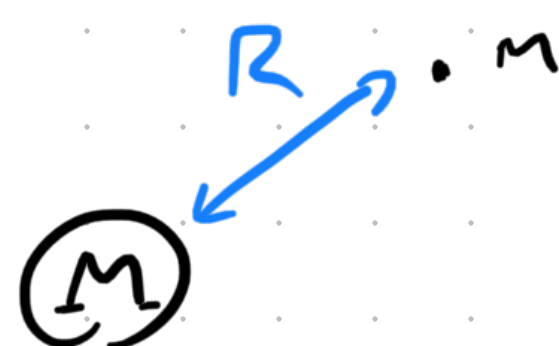


Orbits (Topic 7 in F=ma Physics)

only force is gravity

what is the period of orbit?



$$ma_{\text{cent}} = \frac{mv^2}{R} = F_{\text{radially inward}} = F_{\text{grav}} = \frac{GMm}{R^2}$$

$$\frac{mv^2}{R} = \frac{GMm}{R^2} \Rightarrow \boxed{v = \sqrt{\frac{GM}{R}}}$$

each orbit is $2\pi R$ long,

$$T = \frac{2\pi R}{v} = \frac{2\pi R}{\sqrt{GM/R}} = \boxed{2\pi \sqrt{\frac{R^3}{GM}}} = T$$

orbits are stable when governed by a radially
inwards force that scales as R^{-2}

$$F \propto R^{-2} \quad F = \eta R^{-2} \quad \eta = \text{force constant.}$$

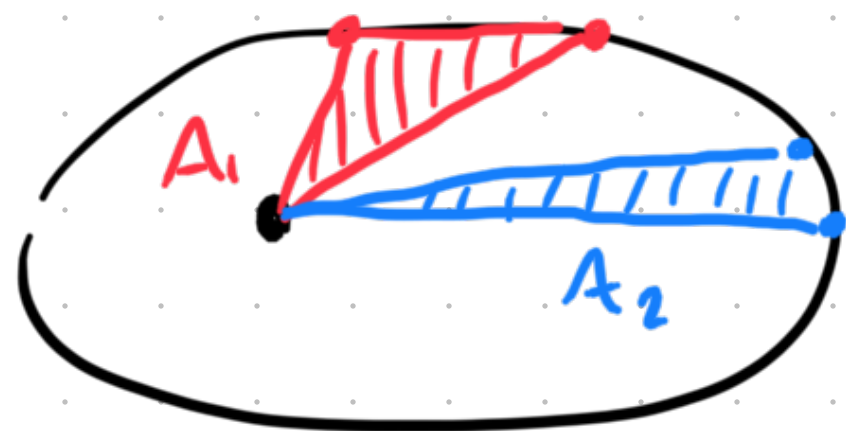
$$\eta R^{-2} = F = \frac{mv^2}{R} \Rightarrow v = \sqrt{\frac{\eta}{mR}}$$

$$\eta = GMm \quad (\text{gravity}) \quad v = \sqrt{\frac{GM}{R}}$$

Kepler's Laws

I. all bound orbits are elliptical, with the orbited body at a focus.

II. a line joining the orbiter and orbited sweeps equal areas in equal time.



$$A_1 = A_2$$

$$t_1 = t_2$$

$$\dot{\theta}_1 > \dot{\theta}_2$$

or

$$\omega_1 > \omega_2$$

or

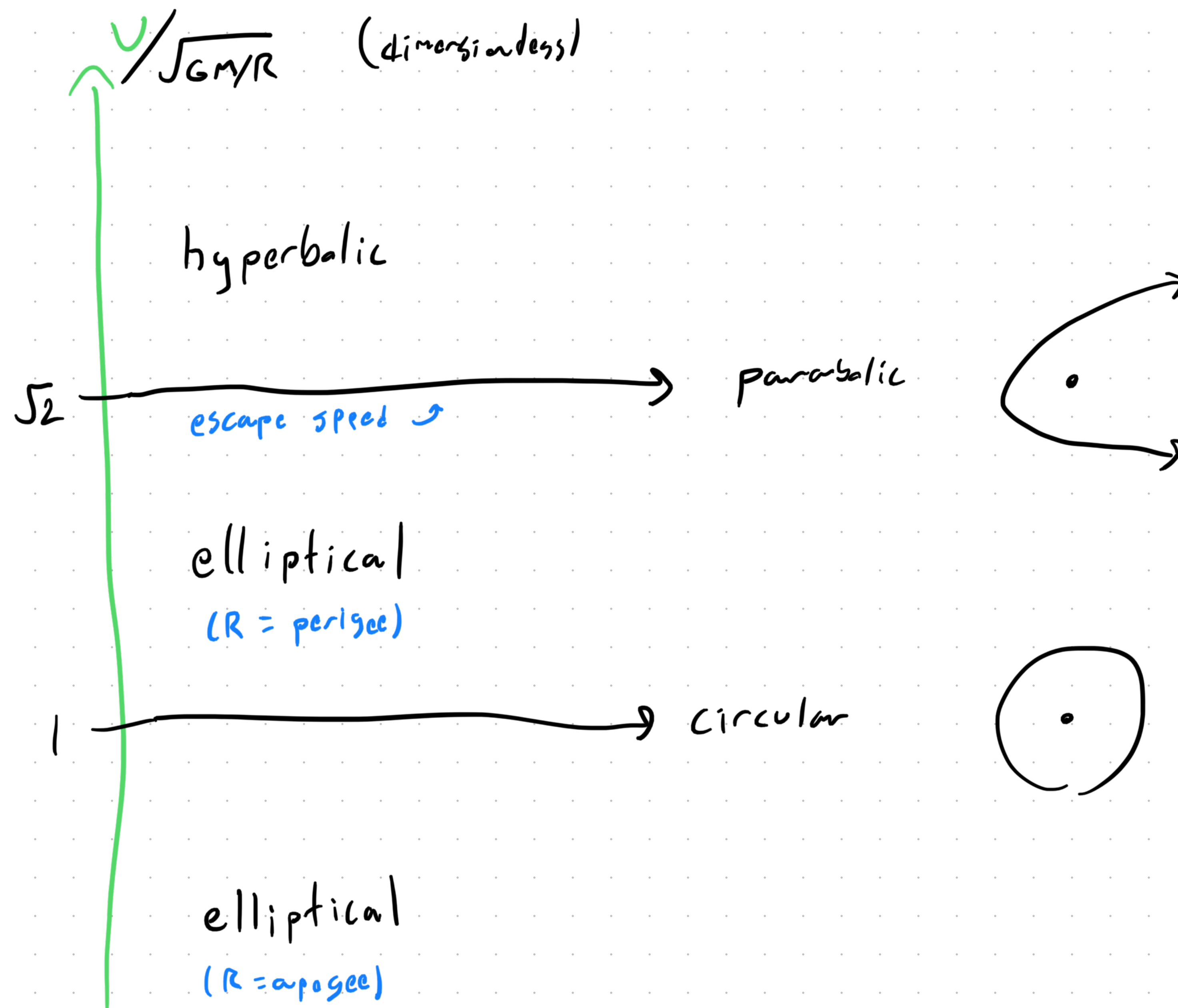
$$v_1 > v_2$$

III.

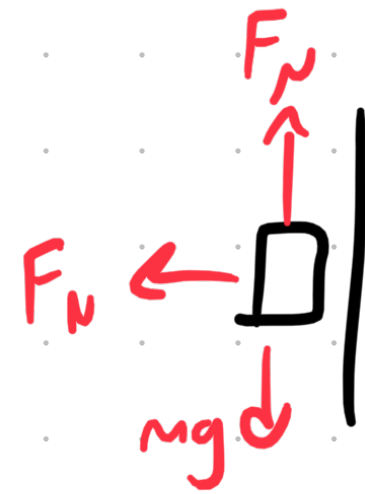
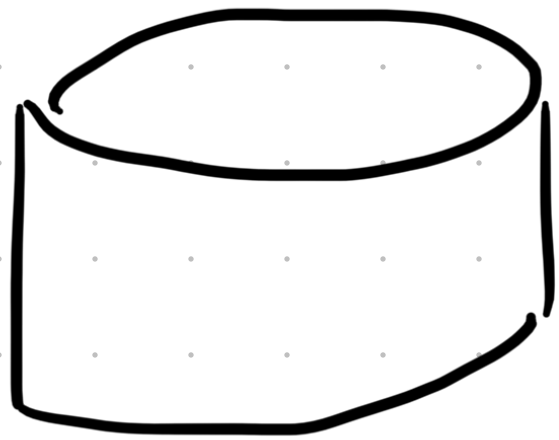
$$T^2 \propto R^3$$

derivation above can be generalized to ellipses.

Non-circular and unbound orbits.



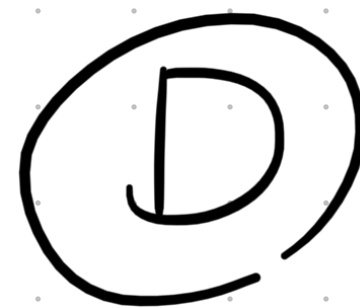
2017 #1



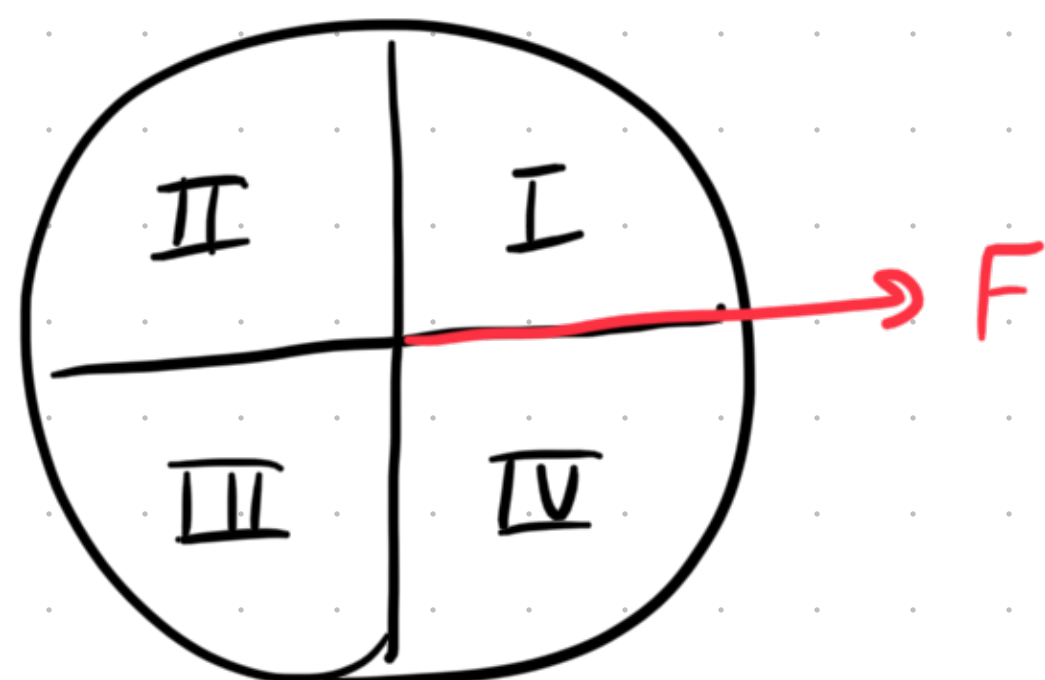
$$F_N = \frac{ms^2}{R}$$

$$F_N = mg = \mu \frac{ms^2}{R}$$

$$\mu \propto s^{-2}$$



2017 #18



any point that has
zero acceleration?

$$\vec{a}_c + \vec{a}_t + \vec{a}_F$$

$$\text{for } \vec{a}_{\text{total}} = 0, \quad \vec{a}_c + \vec{a}_t = -\vec{a}_F$$

so due left.

~~A, B, C~~

r = point radius

R = disk radius

$$\|\vec{a}_c + \vec{a}_t\| = \|\vec{a}_F\|?$$

$$\|\vec{a}_c + \vec{a}_t\| = \sqrt{a_c^2 + a_t^2} = r\sqrt{\omega^4 + d^2}$$

$$\|\vec{a}_F\| = \alpha R$$

$$\sqrt{\omega^4 + d^2} > |\alpha|$$

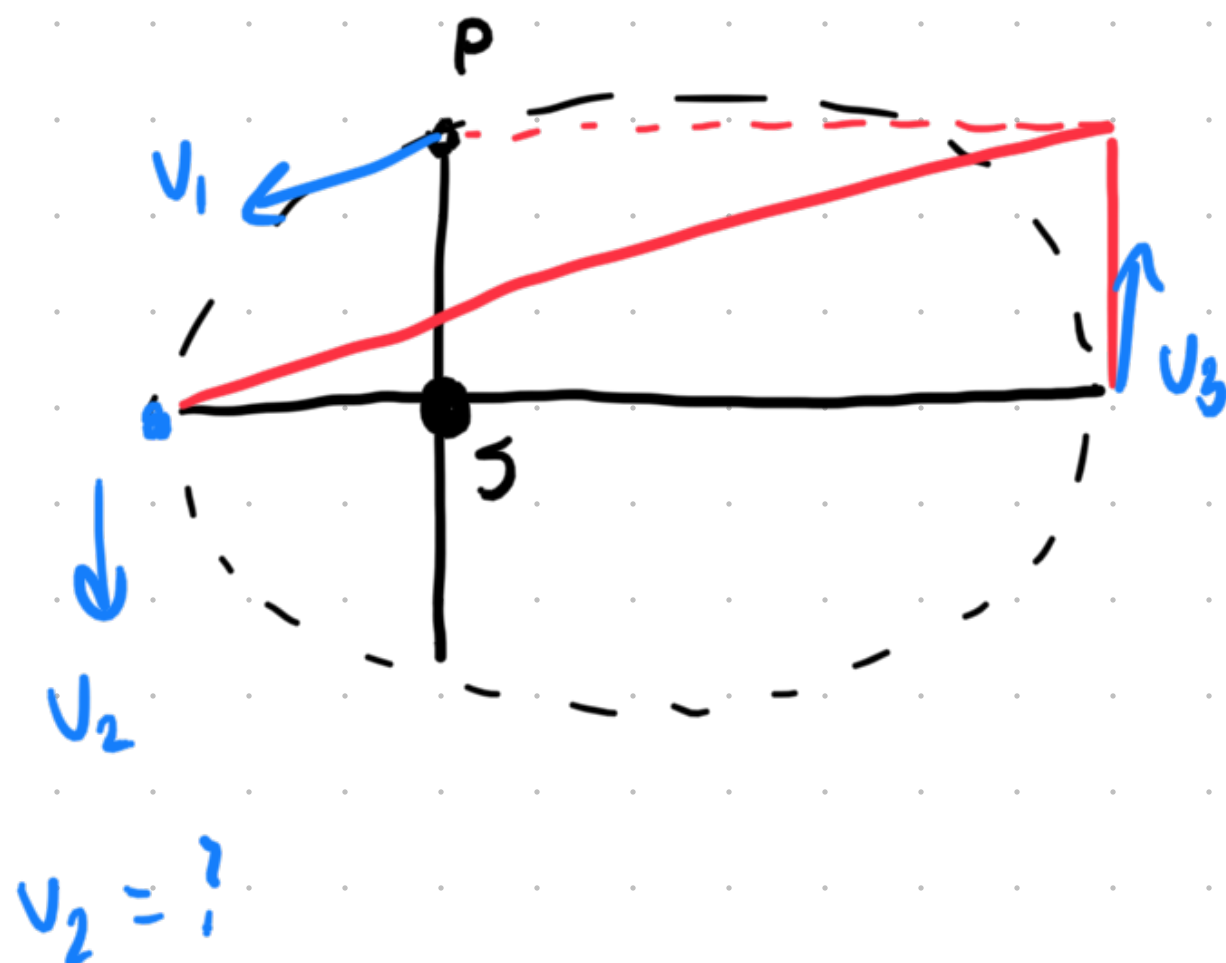
IVT says that $\exists r \in (0, R)$

$$\text{st. } r\sqrt{\omega^4 + d^2} = \alpha R.$$

So, yes!

(D)

2017 #25



$$\overline{SP} = ?$$

$$\overline{SP} + \sqrt{\overline{SP}^2 + a^2} = 2a$$

$$\overline{SP} = \frac{3a}{4}$$

Conservation of energy

$$\frac{1}{2} m v_2^2 - \frac{GMm}{a/2} = \frac{1}{2} m v_3^2 - \frac{GMm}{3a/2}$$

$$K2L: \frac{m v_2 a}{2} = \frac{3 m v_3 a}{2} \Rightarrow v_3 = \frac{1}{3} v_2$$

$$\frac{4}{9} m v_2^2 = \frac{GMm}{a/2} - \frac{GMm}{3a/2} = \frac{4GMm}{3a}$$

$$v_2^2 = \frac{3GM}{a}$$

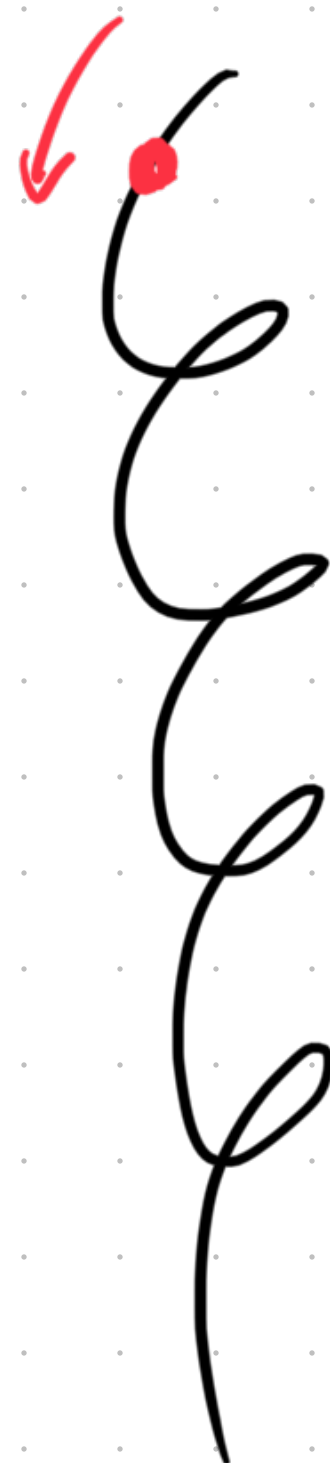
$$\frac{1}{2} m v_2^2 - \frac{GMm}{a/2} = \frac{1}{2} m v_1^2 - \frac{GMm}{3a/4}$$

$$v_1^2 = v_2^2 - \frac{4GM}{a} + \frac{8GM}{3a} = \frac{5GM}{3a}$$

$$v_2^2 / v_1^2 = 9/5 \Rightarrow v_2 / v_1 = 3/\sqrt{5}$$

(A)

2016 #4

radius r angle α

$$a_t = g \sin \alpha$$

$$a_c = \frac{v_h^2}{r}$$

~~A, B, C~~

not constants

$$a(t=0) \neq 0$$

$$a_{tv} = a_t \sin \alpha = g \sin^2 \alpha$$

$$a_{th} = a_t \cos \alpha = g \sin \alpha \cos \alpha$$

$$a = \sqrt{a_t^2 + a_c^2} = \sqrt{g^2 \sin^2 \alpha + \left(\frac{g t \sin \alpha \cos \alpha}{r} \right)^2}$$

$$= \sqrt{P + Q t^4}$$

P, Q are positive constants

NOT LINEAR!

(D)

2016 #8

How would Kepler's Laws

change if $F_{\text{grav}} \propto R^{-3}$ rather than $F_{\text{grav}} \propto R^{-2}$?

III changes because we used gravity to prove it.

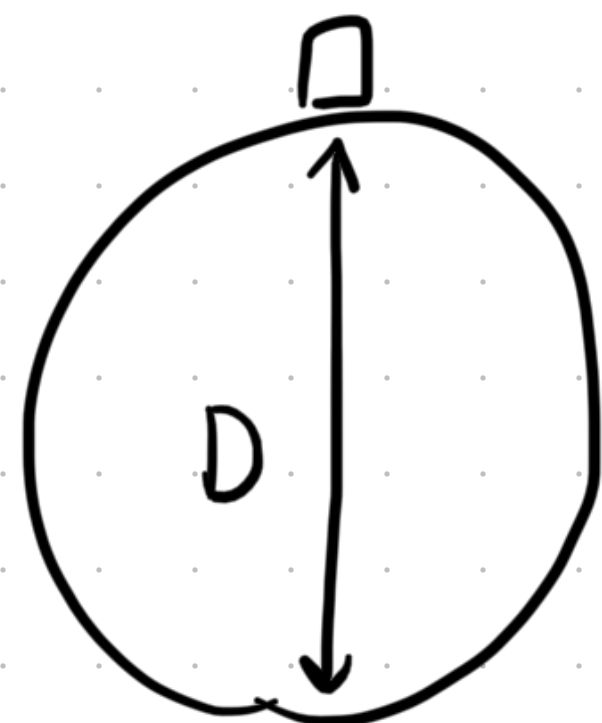
~~C D~~

I is also a consequence of gravity too.

II? that's just angular momentum conservation.

(B)

2016 #9



V when the bead
falls off?

$$mgh = \frac{1}{2}mv^2$$

$$h = R(1 - \cos\theta)$$

$$mg\cos\theta \geq \frac{mv^2}{R}$$

$$\cos\theta = \frac{v^2}{gR}$$

$$v = \sqrt{\frac{2gR}{3}} = \sqrt{\frac{gD}{3}}$$

conservation of energy

geometry of the problem

contact forces

(E)