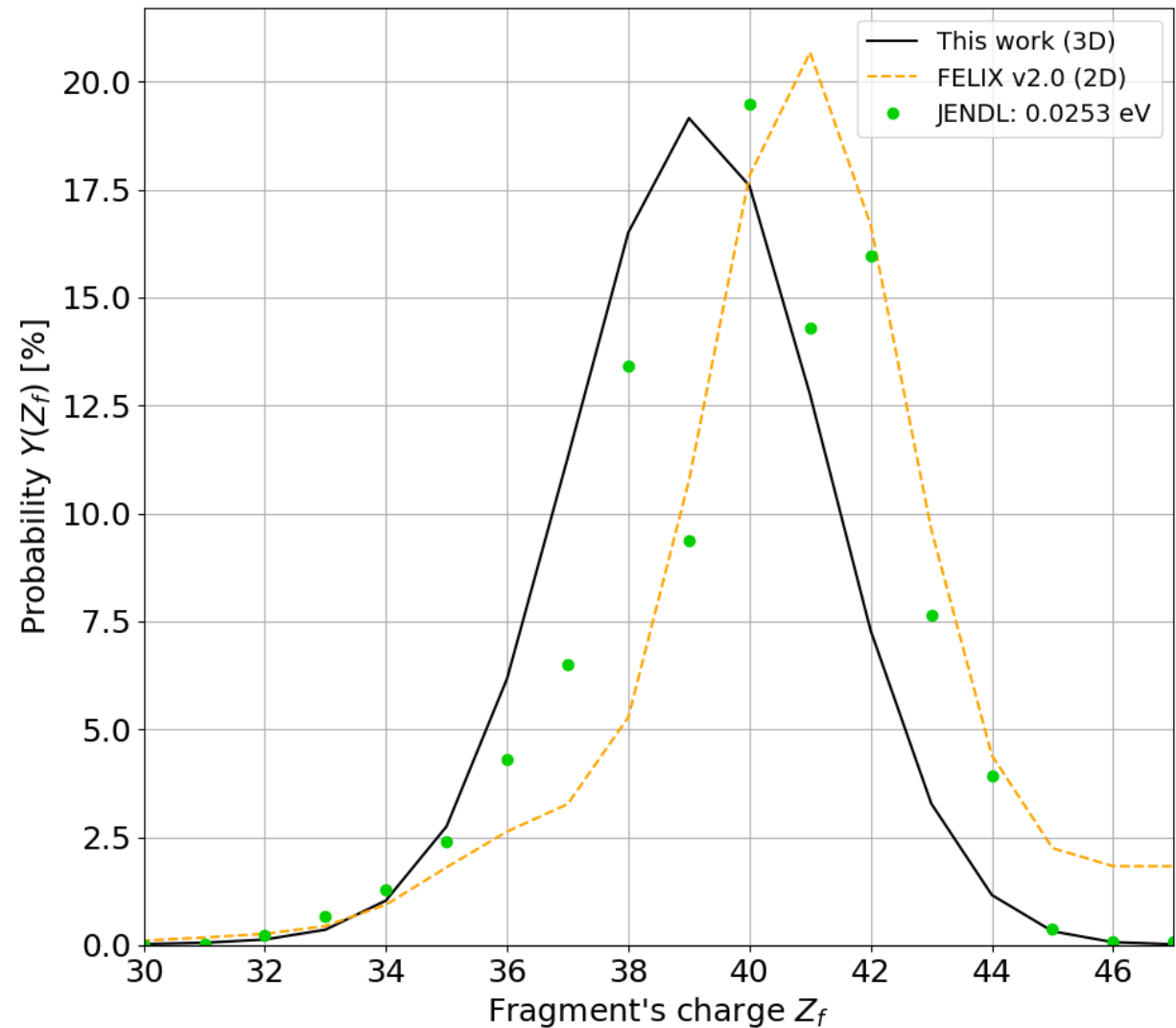
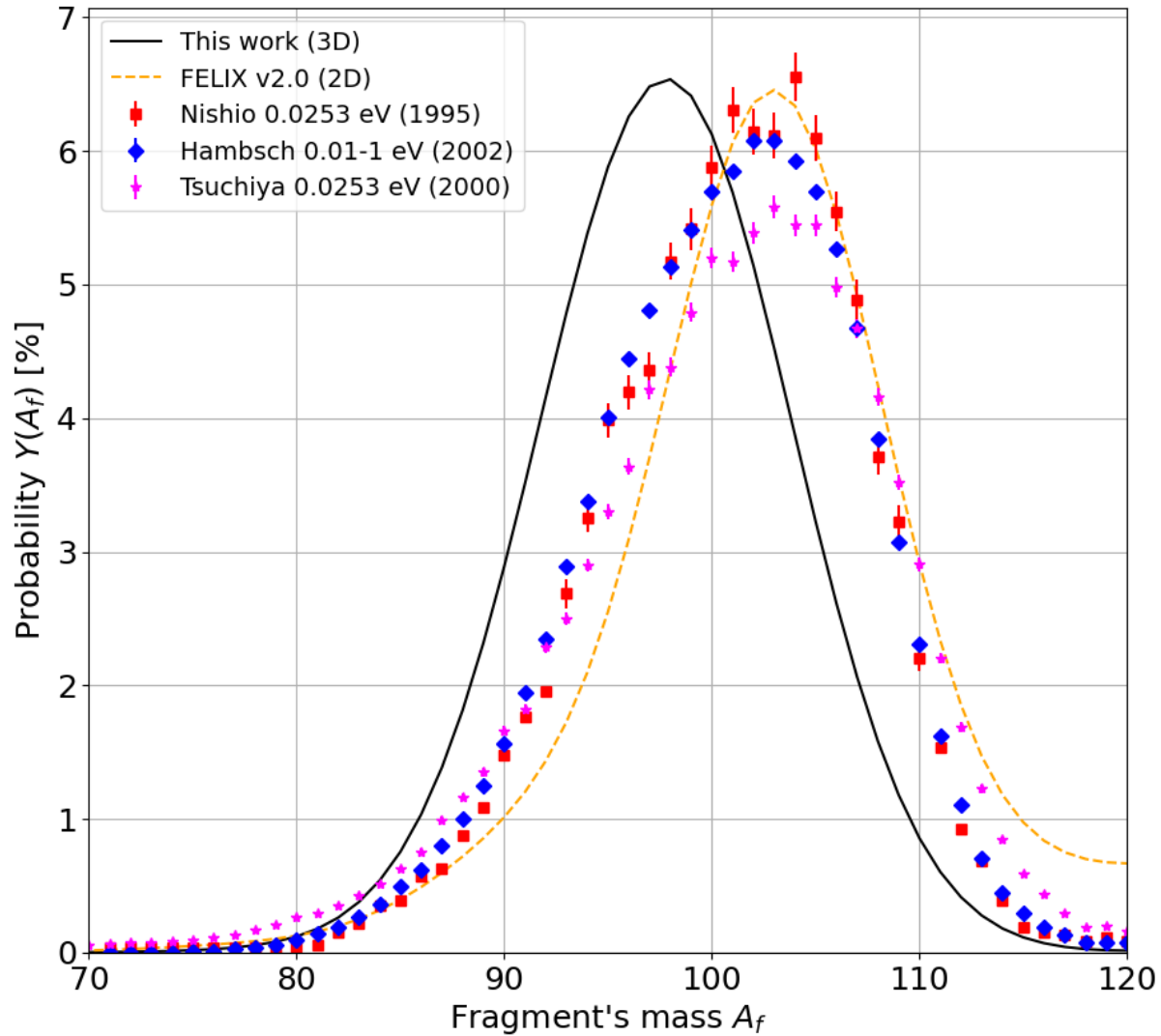
Partial assembly w/ MFEM

→ **≈ 300 kernels** loaded at the same time!

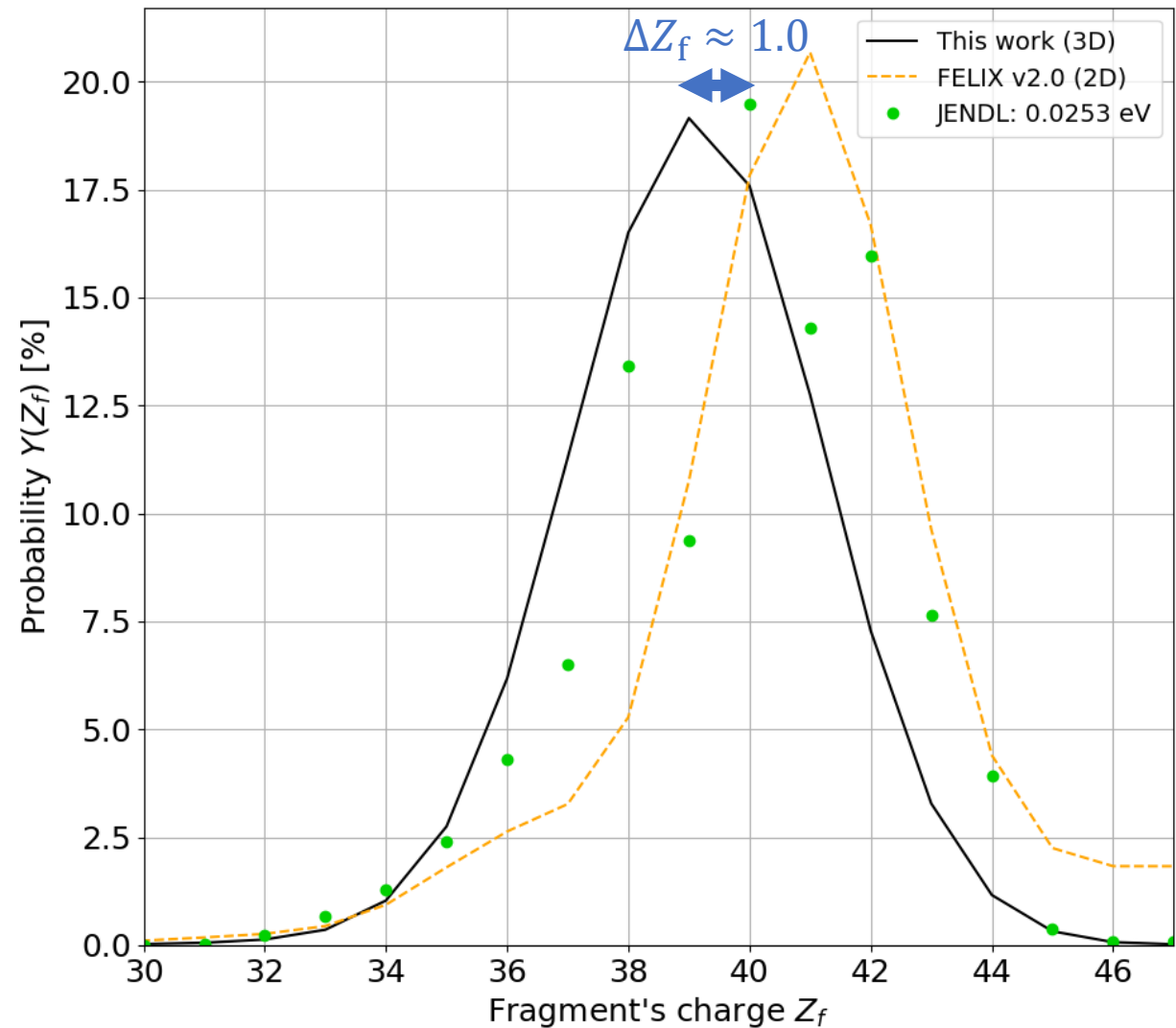
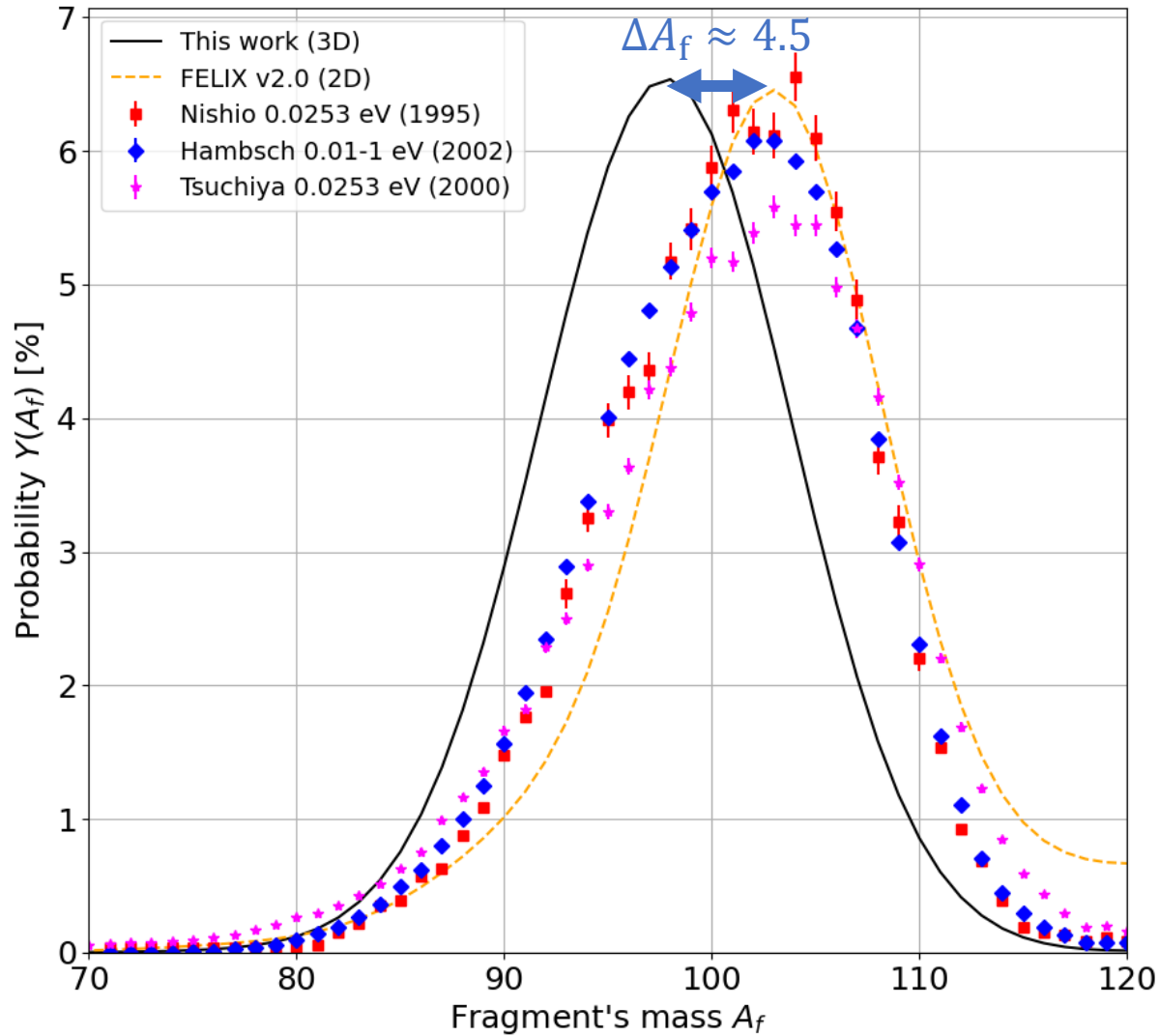
We generated a 3D potential energy surface using a Skyrme functional on ^{240}Pu and transformed it into a MFEM mesh.



Finally, we extracted fission fragment mass and charge yields using our new pipeline.



Our (very) preliminary results are encouraging: symmetric fission is better reproduced, but the mass yields show a bigger deviation.



These preliminary results are promising, and we are working on many further improvements.

- Improve the handling of the potential energy landscape (interpolation).
- Define the initial state from the eigenstates of the extrapolated potential wells.
- Strengthen the order of the time-propagation approach.
- Enable extrapolation of the potential energy landscape.
- Couple Fidelis with particle-number projection in the fragments.
- Study the impact of three degrees of freedom on several fissioning systems.
- Determine fission observables using a Gogny interaction.
- Package the library (CMake, Pybind11, Doxygen, etc).
- ...

Thank you