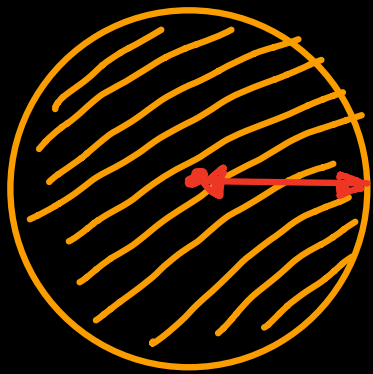


Stellar Structure

• Our Sun:



Radius, $R_{\odot} = 6.96 \times 10^8 \text{ m}$
 Mass, $M = 1.99 \times 10^{30} \text{ kg}$

• Forces present in the Sun:

- outward pressure:

Thermal Pressure

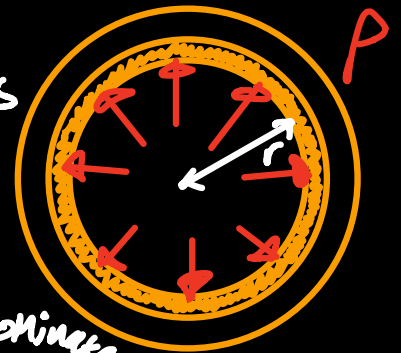
Radiation Pressure

* Degeneracy Pressure *

White D's
and neutron
stars only →

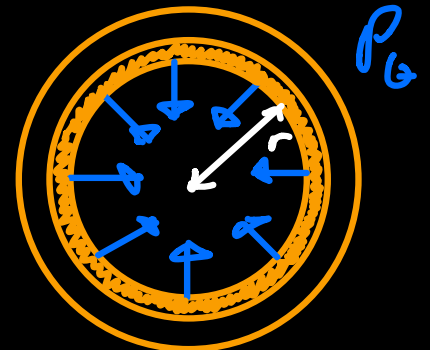
Dominates
MS

Dominates
> 100 M_{\odot}



- Inward Pressure:

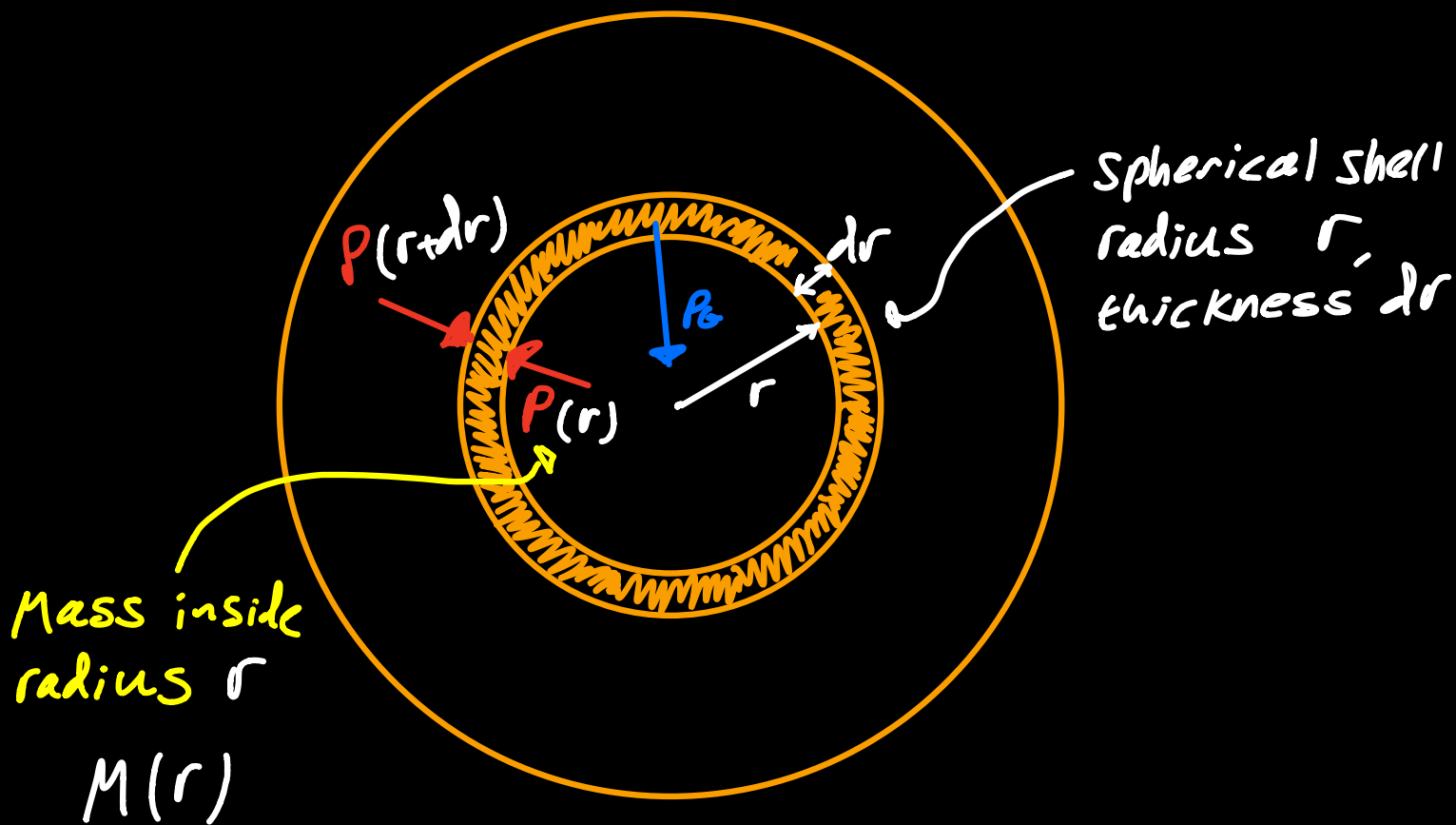
Gravitational Pressure



Most of stars life is spent in 'hydrostatic' equilibrium.

Inward Pressure = outward Pressure

$$P_G = P$$



- force pulling spherical shell inwards;

$$F_G(r) = \frac{M(r) \times \text{Mass of Shell}}{r^2}$$

$$\text{Mass of shell} = \rho(r) 4\pi r^2 dr$$

$$F_G(r) = 4\pi G M(r) \rho(r) dr$$

$$P_G(r) = \frac{F_G(r)}{\text{Area}} = \frac{F_G(r)}{4\pi r^2}$$

$$P_G(r) = \frac{G M(r) \rho(r)}{r^2} dr$$

- Net outward Thermal Pressure:

$$P(r) - P(r+dr)$$

⇒ In Hydrostatic equilibrium:

$$P(r) - P(r+dr) = \frac{G M(r) \rho(r)}{r^2} dr$$

$$\Rightarrow \frac{P(r+dr) - P(r)}{dr} = - \frac{GM(r)\rho(r)}{r^2}$$

$$\frac{dP}{dr} = - \frac{GM(r)\rho(r)}{r^2}$$

Equation of hydrostatic equilibrium

- Mass continuity equation:

$$M(r) = \int_0^r \rho(\bar{r}) 4\pi \bar{r}^2 d\bar{r}$$

$$\Rightarrow \rho(r) = \frac{1}{4\pi r^2} \frac{dM}{dr}$$

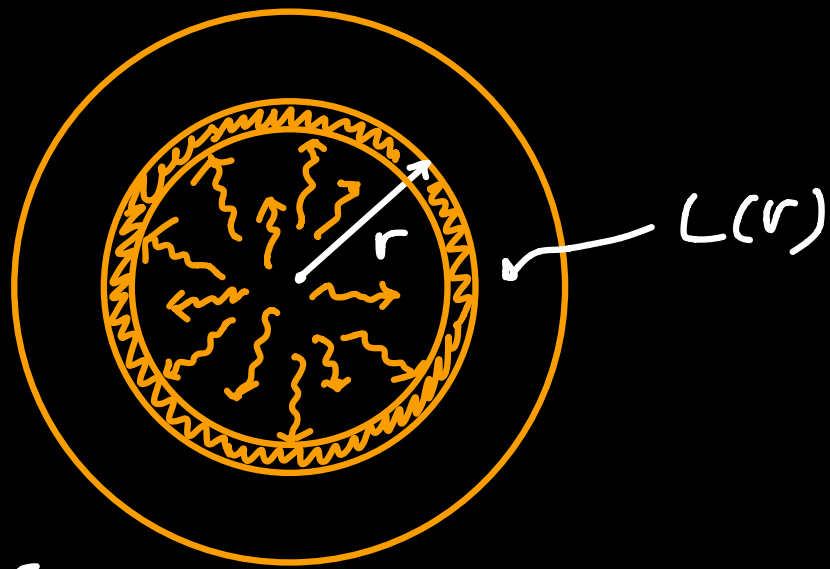
Energy Transport:

Three Mechanisms of energy transport in stars:

- 1) Radiation \rightarrow Dominates for smaller temp gradients
- 2) Convection \rightarrow Dominates for large temp gradients
- 3) Conduction \rightarrow Dominates compact stars (WD's, NS's)

Introduce $\epsilon(r)$ \rightarrow Energy generation per unit mass per unit time

$L(r)$ \rightarrow Luminosity produced inside radius r (Total power produced within radius r)



$$L(r) = \int_0^r \epsilon(\bar{r}) \times 4\pi\bar{r}^2 \rho(\bar{r}) d\bar{r}$$

$$\Rightarrow \frac{dL}{dr} = 4\pi r^2 \rho(r) \epsilon(r)$$

Energy continuity equation

— Radiation dominated energy transport :

Consider the radiation momentum flux
(radiation pressure) on a layer;

Radiation Pressure
from blackbody
source ; $P_{\text{rad}} = \frac{4}{3} \frac{\sigma T^4}{c}$

$\frac{L(r)}{c} \pi r^2 dr$ → Momentum
flux absorbed
by layer

$-\frac{dP_{\text{rad}}}{dr} 4\pi r^2 dr$ → Momentum flux
deposited to
layer

$$\frac{L(r)}{c} \pi r^2 dr = -\frac{dP_{\text{rad}}}{dr} 4\pi r^2 dr$$

$$\Rightarrow L(r) = - \frac{4\pi C r^2}{\kappa \rho(r)} \frac{dP_{\text{rad}}}{dr}$$

$$= - \frac{4\pi C r^2}{\kappa \rho(r)} \frac{d}{dr} \left(\frac{4}{3} \frac{\sigma T^4}{c} \right)$$

$$\Rightarrow \frac{dT}{dr} = - \frac{3 \kappa \rho(r) L(r)}{64 \pi r^2 \sigma T^3}$$

Radiative Energy transport
Equation.

Scaling Relations:

Equations so far:

$$(1) \quad \frac{dP}{dr} = - \frac{GM(r)\rho(r)}{r^2}$$

$$(2) \quad \rho(r) = \frac{1}{4\pi r^2} \frac{dM}{dr}$$

$$(3) \quad \frac{dL}{dr} = 4\pi r^2 \rho(r) \epsilon(r)$$

$$(4) \quad \frac{dT}{dr} = - \frac{3\tau \rho(r) L(r)}{64\pi r^2 \sigma T^3}$$

Ideal gas
equation

$$P U = N K_B T$$

$$\Rightarrow P \propto M \frac{T}{R^3}$$

$$\underline{(1)} \Rightarrow \frac{P}{R} \propto \frac{M^2}{R^5}$$

$$\Rightarrow P \propto \frac{M^2}{R^4}$$

for both to be consistent

$$\Rightarrow M \frac{T}{R^3} \propto \frac{M^2}{R^4}$$

$$\Rightarrow T \propto \frac{M}{R}$$

$$\underline{(4)} \Rightarrow \frac{T}{R} \propto \frac{\frac{M}{R^3} \times L}{R^2 T^3}$$

$$\Rightarrow L \propto \frac{R^4 T^4}{M} \propto \frac{R^4 \left(\frac{M}{R}\right)^4}{M}$$

$$\Rightarrow L \propto M^3$$

• Main Sequence Lifetime:

$$\tau_{MS} \propto \frac{M}{L} \propto M^{-2}$$

Main Sequence
lifetime

\Rightarrow More massive
stars die
younger

• Estimating Core Temperature:

$$\frac{dP}{dr} = - \frac{GM(r)\rho(r)}{r^2}$$

$$\frac{dM}{dr} = 4\pi r^2 \rho(r)$$

$$\Rightarrow \frac{dP}{dM} = - \frac{GM(r)}{4\pi r^4}$$

$$\Rightarrow \int_{P_c}^{P_s} dP = - \int_0^{M_0} \frac{GM(r)}{4\pi r^4} dM$$

$$\Rightarrow P_c = - \int_0^{M_s} \frac{GM(r)}{4\pi r^4} dM \geq \int_0^{M_s} \frac{GM}{4\pi R_s^4}$$

$$\rho_c \gtrsim \frac{GM_s^2}{8\pi R_s^4}$$

Ideal gas equation

$$\rho = \frac{\rho}{\bar{M}} k_B T$$

\leftarrow Mean mass
of gas constituents

$$\bar{M} = \frac{1}{2} M_H \text{ for fully ionized Hydrogen Plasma}$$

$$\rho = \frac{M_s}{\frac{4}{3}\pi R_s^3} \quad (\text{assuming negligible mass for electrons})$$

(Heavy approximations...)

$$T_c \gtrsim \frac{GM_s^2}{8\pi R_s^4} \frac{\bar{M} \frac{4}{3}\pi R_s^3}{M_s k_B}$$

$$T_c \sim \frac{GM_S \bar{M}}{6 k_B R_S}$$

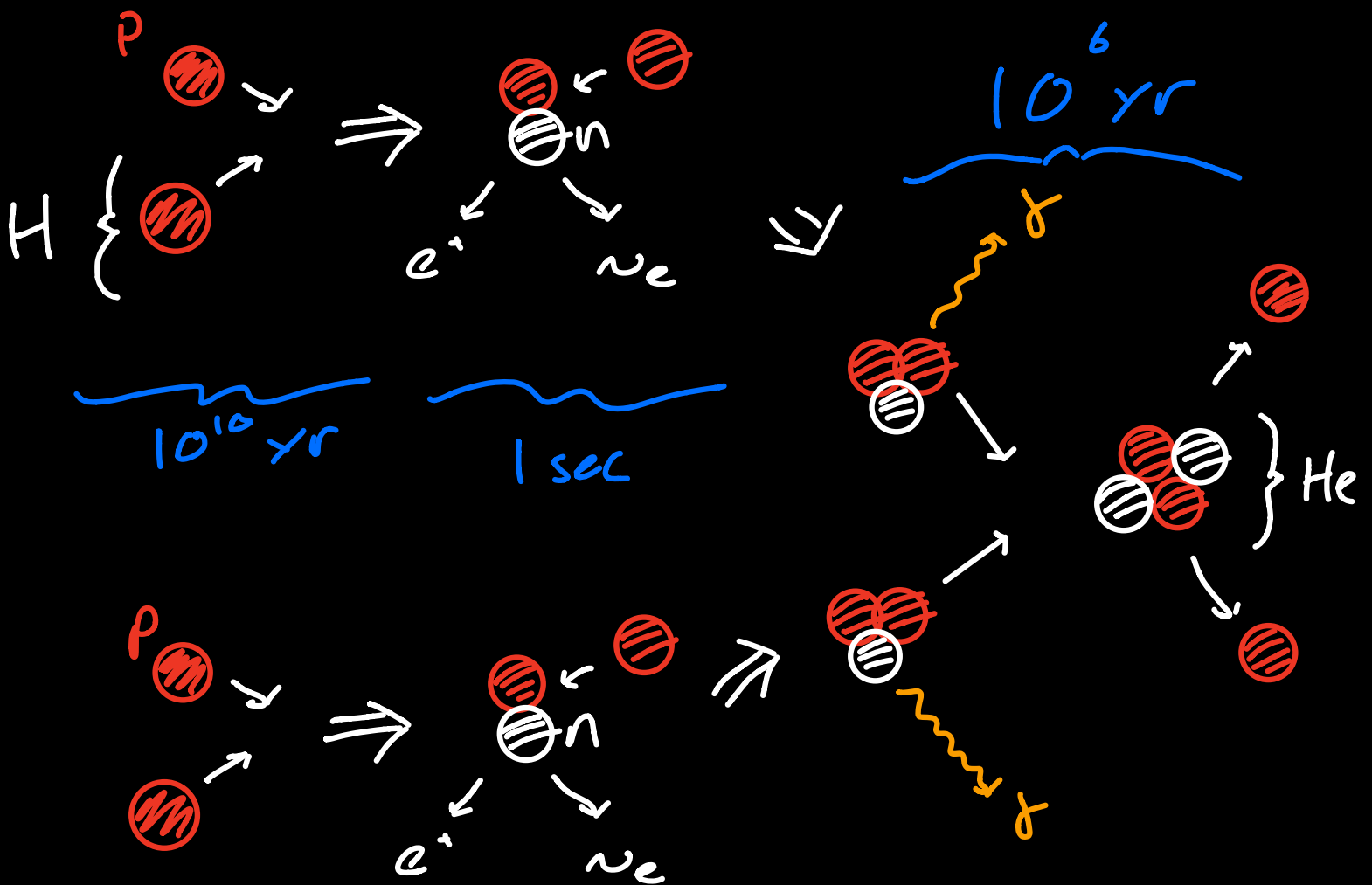
$$= \frac{(6.67 \times 10^{-11}) \times (1.99 \times 10^{30}) \times \left(\frac{1}{2} \times 1.67 \times 10^{-27}\right)}{6 \times (1.38 \times 10^{-23}) \times (6.96 \times 10^8)}$$

$$\Rightarrow T_c \sim 2 \times 10^6 \text{ K}$$

More sophisticated estimates give around 15 million Kelvin. So not bad for 'back of the envelope'!

Energy Generation in the Sun:

Energy generated by Nuclear fusion in the stars core. For our Sun this is mostly Hydrogen fusion via the PP-1 chain (PP-1 accounts for ~85%) energy generation in the Sun.



Net reaction:



Nuclear
Binding
Energy

$$\Delta E = \Delta M c^2$$

$$\Delta M = 4M_H - M_{He} \approx 4.77 \times 10^{-29} \text{ kg}$$

$$\Rightarrow \Delta E = 4.3 \times 10^{-12} \text{ J} = 27 \text{ MeV}$$

- Practicalities of solar fusion:

Protons must be brought close enough so that the strong force can take over. ($\sim 2 \text{ fm} = 2 \times 10^{-15} \text{ m}$)

$$\text{Coulomb repulsion} \sim \frac{e^2}{4\pi\epsilon_0 (2 \times 10^{-15})}$$

$$\approx 1.15 \times 10^{-13} \text{ J} = 0.72 \text{ MeV}$$

Recall, $T_c \approx 2 \times 10^6 \text{ K}$

$$\Rightarrow \text{Kinetic energy } \mathcal{J} \text{ of nuclei} = \frac{3}{2} k_B T_c$$

$$\approx 4.14 \times 10^{-17} \text{ J} = 0.00026 \text{ MeV}$$

Energy needed
to overcome
Coulomb force

$$\sim 0.72 \text{ MeV}$$

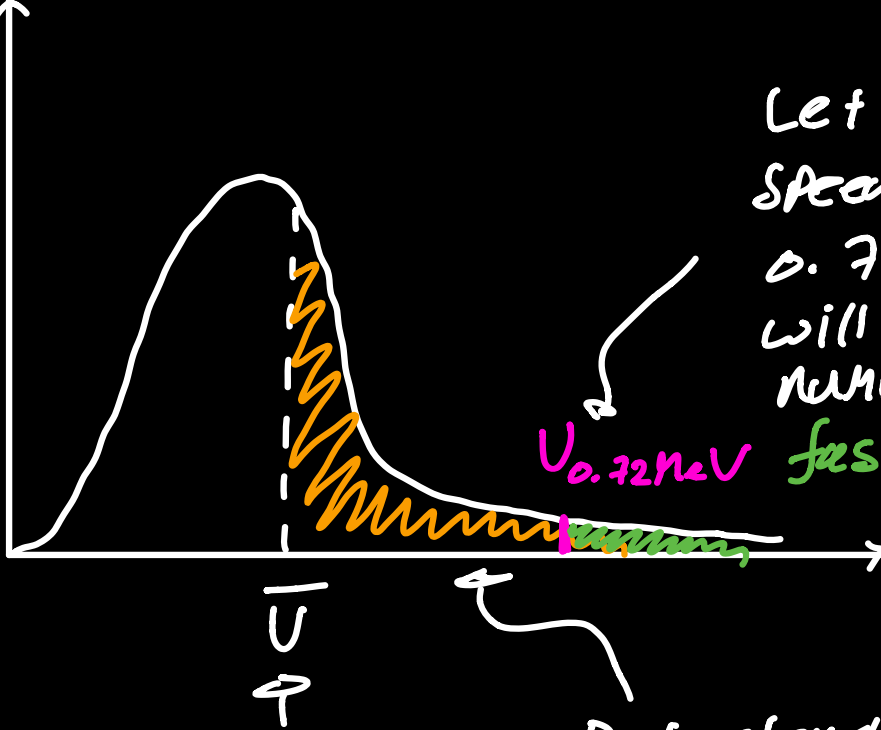
Energy of
nuclei in core

$$\sim 0.00026 \text{ MeV}$$

So why does fusion happen;

1. Boltzmann tail.
2. Quantum tunneling.

1. $f(u)$

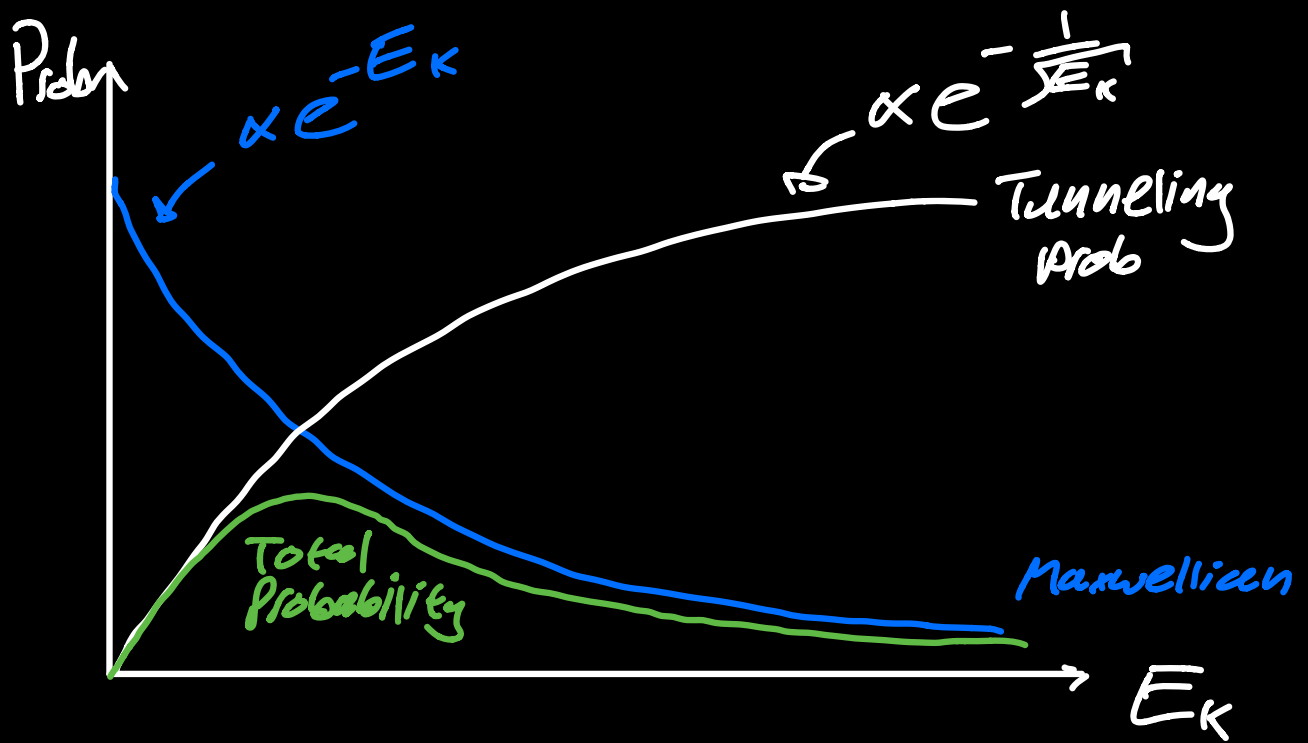


Let this be the speed with KE 0.72 MeV. There will be a small number of particles faster than this speed, which will fuse

This is what we calculated above

But clearly a number of nuclei have velocities higher than this.

2. Nuclei can tunnel through the coulomb barrier quantum mechanically, with some tunnelling probability $P_{\text{tunn}} \sim e^{-\frac{1}{\sqrt{E_k}}}$



Quantum mechanical tunneling dominates.

The sun wouldn't shine without quantum mechanics.